Dispersing billiards with cusps and tunnels

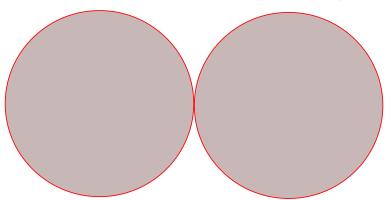
Péter Bálint work in progress with N. Chernov and D. Dolgopyat

Institute of Mathematics
Budapest University of Technology and Economics

HDSS, Corinaldo, June 1, 2010

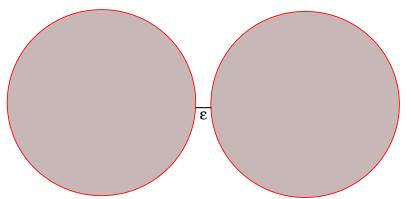
In a nutshell

- Billiards with cusps: slow decay of correlations, non-standard limit theorem;
- Billiards with tunnels: CLT, but variance blows up as $\varepsilon \to 0$



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- Billiards with cusps: slow decay of correlations, non-standard limit theorem;
- Billiards with tunnels: CLT, but variance blows up as $\varepsilon \to 0$.



Outline

Known results

Dispersing billiards in 2D Dispersing billiards with cusps

New "results"

Cusp case Tunnel case

Skeletons of arguments

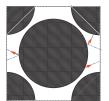
Skeleton for cusp Skeleton for tunnel

Some words on the phenomena

Rough description for cusp Rough description for tunnel Known results

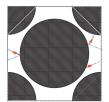
Billiards

- Billiard flow : $S^t : \mathcal{M} \to \mathcal{M}$, $(q, v) \in \mathcal{M} = Q \times S^1$, |v| = 1Uniform motion within Q, elastic reflection at the boundaries
- Billiard map phase space: $M = \bigcup_{k=1}^{K} M_k$
- $(r,\phi) \in M_k$, r: arclength along ∂C_k , $\phi \in [-\pi/2,\pi/2]$
- invariant measure $d\mu = c \cos\phi dr d\phi$



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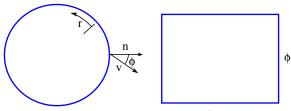
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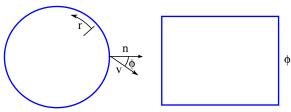
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- Billiard map is ergodic, K-mixing (Sinai '70)
- EDC: $f,g:M\to\mathbb{R}$ Hölder continuous, $\int f d\mu = \int g d\mu = 0$ let $C_n(f,g) = \mu(f\cdot g\circ T^n)$, then $|C_n(f,g)| \leq C\alpha^n$ for suitable C>0 and $\alpha<1$
 - Young '98 tower construction with exponential tails,
 - Chernov & Dolgopyat '06 standard pairs
- CLT: let $S_n f = f + f \circ T + ... + f \circ T^{n-1}$, then $\frac{S_n f}{\sqrt{n}} \stackrel{\mathcal{D}}{\Longrightarrow} \mathcal{N}(0, \sigma)$ where $\sigma = \int f^2 d\mu + 2 \sum_{n=1}^{\infty} C_n(f, f)$.
- Billiard flow: $F, G : \mathcal{M} \to \mathbb{R}, C_t(F, G)$:
 - stretched exponential bound, Chernov '07 (approximate Markov partitions)

 C_k are C^3 smooth and disjoint (no corner points); finite horizon: flight length uniformly bounded from above

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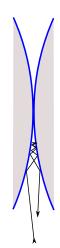
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Cusp map



 C_1 and C_2 touch tangentially – unbounded series of consecutive reflections in the vicinity of the cusp

- Reháček '95 ergodicity
- Machta '83 numerics and heuristic reasoning for $C_n(f,g) \simeq 1/n$
- Chernov & Markarian '07: $C_n(f,g) \le C \frac{\log^2 n}{n}$
- Chernov & Zhang '08: $C_n(f,g) \leq C_n^{\frac{1}{n}}$

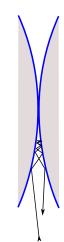
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long collision series near the cusp correspond to bounded flow time – flow mixes faster?

Melbourne & B. '08

- C_t(F, G) decays faster than any polynomial
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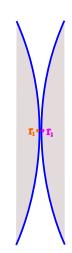
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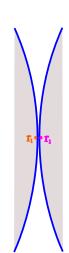
"Result" (C)

• Let
$$I_f = \int\limits_{-\pi/2}^{\pi/2} (f(r_1,\phi) + f(r_2,\phi)) \rho(\phi) d\phi$$

with $\rho(\phi) = \frac{\sqrt{\cos\phi}}{\int\limits_{-\pi/2}^{\pi/2} \sqrt{\cos\phi} d\phi}$

- if $I_f \neq 0$ then $\frac{S_n f}{\sqrt{n \log n}} \stackrel{\mathcal{D}}{\Longrightarrow} \mathcal{N}(0, D_f)$ where $D_f = c^* I_f^2$ and c^* is some numerical constant.
- if $I_f = 0$ then $S_n f$ satisfies standard CLT.



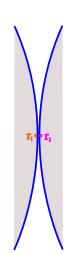


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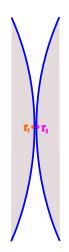
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Remarks concerning the cusp flow

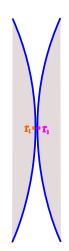


• if $G: \mathcal{M} \to \mathbb{R}$ Hölder, then let $\tau(x)$

$$g(x) = \int_{0}^{\tau(x)} G(x,t) dt,$$

- and we have $I_g = 0$ (as $\tau(x) = 0$ for $x = (r_1, \phi)$),
- hence CLT and invariance principle are reasonable.

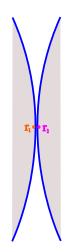
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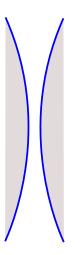
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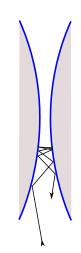
- for fixed $\varepsilon > 0$ this is a Sinai billiard, hence *CLT*:
- $\frac{S_n f}{\sqrt{n}} \stackrel{\mathcal{D}}{\Longrightarrow} \mathcal{N}(0, D_{f,\varepsilon})$ with
- $D_{f,\varepsilon} = D_f |\log \varepsilon| (1 + o(1))$



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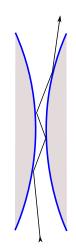
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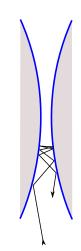
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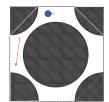


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1. Brownian Brownian motion – Chernov & Dolgopyat '09



 $m \ll M$ (separation of time scales)

Motivation

1. Brownian Brownian motion - Chernov & Dolgopyat '09



 $m \ll M$ (separation of time scales) SDE for large particle:

 $dV = \sigma_{Q}(f)dW$

collisions of the heavy particle with the wall?

2. Triangular lattice with small opening

How does the planar

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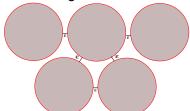
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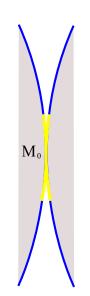
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How does the planar diffusion depend on ε ?

The first return map



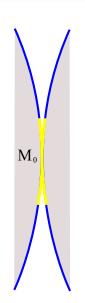
Let $\hat{M} = M \setminus M_0$ where M_0 is a fixed small nbd. of the cusp.

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limit law for $\hat{S}_n \hat{f}$ implies limit law for $S_n f$ (eg. Gouëzel '04)

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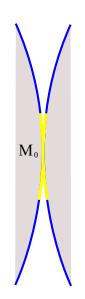


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The map $\hat{T}: \hat{M} \to \hat{M}$ is uniformly hyperbolic and it satisfies the Growth Lemma ("Expansion prevails fractioning")

so that

- Young tower with exponential tails can be constructed
- standard pairs can be coupled at an exponential rate

Hence: EDC for Hölder observables

Lemma (C2)

 $|\hat{\mu}(\hat{f} \cdot \hat{f} \circ \hat{T}^n)| \le Ce^{-\alpha n}$ with $C > 0, \alpha < 1$ for $n \ge 1$ Not for n = 0 as \hat{f} is not Hölder and not in L^2

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Not for $n = 0$ as \hat{f} is not Hölder and not in L^2

- $M_n = \{x \in \hat{M} | R(x) = n\}$ *n*-cell
- $L_n = \bigcup_{i < n} M_i$ low cells, $H_n = \bigcup_{i > n} M_i$ high cells

•
$$\hat{f}|_{M_n} = nI(1 + o(1))$$

 $(recall\ I = c_1 \int\limits_{-\pi/2}^{\pi/2} (f(r_1, \phi) + f(r_2, \phi)) \sqrt{\cos(\phi)} d\phi$

- $\hat{\mu}(H_n) = \frac{c_2}{R^2}(1 + o(1))$ (here c_1 , c_2 are numerical constants)
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Blow-up of \hat{f}^2

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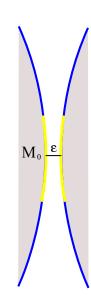
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if $\hat{f} \circ \hat{T}^n$ were i.i.d, it would belong to the non-standard domain of attraction of the normal law



 $T_{\varepsilon}: M \to M, M_0$: same nbd. for any ε ,

 $\hat{M} = M \setminus M_0$

Return map $\hat{T}_{\varepsilon}: \hat{M} \to \hat{M}$ and return time R_{ε} depend on ε

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The map $\hat{T}_{\varepsilon}: \hat{M} \to \hat{M}$ satisfies the Growth Lemma and EDC for Hölder observables uniformly in ε .

Lemma (T2)

 $|\hat{\mu}(\hat{t}_{\varepsilon} \cdot \hat{t}_{\varepsilon} \circ \hat{T}_{\varepsilon}^n)| \leq Ce^{-\alpha n}$ with $C > 0, \alpha < 1$ independent of ε

Hence CLT for $\hat{S}_n \hat{f}_{\varepsilon}$ with variance $D_{\hat{\varepsilon}} = \hat{u}(\hat{f}_{\varepsilon}^2) + \mathcal{O}(1)$:

correlations do not contribute to the main term

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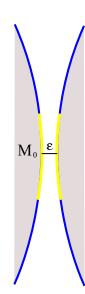
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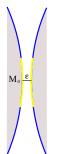
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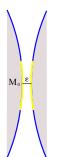
Blow-up of \hat{f}_{ε}^2



Lemma (T3)
$$\hat{\mu}(\hat{f}_{\varepsilon}^2) = |\log \varepsilon| D_{\hat{f}}(1 + o(1))$$

All these Lemmas require: detailed geometric analysis of the cells M_k (measures, unstable and stable dimensions etc...)

- For cusp, mostly (but not completely) done by Chernov & Markarian
- For tunnel, requires new ideas & technical work (in progress)



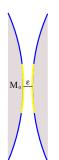
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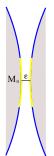
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Blow-up of \hat{f}_{s}^{2}



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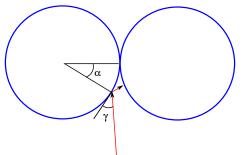
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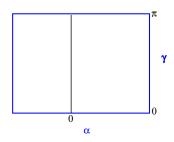
For simplicity assume that C_1 and C_2 are circles of radius 1.

- while going down the cusp: α decreases, $\gamma:0\longrightarrow\frac{\pi}{2}$
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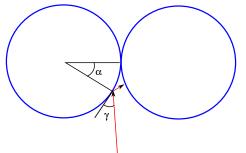
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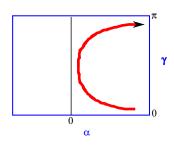
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$$\gamma' - \alpha' = \alpha + \gamma$$

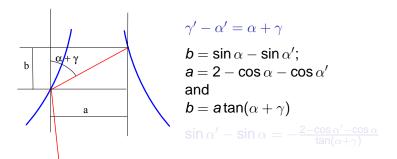
$$b = \sin \alpha - \sin \alpha';$$

$$a = 2 - \cos \alpha - \cos \alpha'$$
and
$$b = a \tan(\alpha + \gamma)$$

$$\sin \alpha' - \sin \alpha = -\frac{2 - \cos \alpha' - \cos \alpha}{2}$$

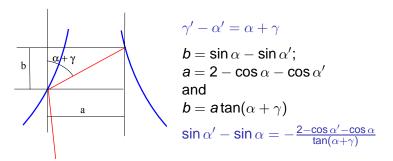
- Throughout the corner series: $\alpha \ll 1$, $\alpha < \gamma$;
- in a "large part" of the corner series: $\alpha \ll \gamma$.

$$\gamma' - \gamma \approx 2\alpha;$$
 $\alpha' - \alpha \approx -\frac{\alpha^2}{\tan(\gamma)}$



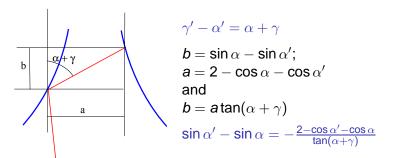
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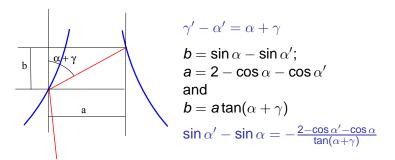
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 $J = \alpha^2 \sin \gamma$ is first integral, so $\dot{\gamma} = 2\sqrt{\frac{J}{\sin \gamma}}, \ dt = \frac{2\sqrt{\sin \gamma}}{\sqrt{J}}d\gamma$

proportion of time between γ_1 and $\gamma_2 \approx \int_{\gamma_1}^{\gamma_2} \sqrt{\sin \gamma}d\gamma.$

Recall $I_f = c \int_{-\pi/2}^{\pi/2} (f(r_1, \phi) + f(r_2, \phi)) \sqrt{\cos(\phi)}d\phi.$

length of the excursion $R = cJ^{-\frac{1}{2}} \int_{0}^{\pi} \sqrt{\sin \gamma}d\gamma = cJ^{-\frac{1}{2}}$

hence $u(H_2) = u(R > p) = u(J < \frac{c}{\gamma}) = u(\alpha^2 < \frac{c}{\gamma}) = \frac{c}{\gamma}$

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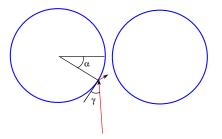
Corner series for tunnel

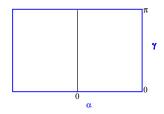
Coordinates: α, γ as for cusp

$$\gamma' - \alpha' = \alpha + \gamma$$

$$a = 2 - \cos \alpha - \cos \alpha' + \varepsilon$$

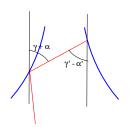
$$\sin \alpha' - \sin \alpha = -\frac{2 - \cos \alpha' - \cos \alpha + \varepsilon}{\tan(\alpha + \gamma)}$$





Corner series for tunnel

Coordinates: α, γ as for cusp



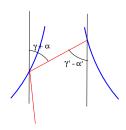
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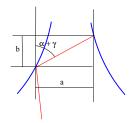
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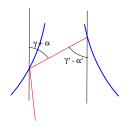


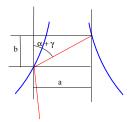
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Flow approximation for tunnel

$$\dot{\gamma} = 2\alpha; \qquad \dot{\alpha} = -\frac{\alpha^2 + \varepsilon}{\tan(\gamma)}.$$

$$J=(lpha^2+arepsilon)\sin\gamma$$
 is first integral, so $\dot{\gamma}=2lpha=\pm2\sqrt{rac{J}{\sin\gamma}-arepsilon}$.

Fix some small δ_0 . We distinguish three cases:

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Cusp case

$$J = (\alpha^2 + \varepsilon) \sin \gamma, \qquad \dot{\gamma} = 2\alpha = \pm 2\sqrt{\frac{J}{\sin \gamma} - \varepsilon}$$

 $J > \varepsilon/\delta_0$:

- $\alpha > 0$ and $\alpha^2 \gg \varepsilon$ throughout the excursion
- cusp estimates apply, however $R = CJ^{-1/2} \le \frac{C}{\sqrt{c}}$

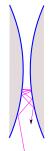
Contribution to the variance: $\hat{\mu}(\hat{f}^2 \cdot \mathbf{1}_{L_{\underline{c}}}) = D_{\hat{f}} |\log \varepsilon|$

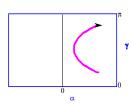
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- $\alpha > 0$ and $\alpha^2 \gg \varepsilon$ throughout the excursion
- cusp estimates apply, however $R = CJ^{-1/2} \le \frac{C}{\sqrt{\epsilon}}$

Contribution to the variance: $\hat{\mu}(\hat{f}^2 \cdot \mathbf{1}_{L_{\frac{c}{-}}}) = D_{\hat{f}} |\log \varepsilon|$





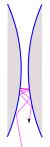
Cusp case

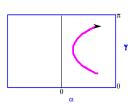
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 $\mathcal{O}(1)$ contribution to the variance.

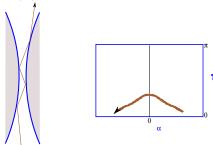
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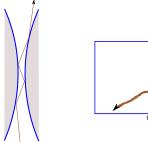


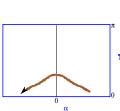
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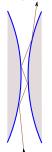
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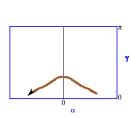
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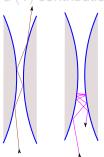
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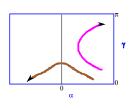




What is in between?

 $\alpha = 0, \gamma = \pi/2$ is a hyperbolic fixed point (period two orbit) Saddle case: if $J \approx \varepsilon$, R can be arbitrary large, however, it is dominated by the hyperbolic periodic orbit $\mathcal{O}(1)$ contribution to the variance

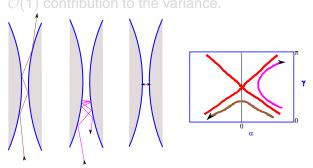




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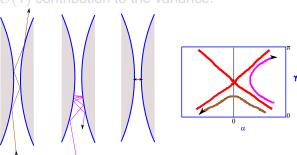
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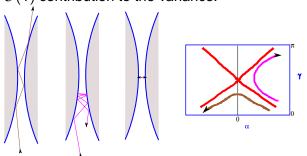
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• Cusp:
$$\frac{S_n f}{\sqrt{n \log n}} \stackrel{\mathcal{D}}{\Longrightarrow} \mathcal{N}(0, D_f)$$
 with explicit D_f

• Tunnel:
$$\frac{S_n f}{\sqrt{n}} \stackrel{\mathcal{D}}{\Longrightarrow} \mathcal{N}(0, D_{f, \varepsilon})$$
 with $D_{f, \varepsilon} = |\log \varepsilon| D_f (1 + o(1))$

Related models:

1. Infinite horizon Lorentz gas and field of strength ε

- short-range correlations
- what is ε?

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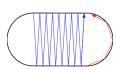
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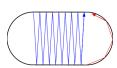
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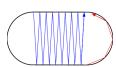


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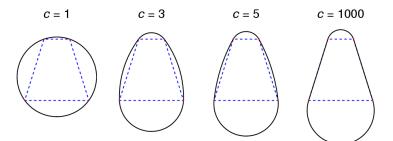
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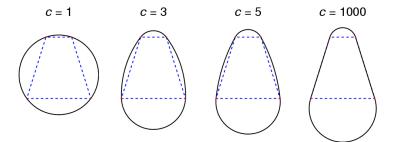


Generalized squashes



Numerics and heuristic reasoning: Ergodicity for large enough finite *c* (Halász, Sanders, Tahuilán, B.)

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