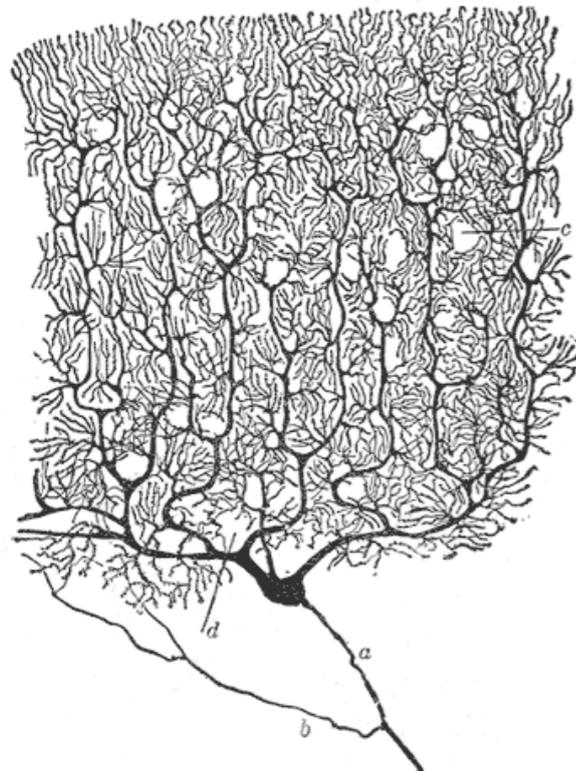
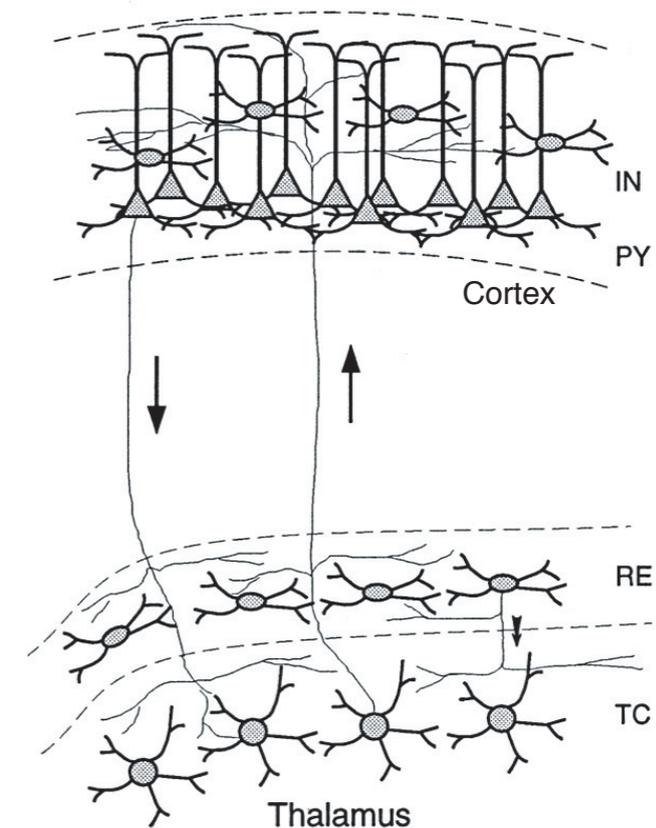


Mathematical Neuroscience: from neurons to networks



Part I

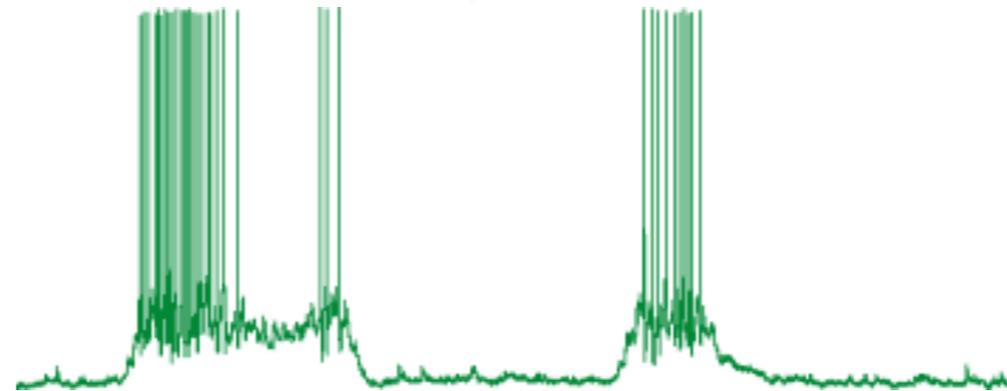
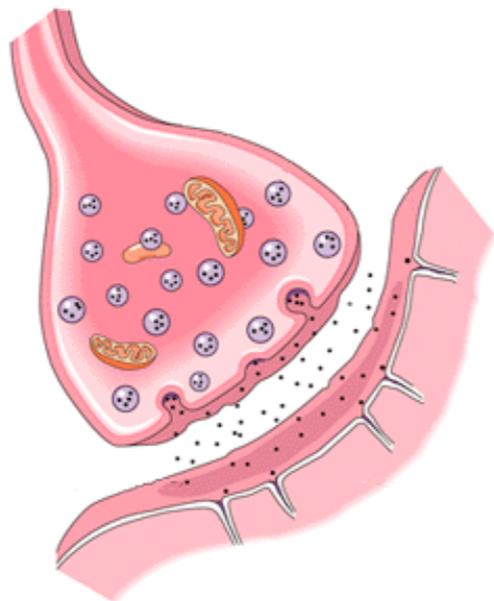
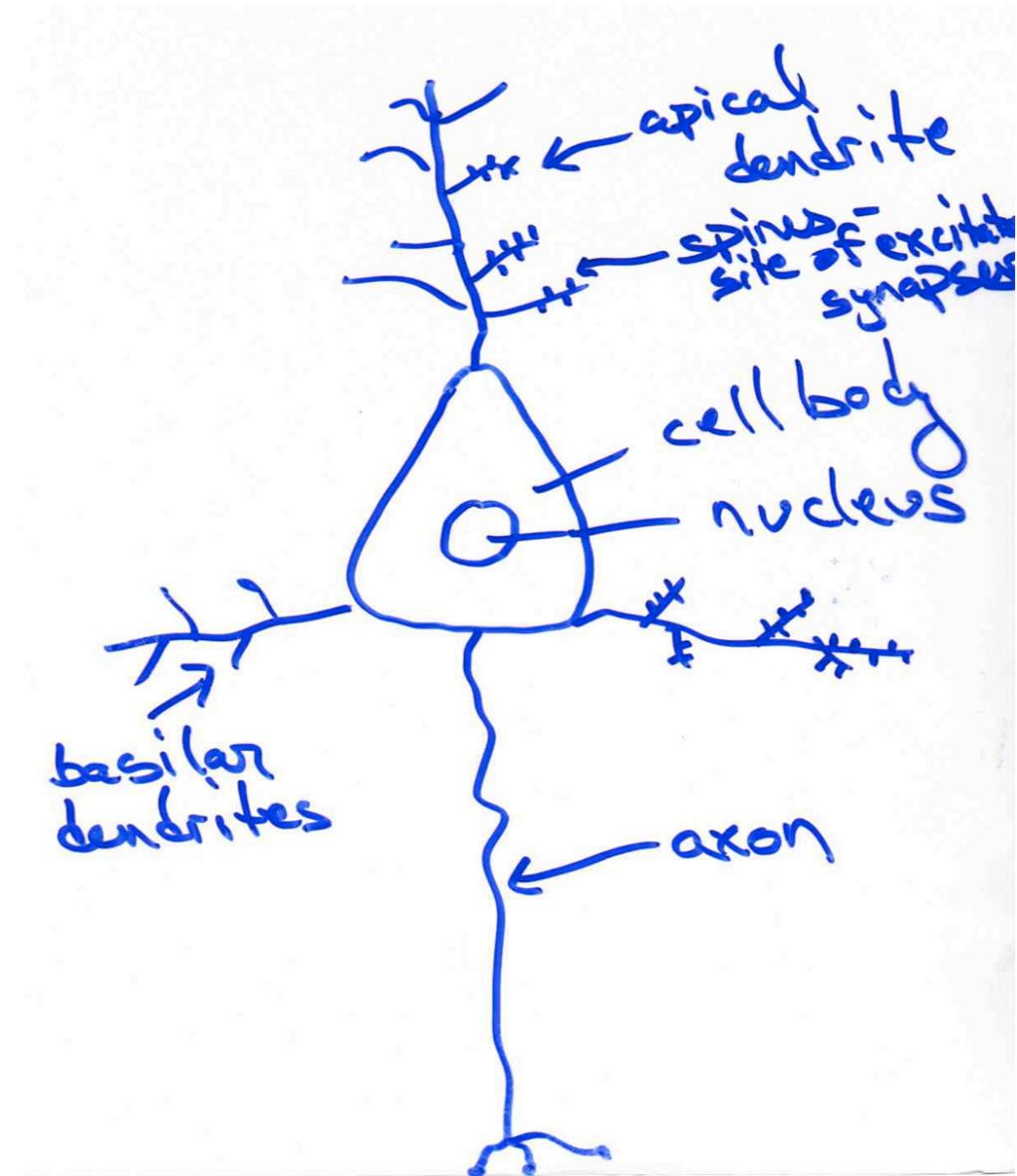
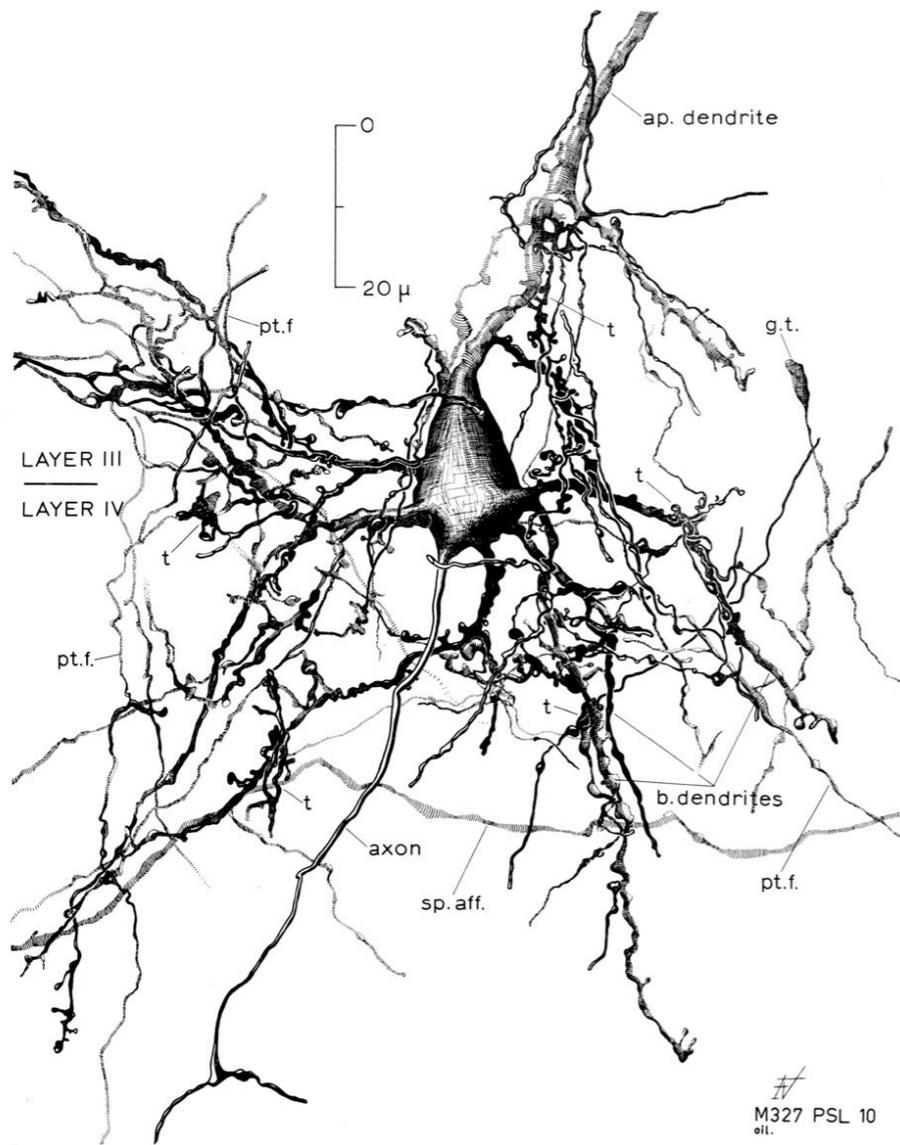
Steve
Coombes



The University of
Nottingham

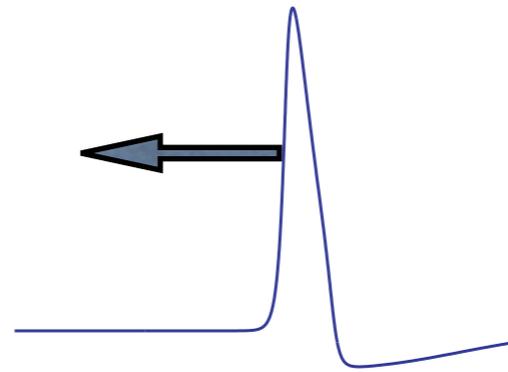
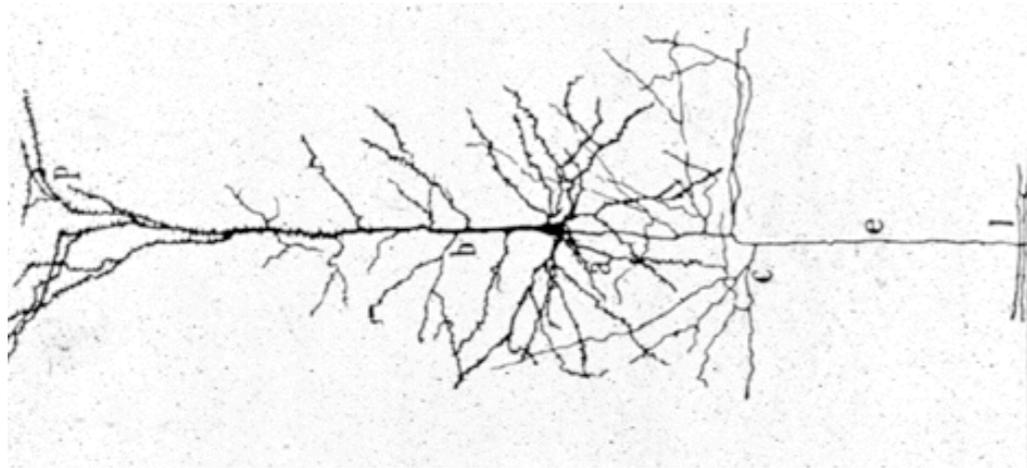
School of Mathematical
Sciences

Neurons: pyramidal cells

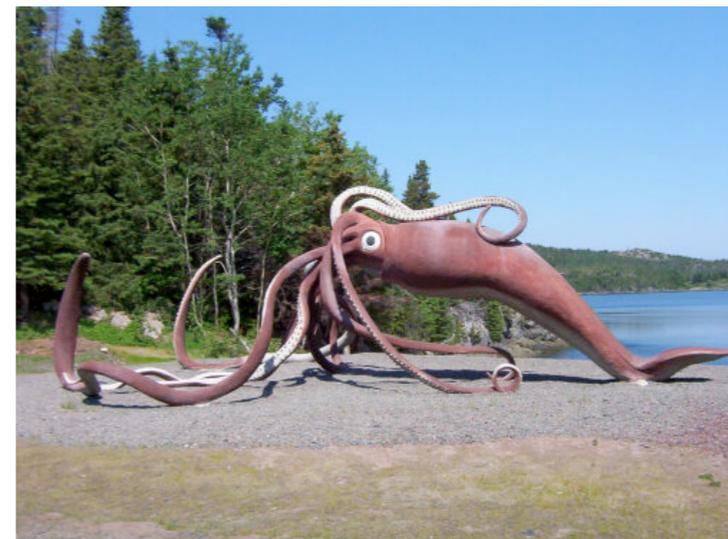
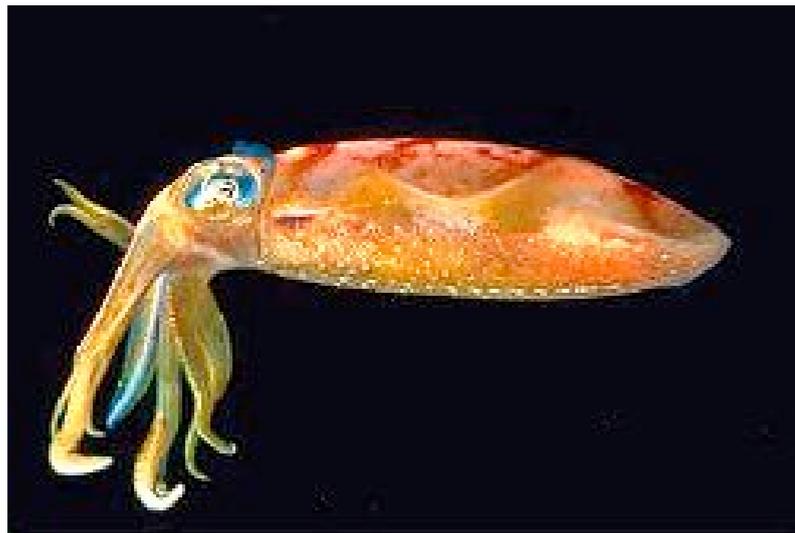


Hodgkin and Huxley

(1950s) express (and subsequently fit) the dynamics of gating variables (representing membrane channels) using the mathematical language of nonlinear ODEs.



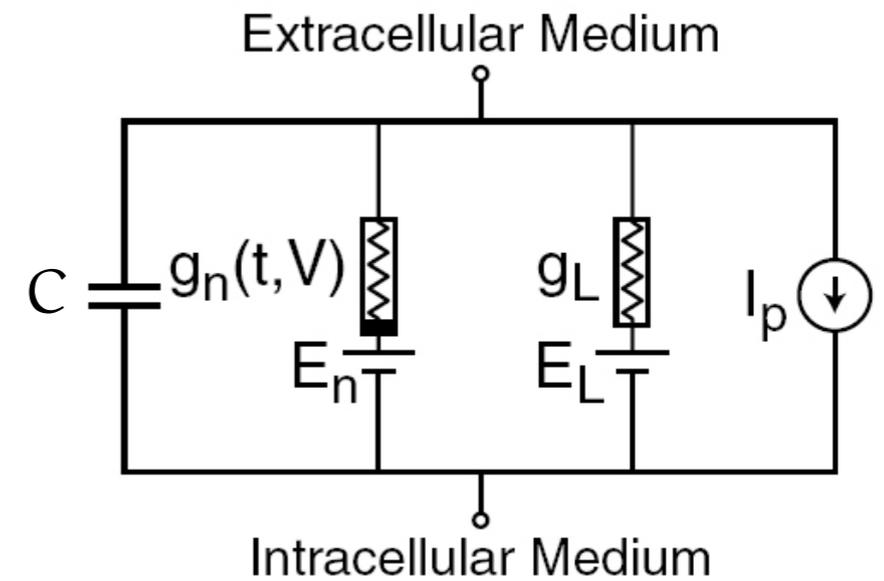
Action
potentials
m/s



Active membrane models

$$C \frac{dv}{dt} = - \sum_k g_k m_k^{p_k} h_k^{q_k} (v - v_k) + I$$

$$p_k, q_k \in \mathbb{Z}$$



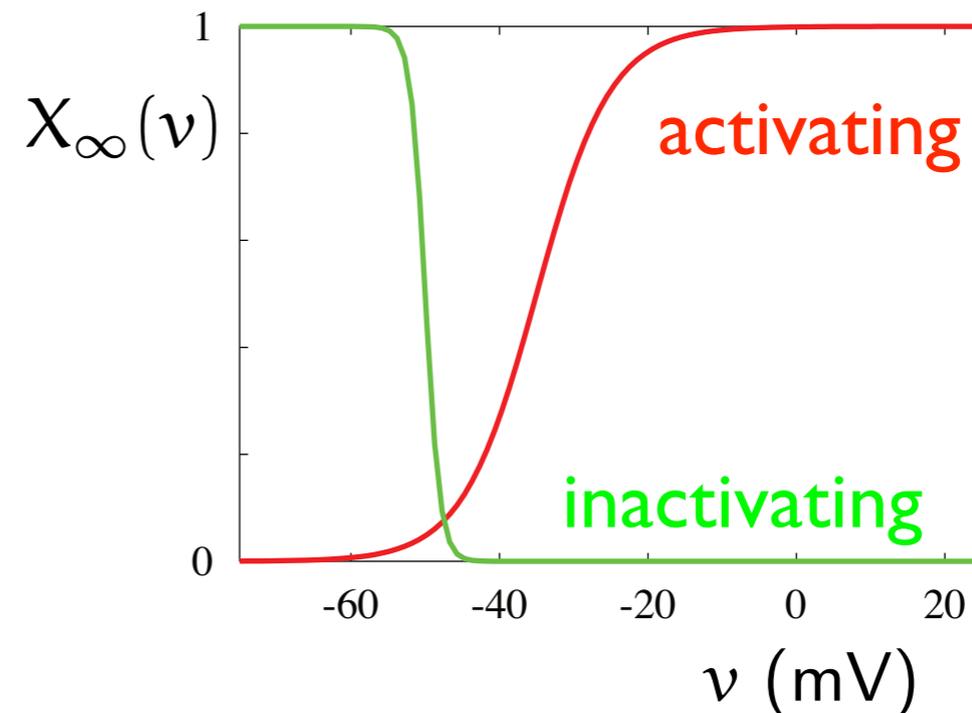
v — membrane potential

m_k, h_k — gating variables

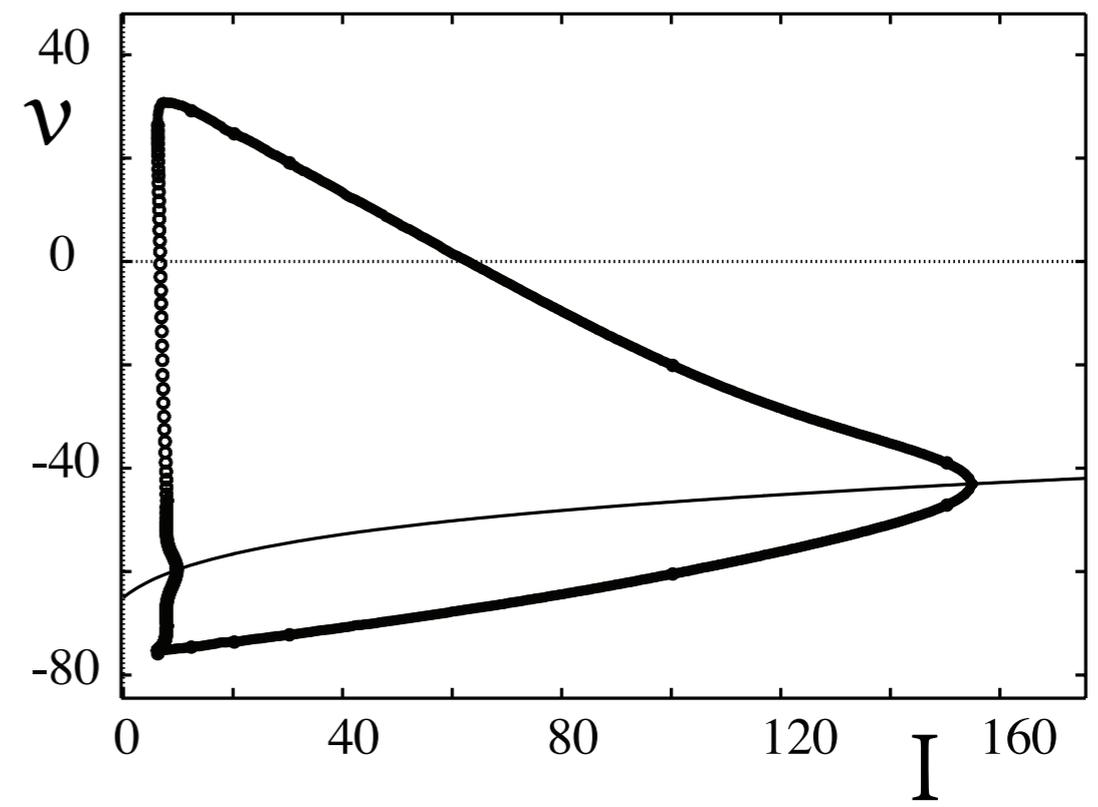
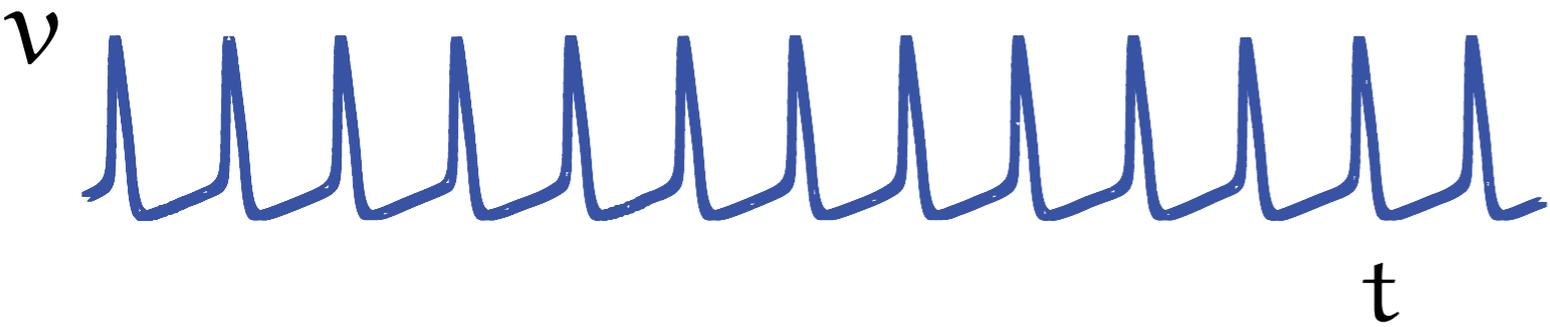
v_k — reversal potentials

g_k — conductances

$$\frac{dX}{dt} = \frac{X_\infty(v) - X}{\tau_x(v)}$$

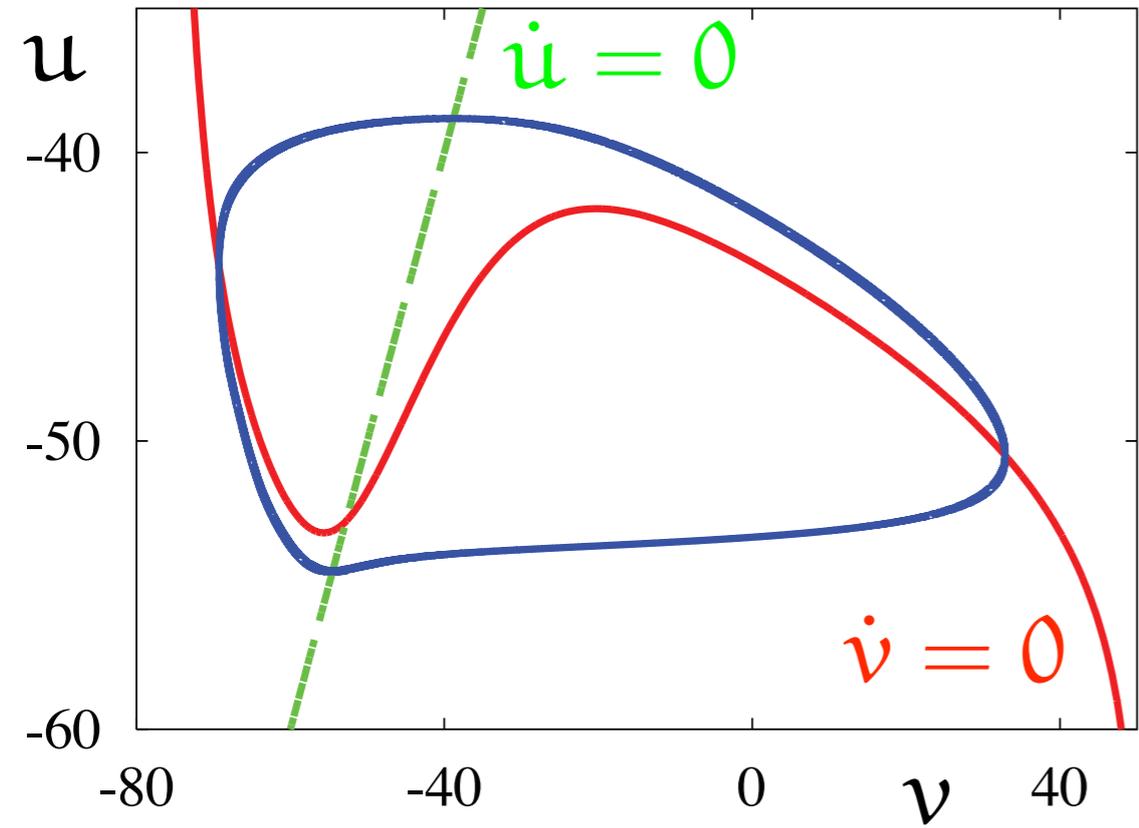


Hodgkin and Huxley: (v, m, n, h)



Reduction: $m \rightarrow m_\infty(v)$

$(n, h) \rightarrow (n_\infty(u), h_\infty(u))$

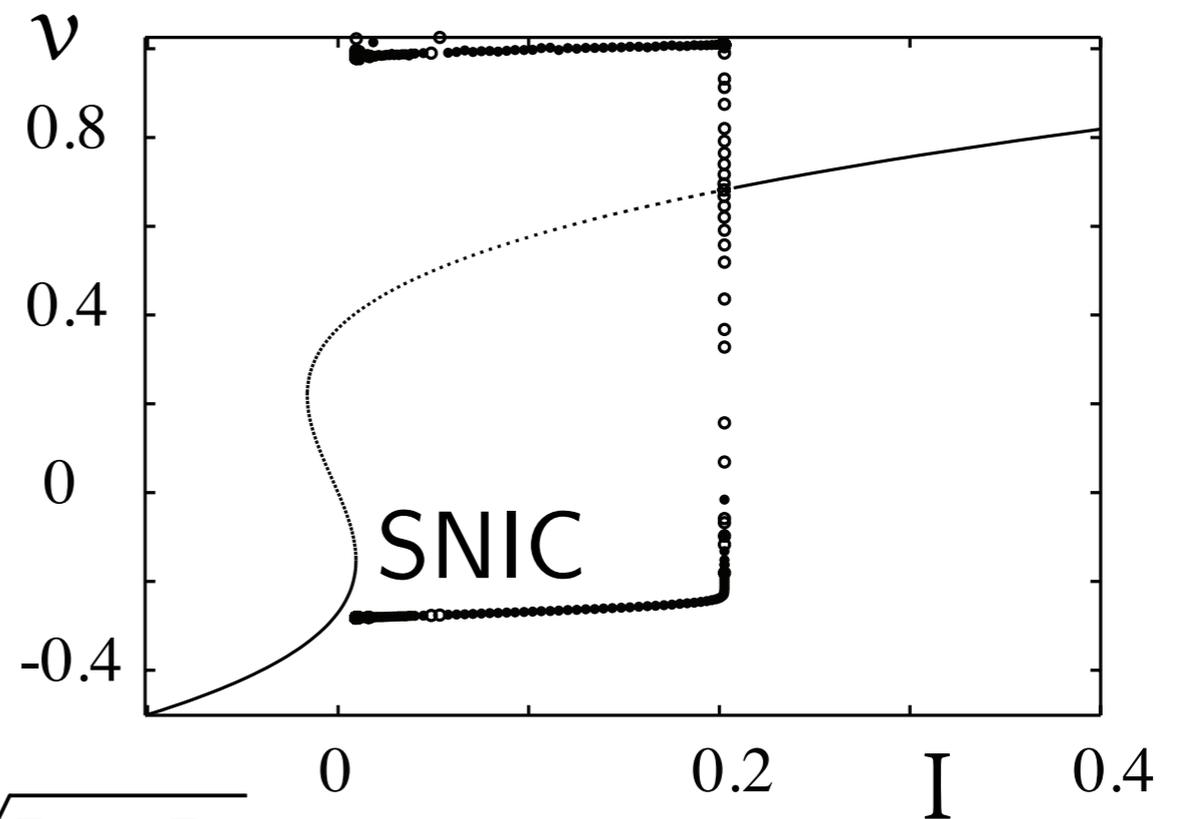
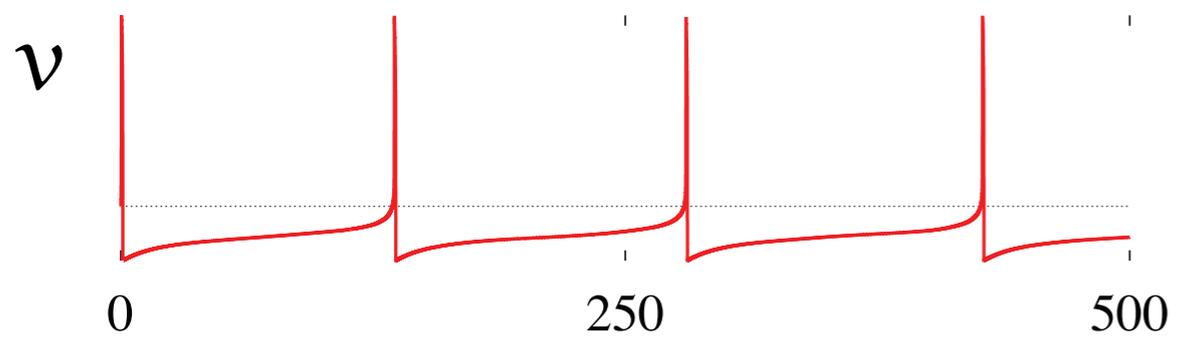
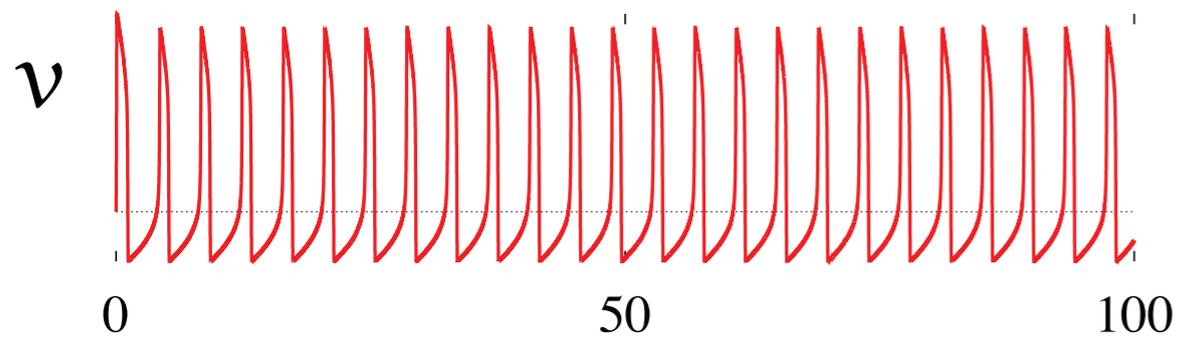
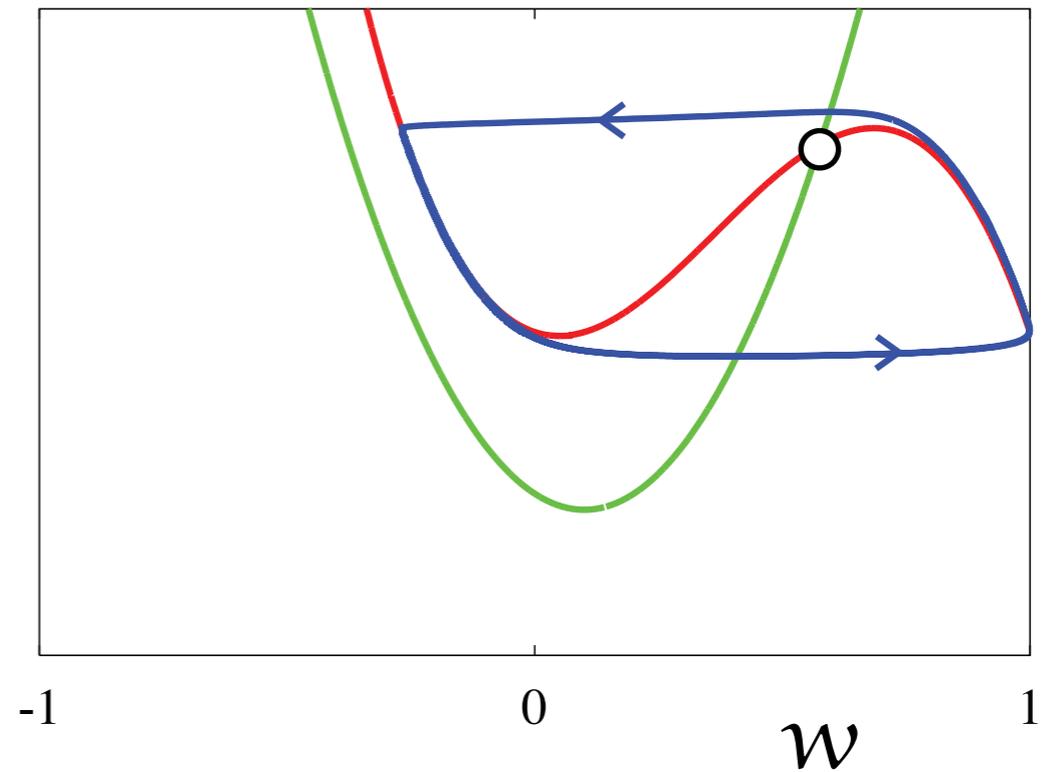
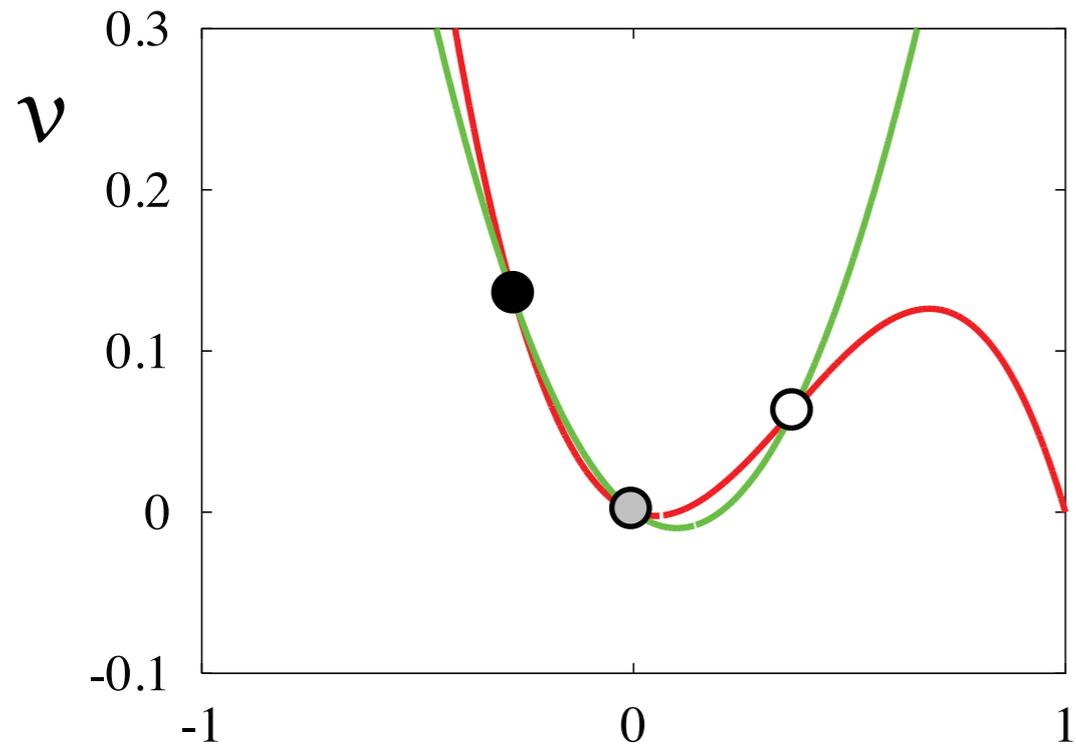


$$C \frac{dv}{dt} = f(v, u) + I$$

$$\frac{du}{dt} = g(v, u)$$

Method of equivalent potentials gives f and g in terms of HH model - Abbott and Kepler 1990

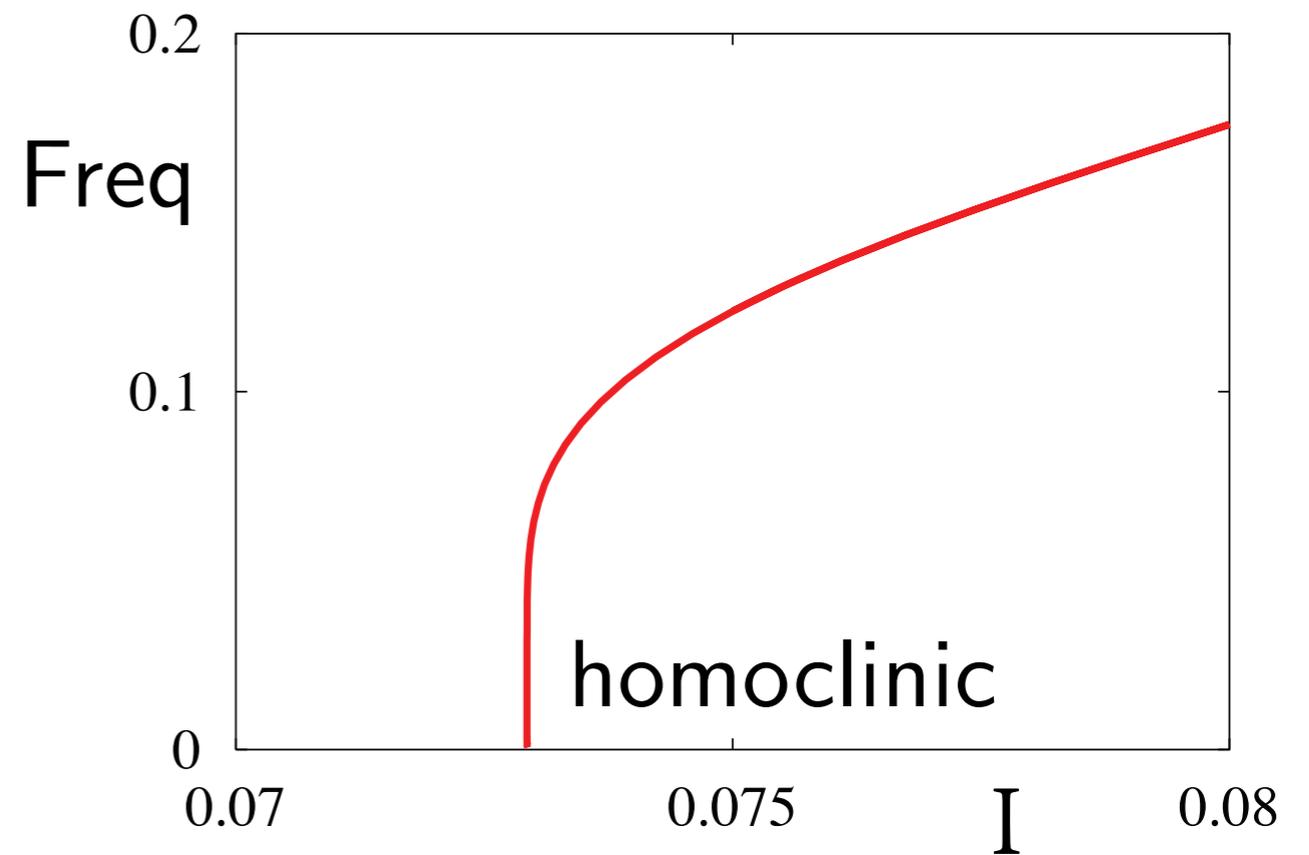
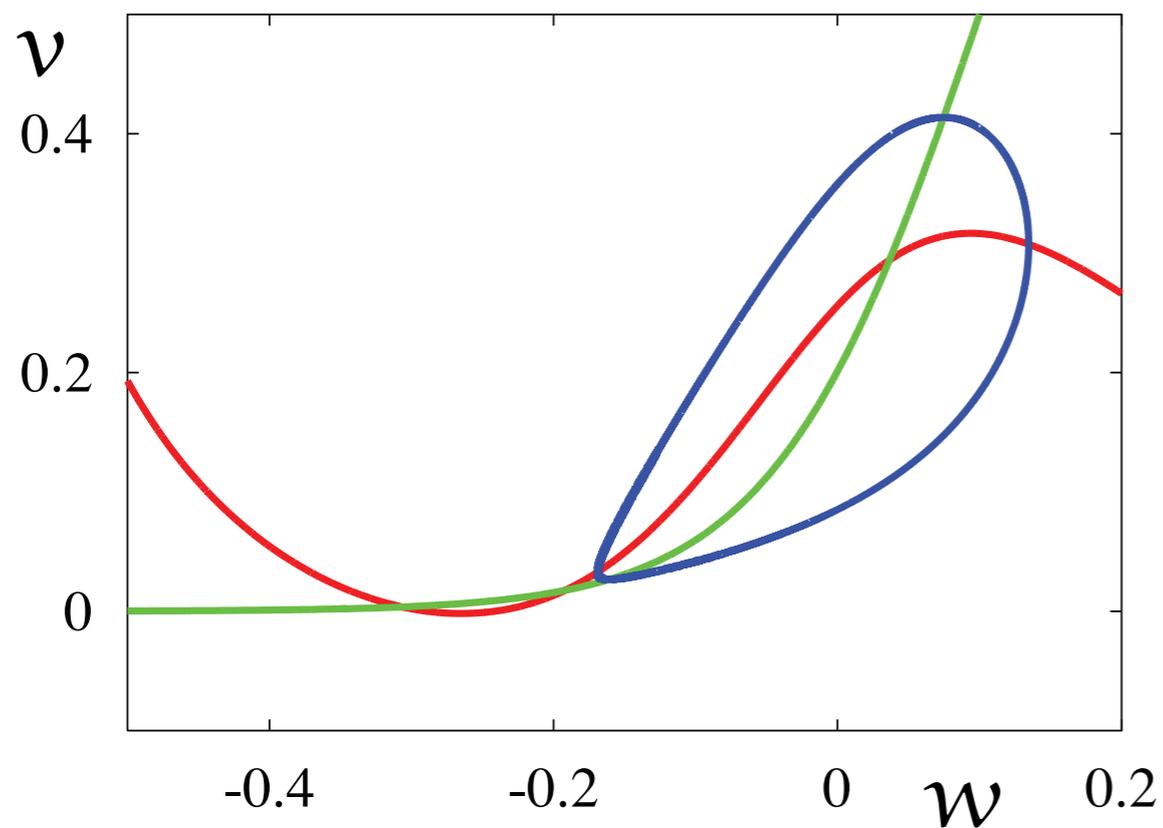
Cortical model (slow firing)



$$\text{Freq} \sim \sqrt{I - I_c}$$

Morris-Lecar model (slow firing) (v, w)

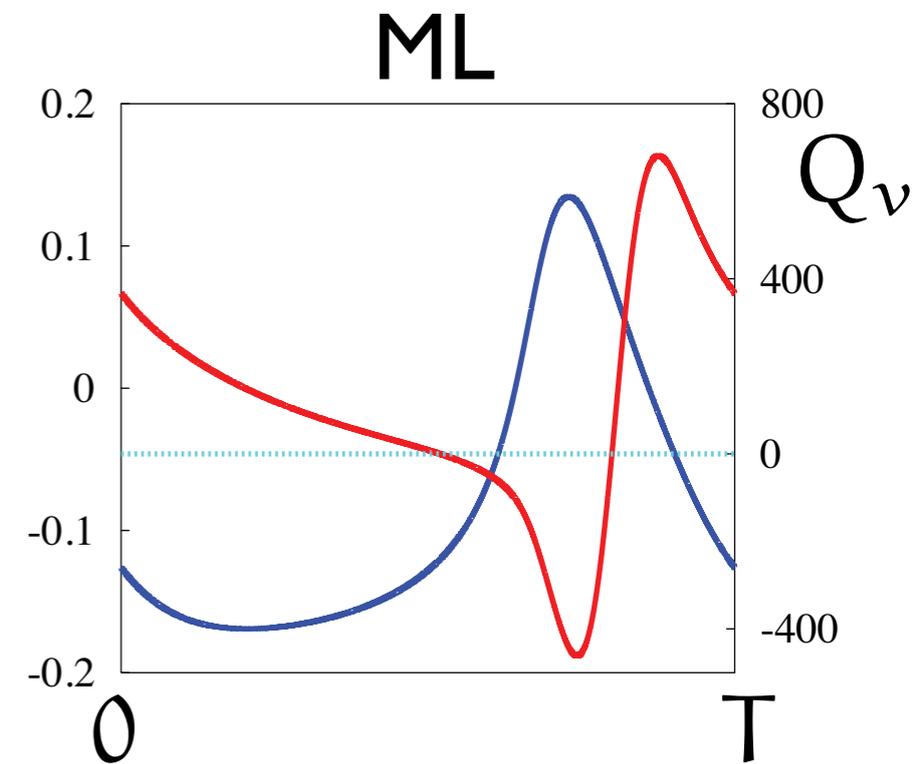
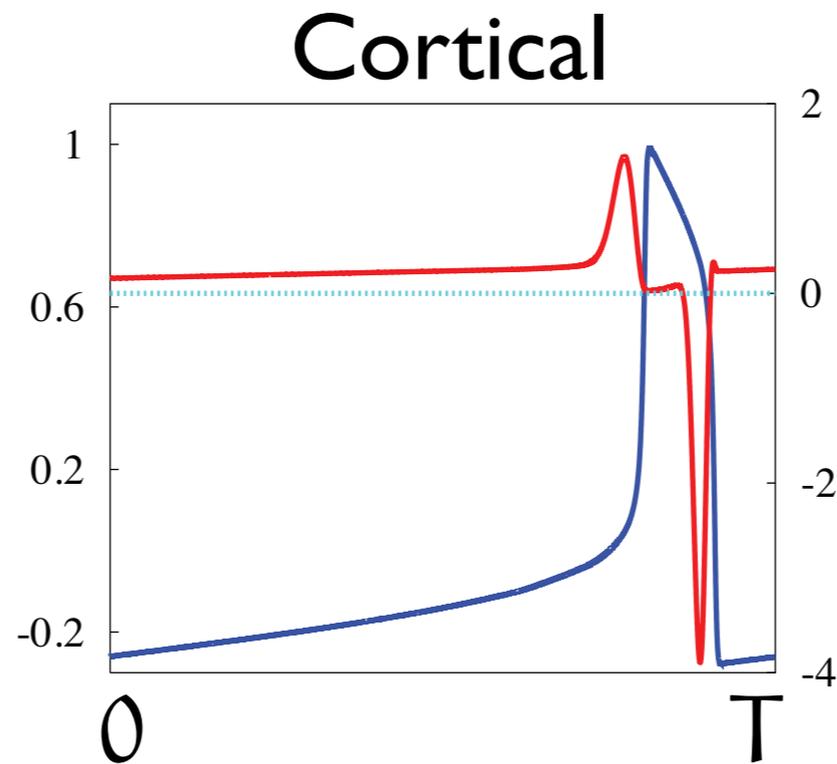
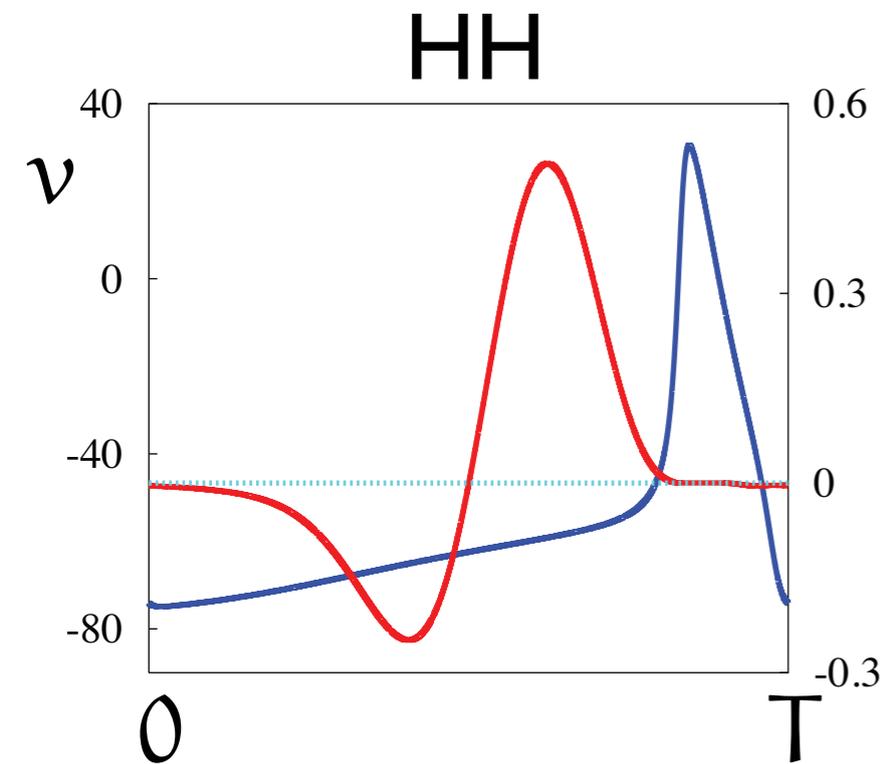
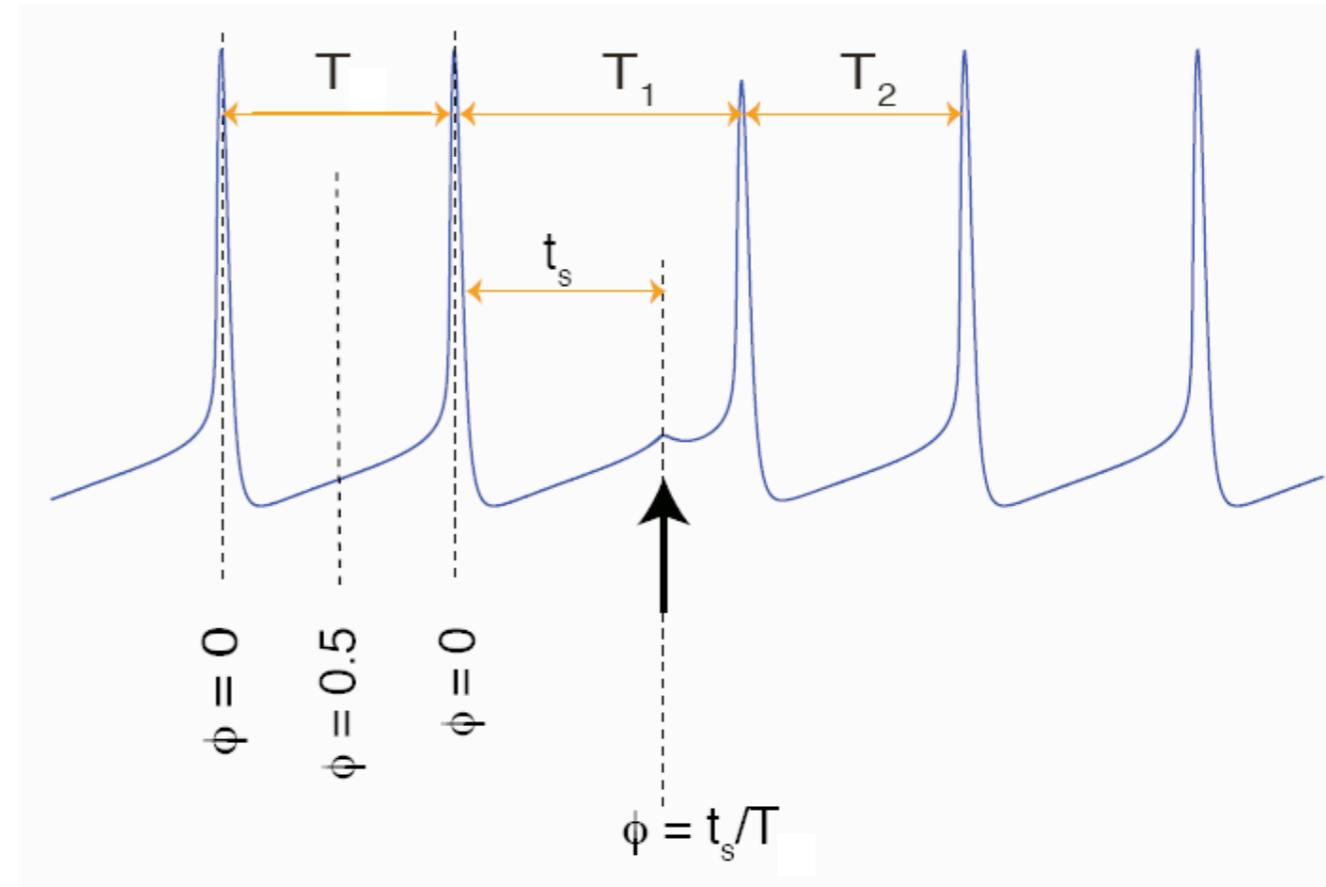
Originally a model of the barnacle giant muscle fiber



$$\text{Freq} \sim \frac{1}{\ln(I - I_c)}$$

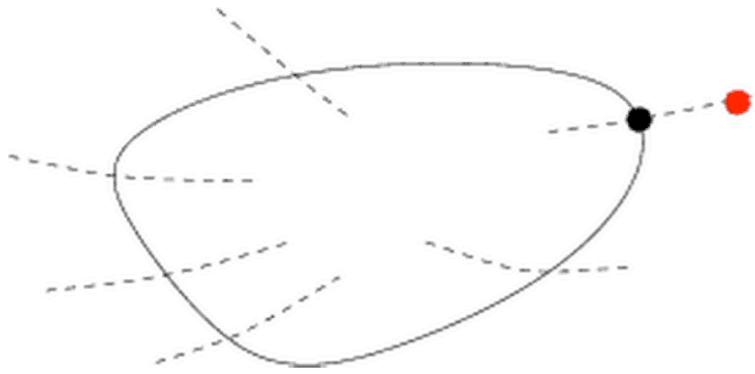
Phase Response Curve (PRC)

A **PRC** tabulates the transient change in the cycle period of an oscillator induced by a perturbation as a function of the phase at which it is received.

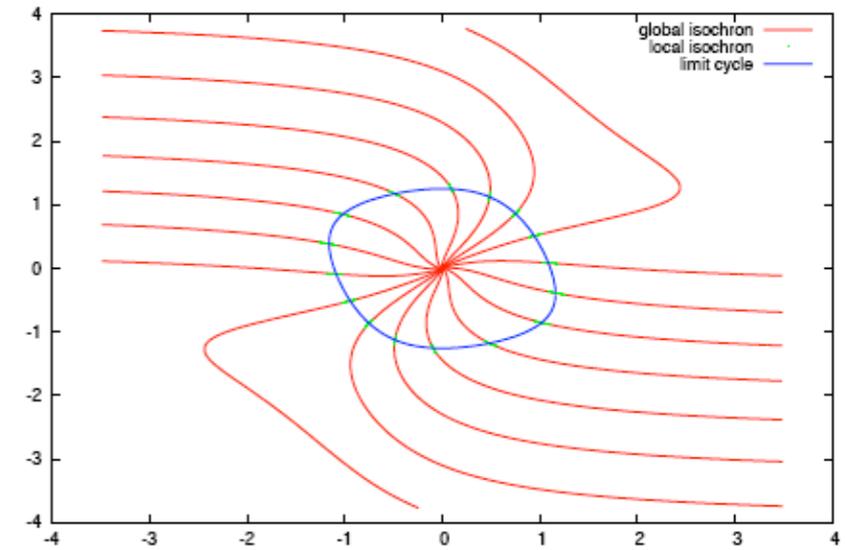


obtained numerically

$$Q = \nabla_z \theta$$



Isochrons as leaves of the stable manifold of a hyperbolic limit cycle



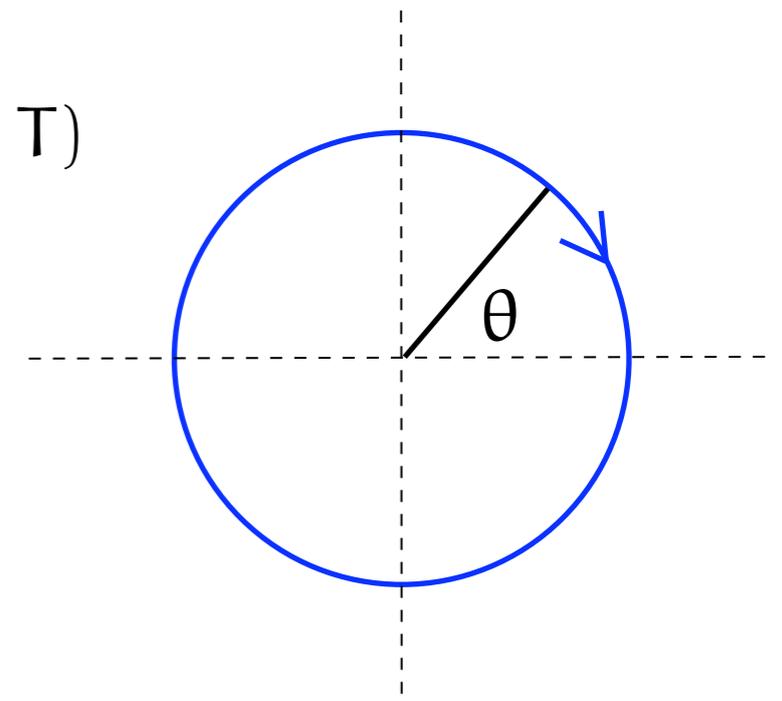
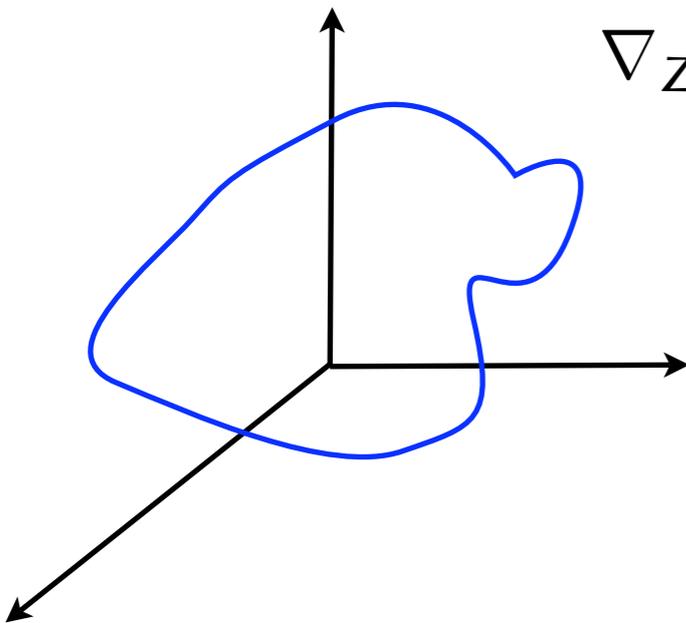
Call the orbit $z = Z(t)$ where $\dot{z} = F(z)$

Introduce a phase (isochronal coordinates) θ

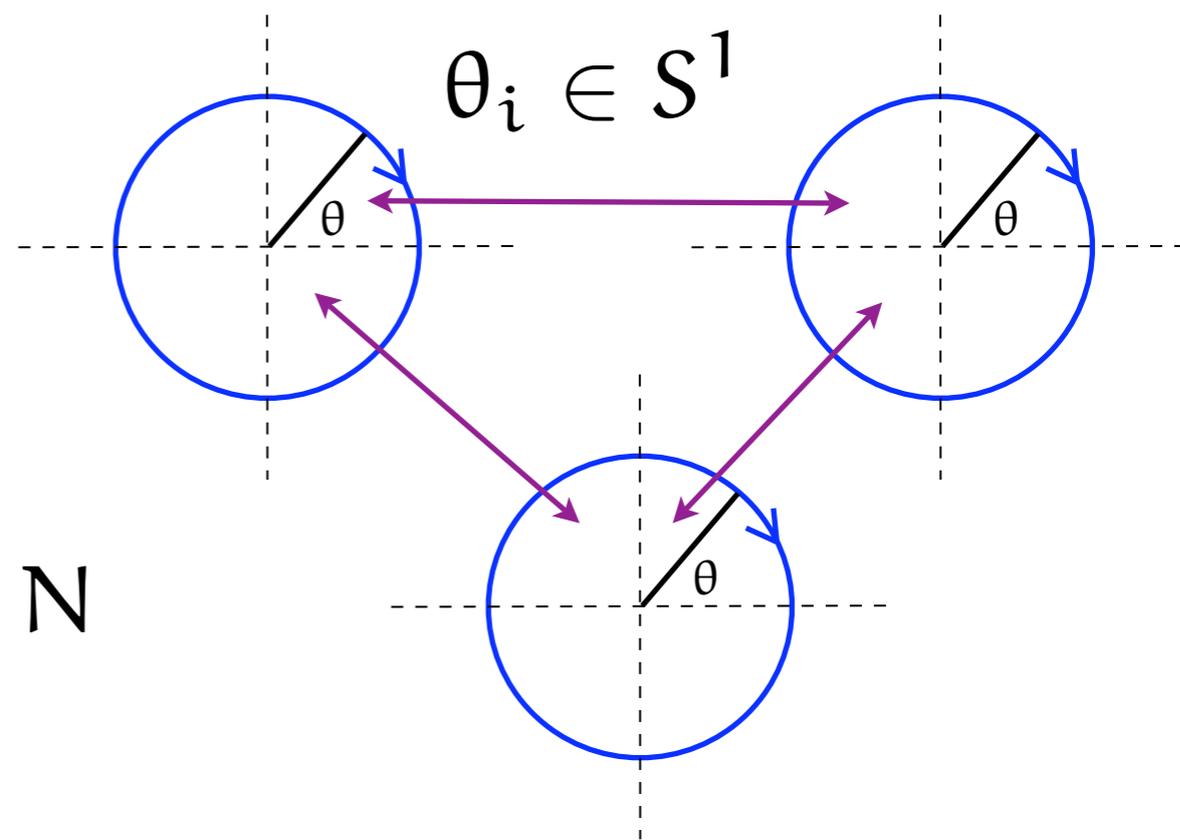
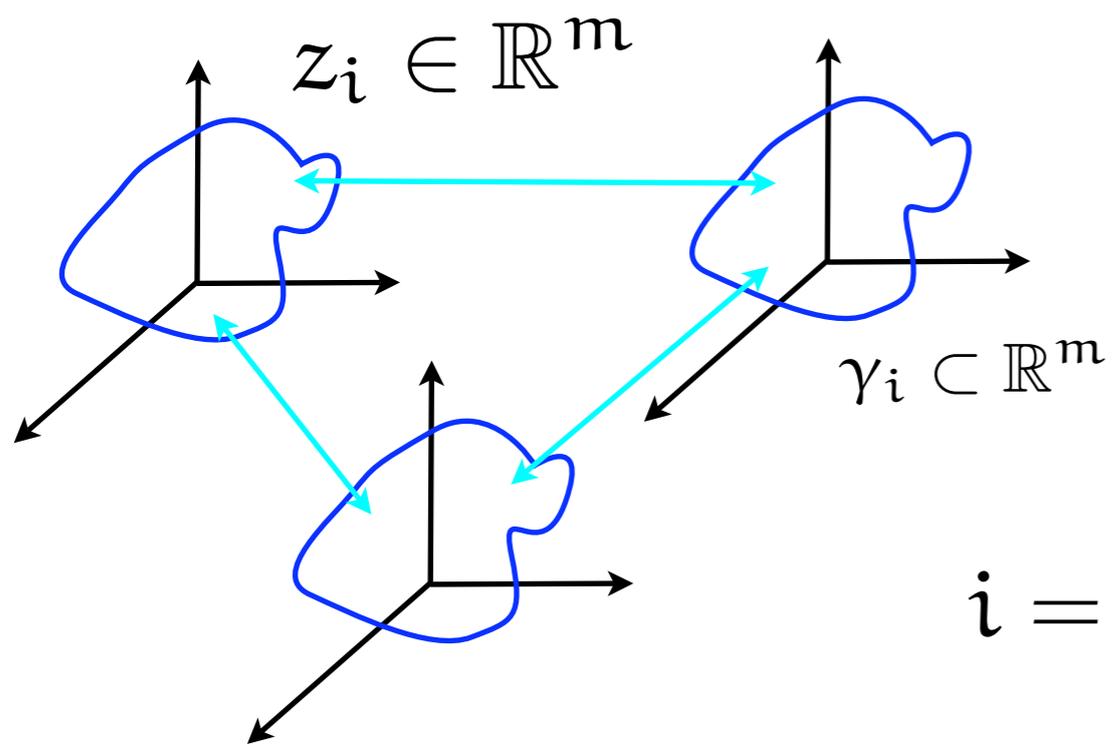
$$\frac{dQ}{dt} = D(t)Q, \quad D(t) = -DF^T(Z(t))$$

$$\nabla_{Z(0)} \cdot F(Z(0)) = \frac{1}{T} \text{ and } Q(t) = Q(t + T)$$

$$\dot{\theta} = \frac{1}{T}$$



Weak Coupling



$i = 1, \dots, N$

$$\dot{z}_i = F(z_i) + \epsilon G_i(z_1, \dots, z_N)$$

Uncoupled system has an exponentially stable limit cycle γ_i

Direct product of hyperbolic limit cycles is a normally hyperbolic invariant manifold

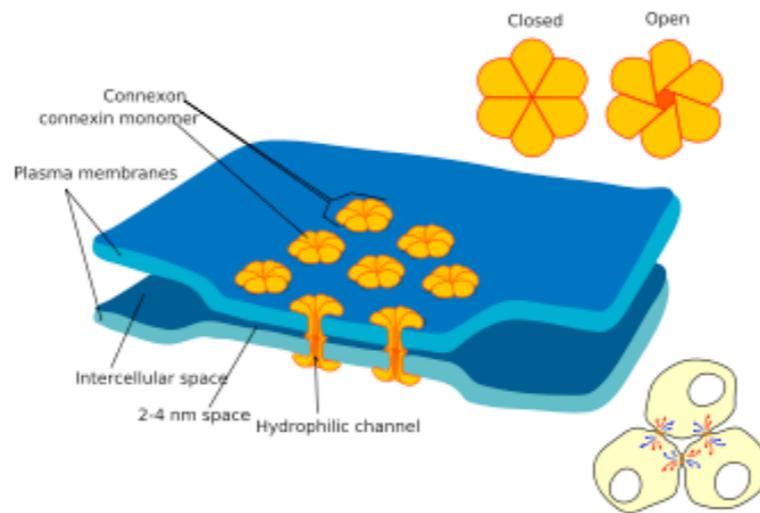
$$\dot{\theta}_i = \frac{1}{T} + \epsilon \langle Q(\theta_i), G_i(\Gamma(\theta)) \rangle$$

Drive

PRC

Coupled oscillator networks

An example:
gap junction
coupling



$$\frac{1}{N} \sum_{j=1}^N (v_j - v_i)$$

Averaging gives

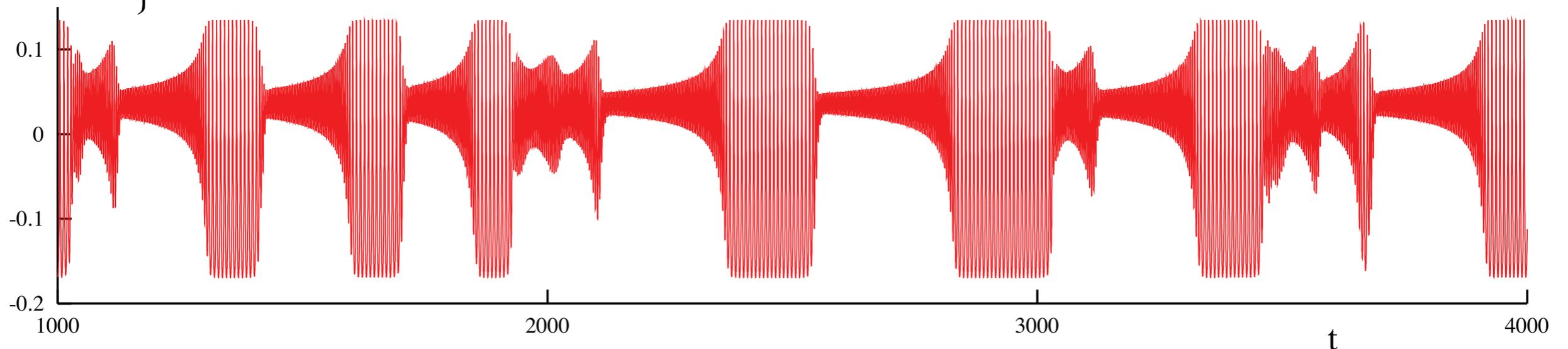
$$H(\theta) = \frac{1}{T} \int_0^T \langle Q(t), (v(t + \theta T) - v(t), 0) \rangle dt$$

Kopell and Ermentrout

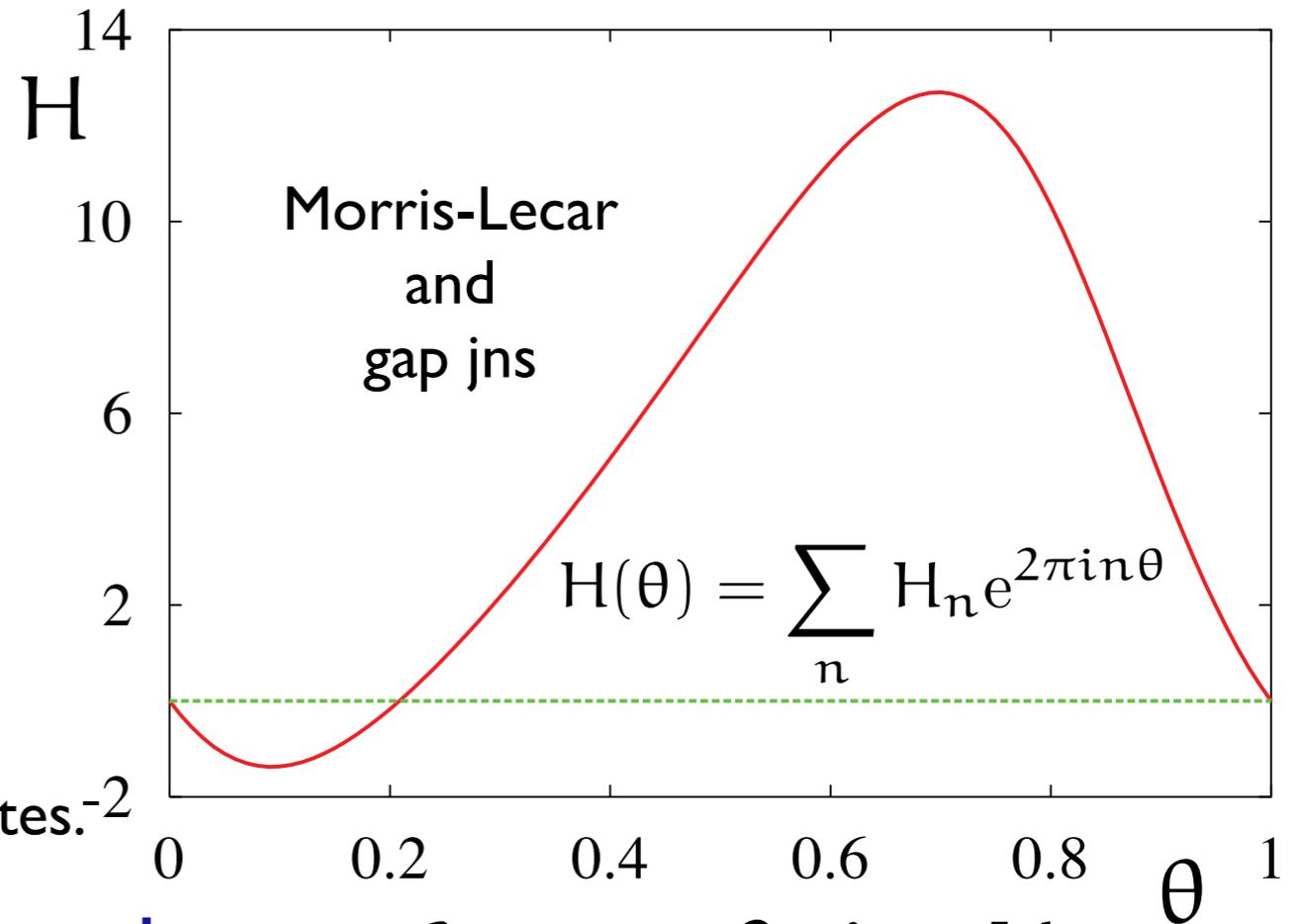
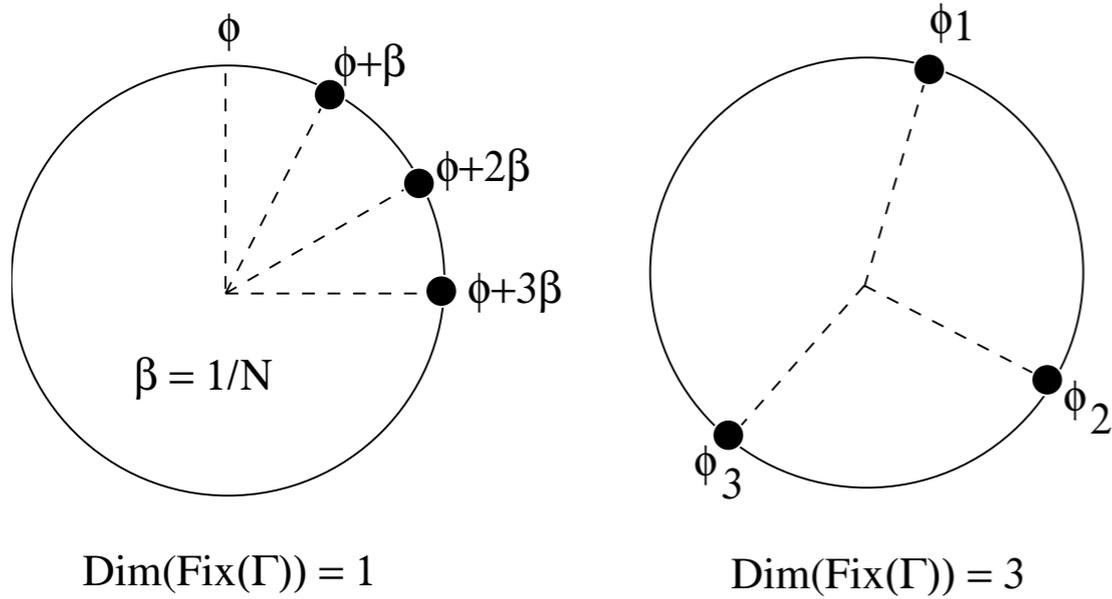
$$\dot{\theta}_i = \frac{1}{T} + \frac{\epsilon}{N} \sum_j H(\theta_j - \theta_i)$$

$$E = \frac{1}{N} \sum_j v_j$$

Morris-Lecar



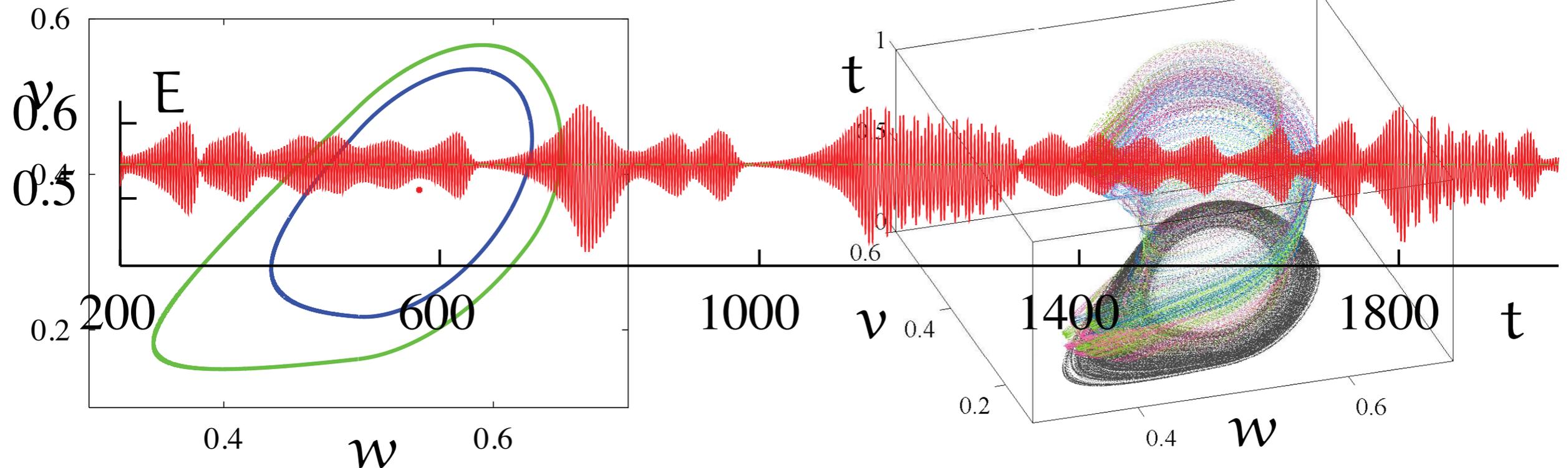
Stability of phase-locked states



Bifurcations from maximally symmetric solutions to ones with smaller isotropy groups. eg. cluster states.

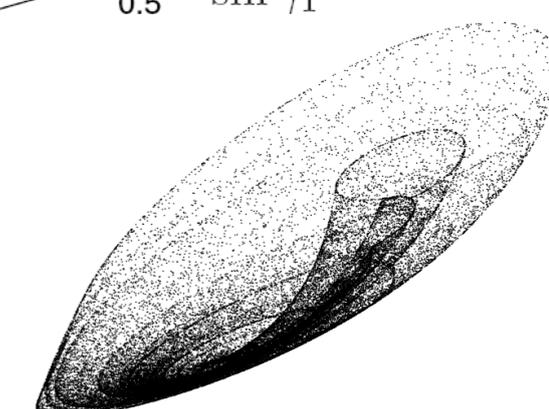
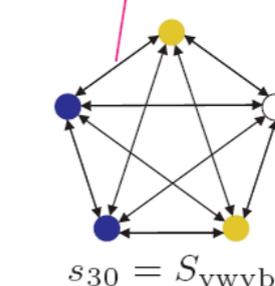
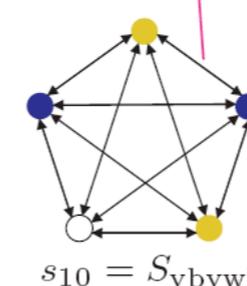
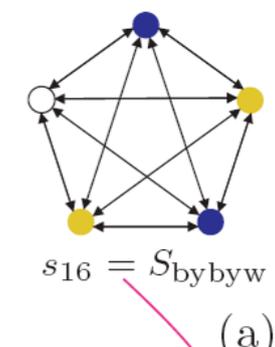
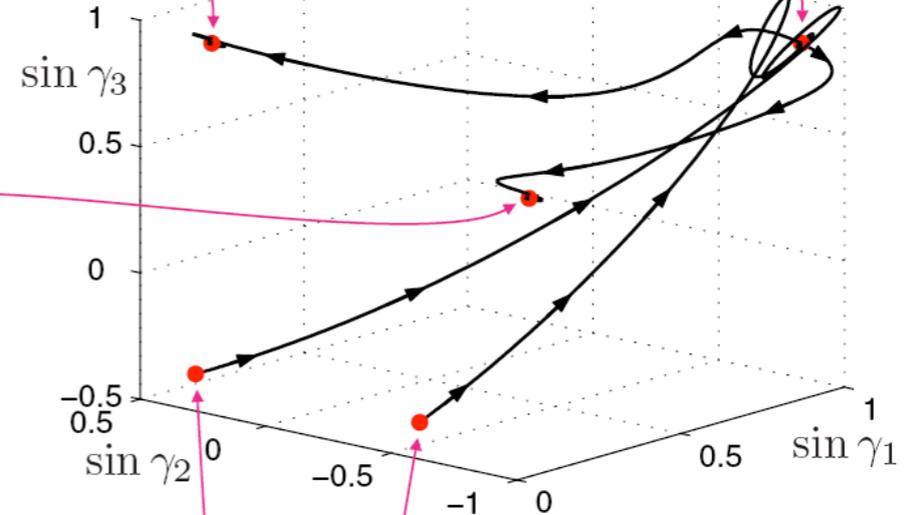
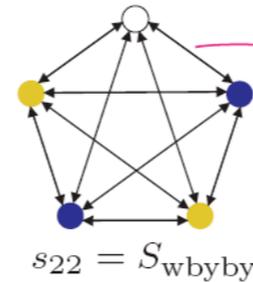
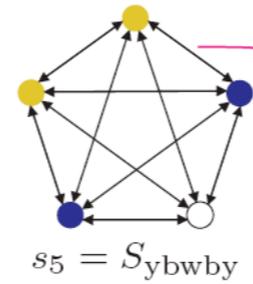
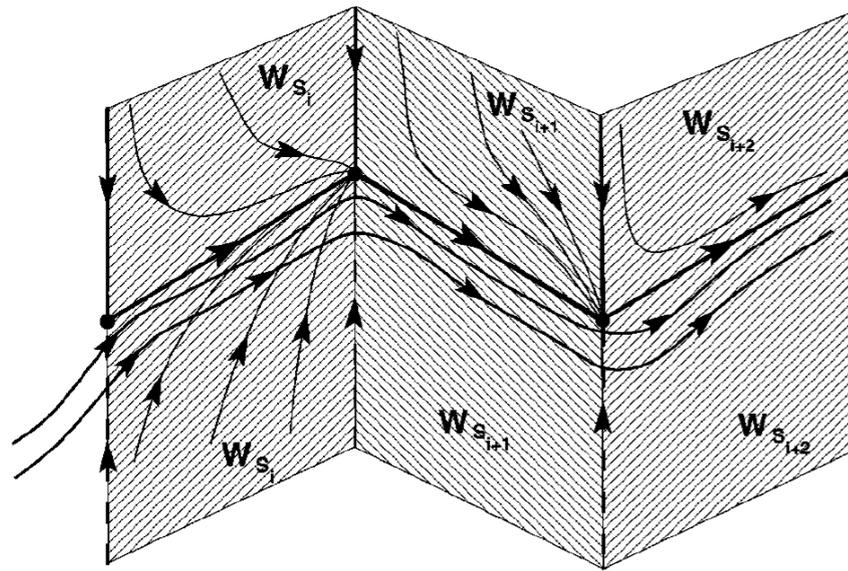
Synchrony $\lambda = -\epsilon H'(0)$

Asynchrony $\lambda_n = -2\pi i n \epsilon H_{-n}$



Heteroclinic cycles

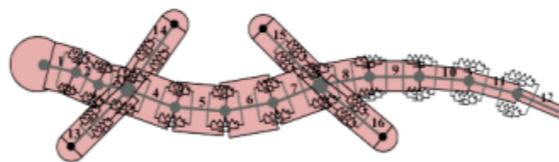
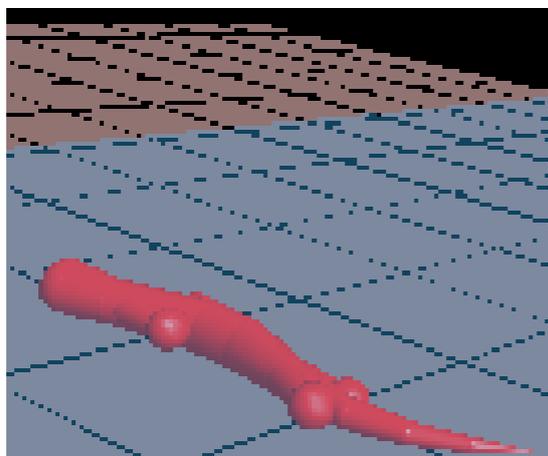
Winnerless networks



Rabinovich *et al.* Dynamical principles in neuroscience, Rev. Mod. Phys., 78, 2006.

Ashwin *et al.* SIADS, Dynamics on networks of cluster states for globally coupled phase oscillators, 6, 2007.

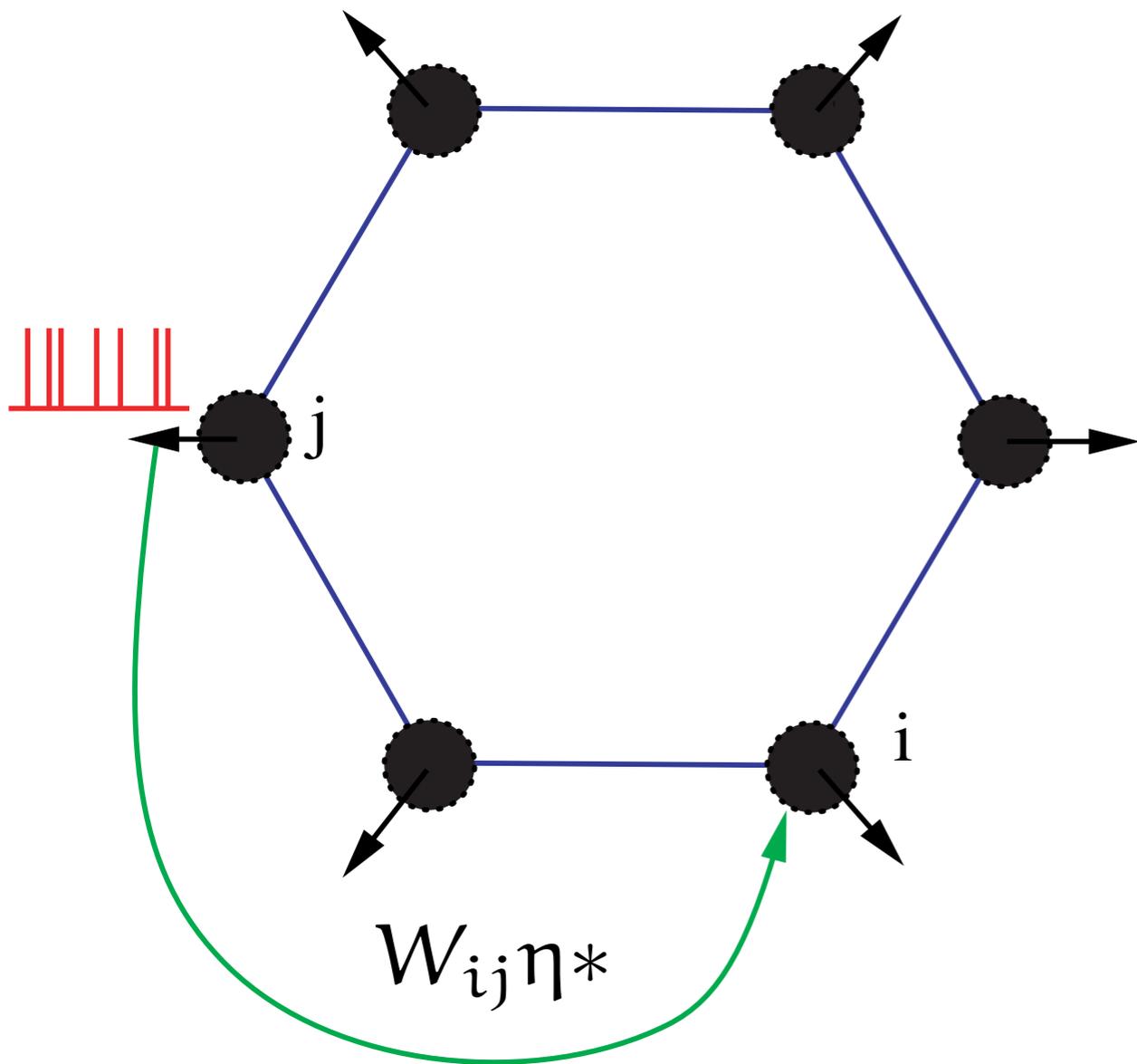
Applications of weakly coupled oscillator theory to CPGs, robot control, ...



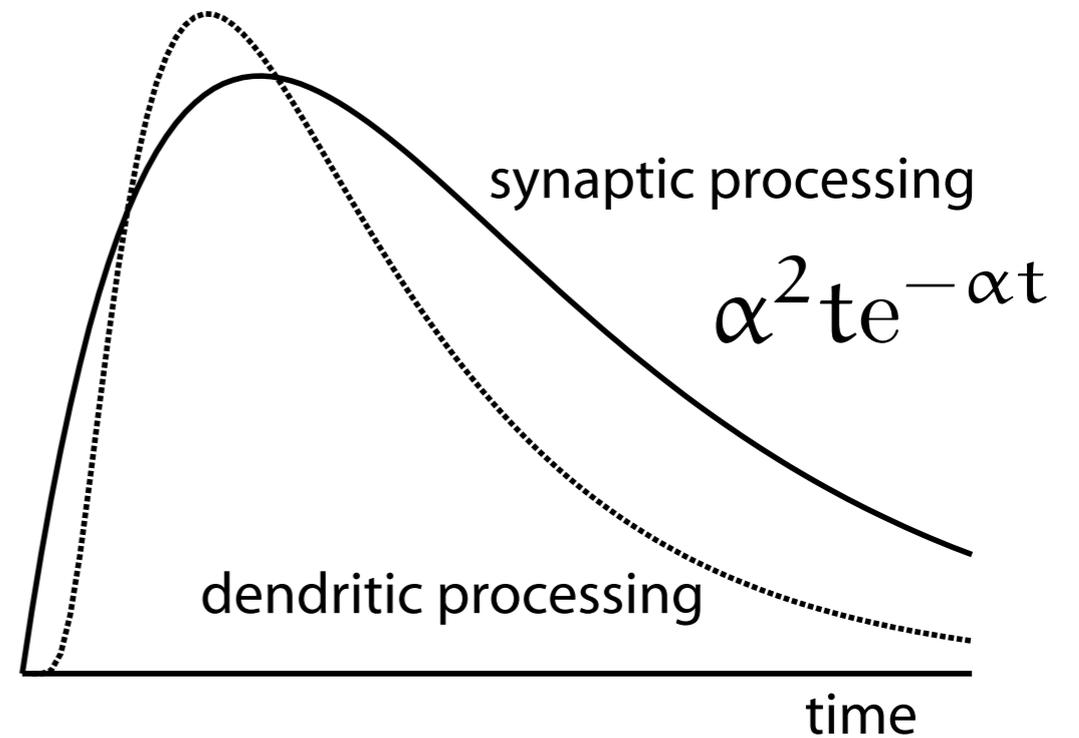
Biorobotics lab at EPFL
<http://biorob.epfl.ch/>



Strongly coupled synaptic networks



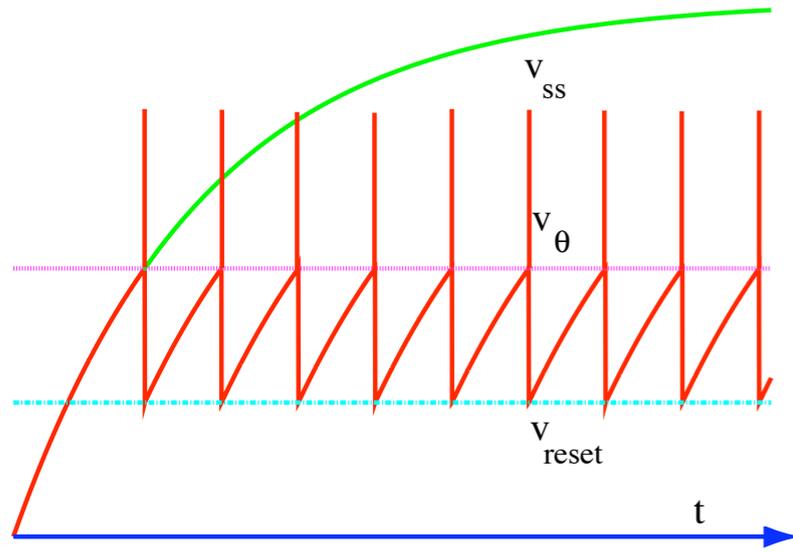
PSP



$$s_i(t) = g_s (v_s - v_i(t)) \sum_{j=1}^N W_{ij} \sum_{m \in \mathbb{Z}} \eta(t - T_j^m)$$

$$T_i^m = \inf\{t \mid v_i(t) > h, \dot{v}_i > 0, t > T_i^{m-1}\}$$

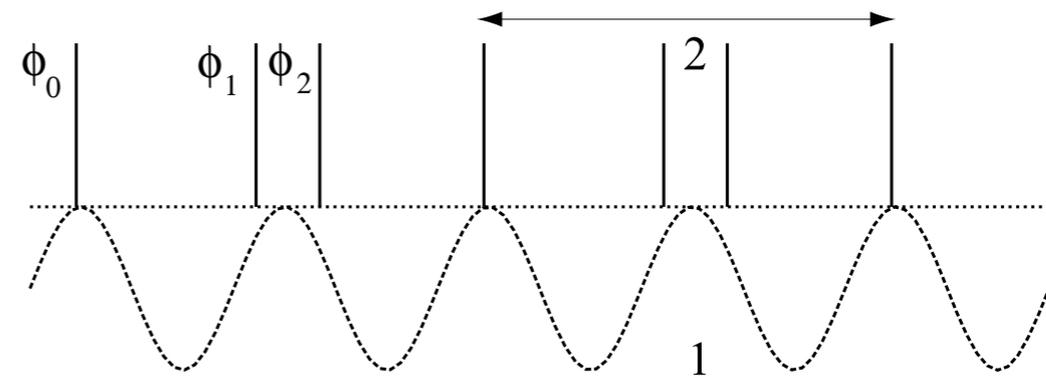
Integrate-and-fire neurons



$$\frac{dv}{dt} = -\frac{v}{\tau} + A(t), \quad t \in (T^m, T^{m+1})$$

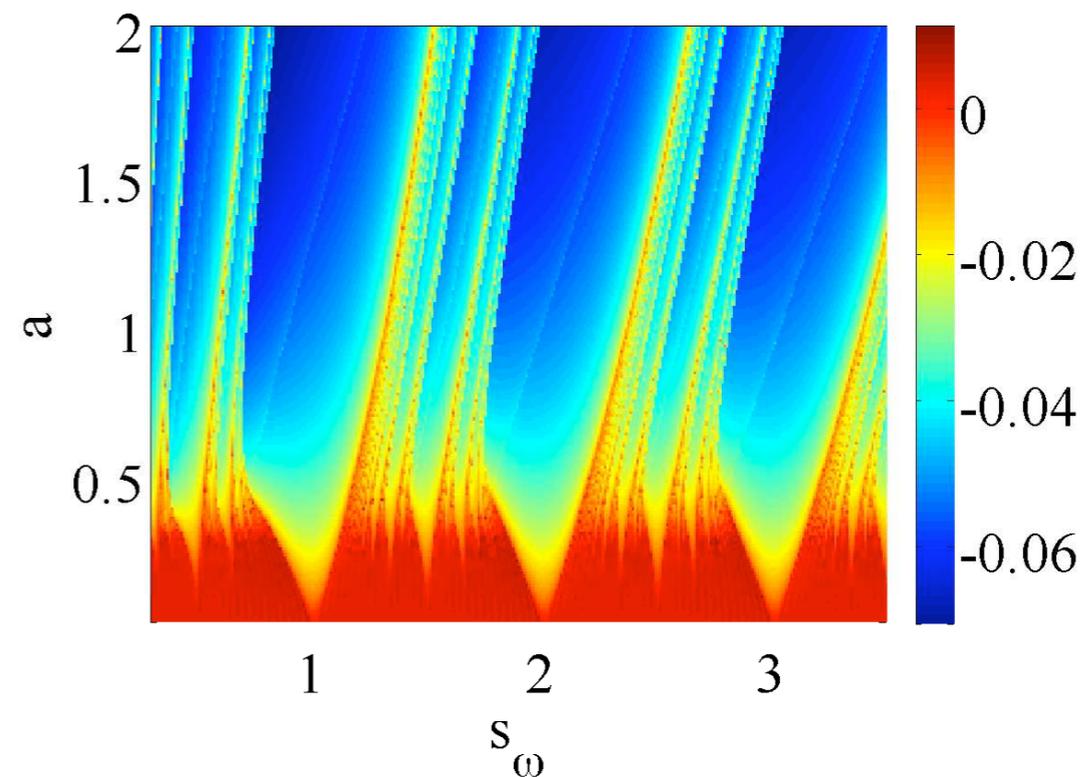
subject to nonlinear reset

Periodic forcing gives $p : q$ mode-locked states



Implicit map of firing times

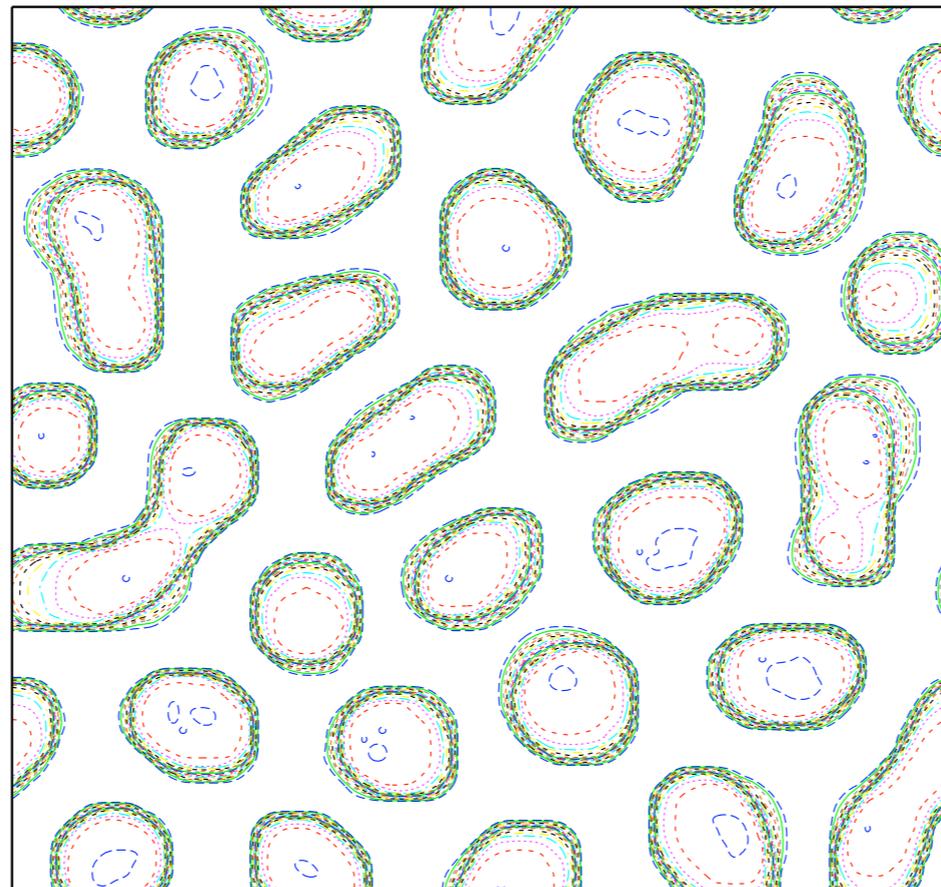
Arnol'd tongue structure dominated by non-smooth bifurcations



CML - discrete time IF

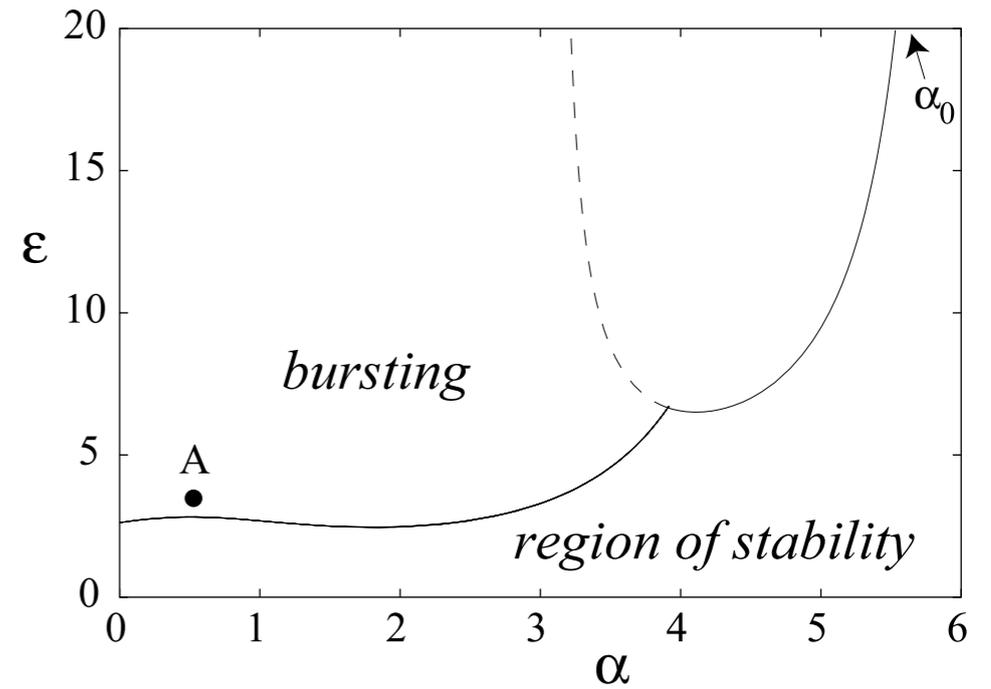
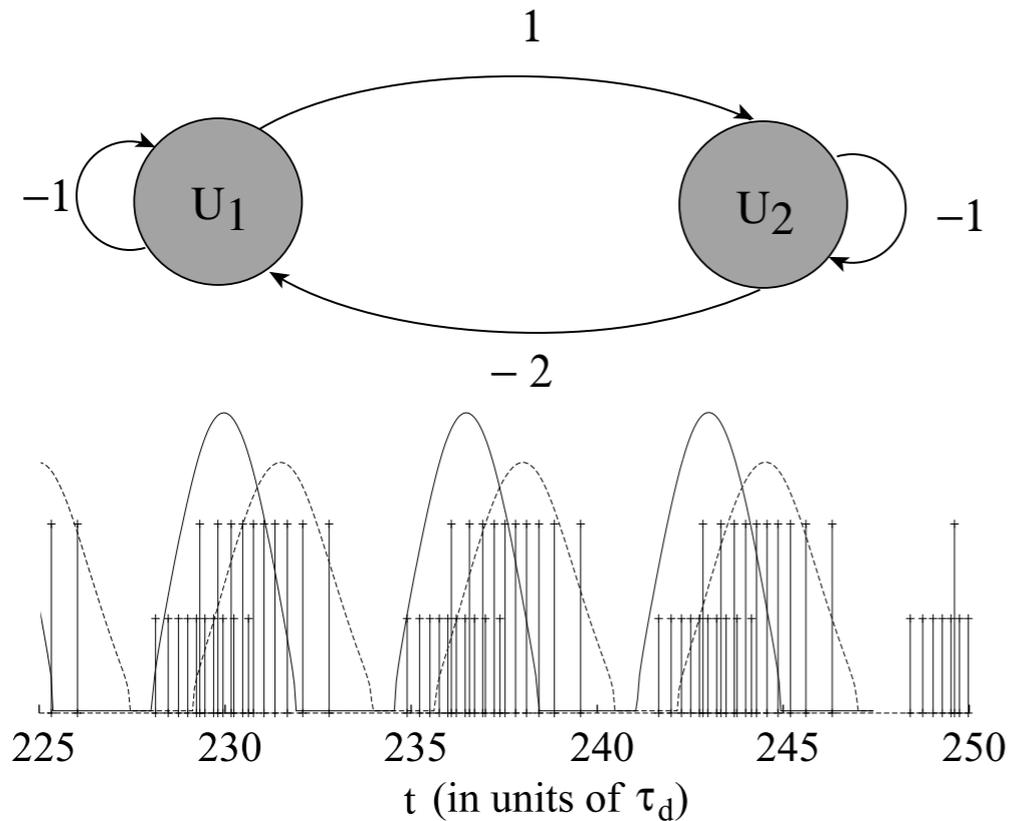
$$V_i(\mathbf{n} + 1) = [\gamma V_i(\mathbf{n}) + \epsilon \sum_j W_{ij} a_j(\mathbf{n})] \Theta(1 - V_i(\mathbf{n}))$$

$$a_i(\mathbf{n}) = \Theta(V_i(\mathbf{n}) - 1)$$



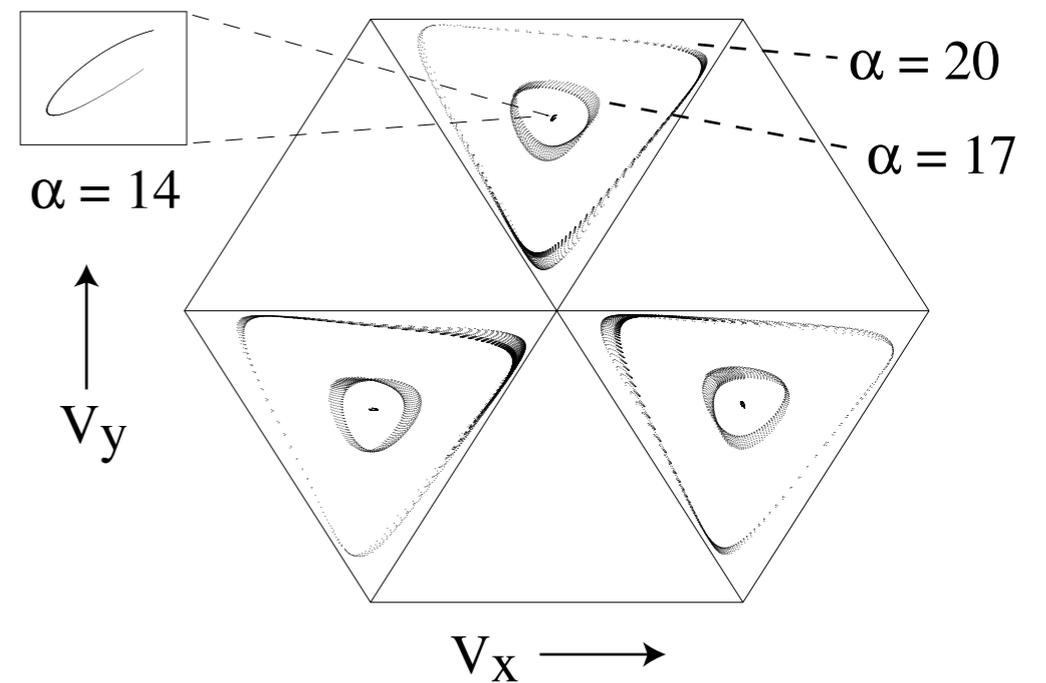
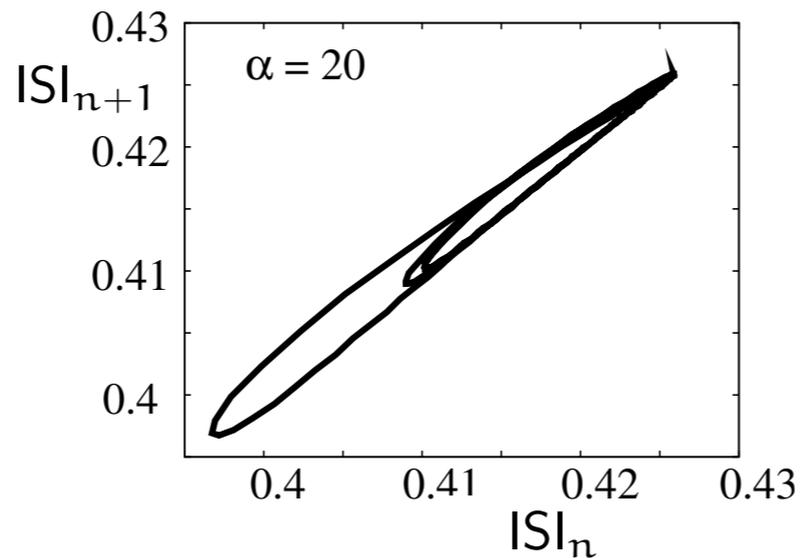
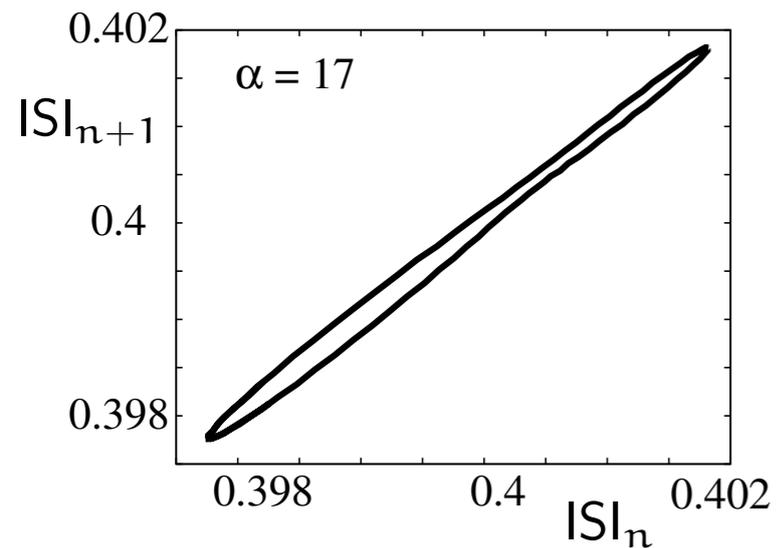
Mexican hat interaction

Network firing maps

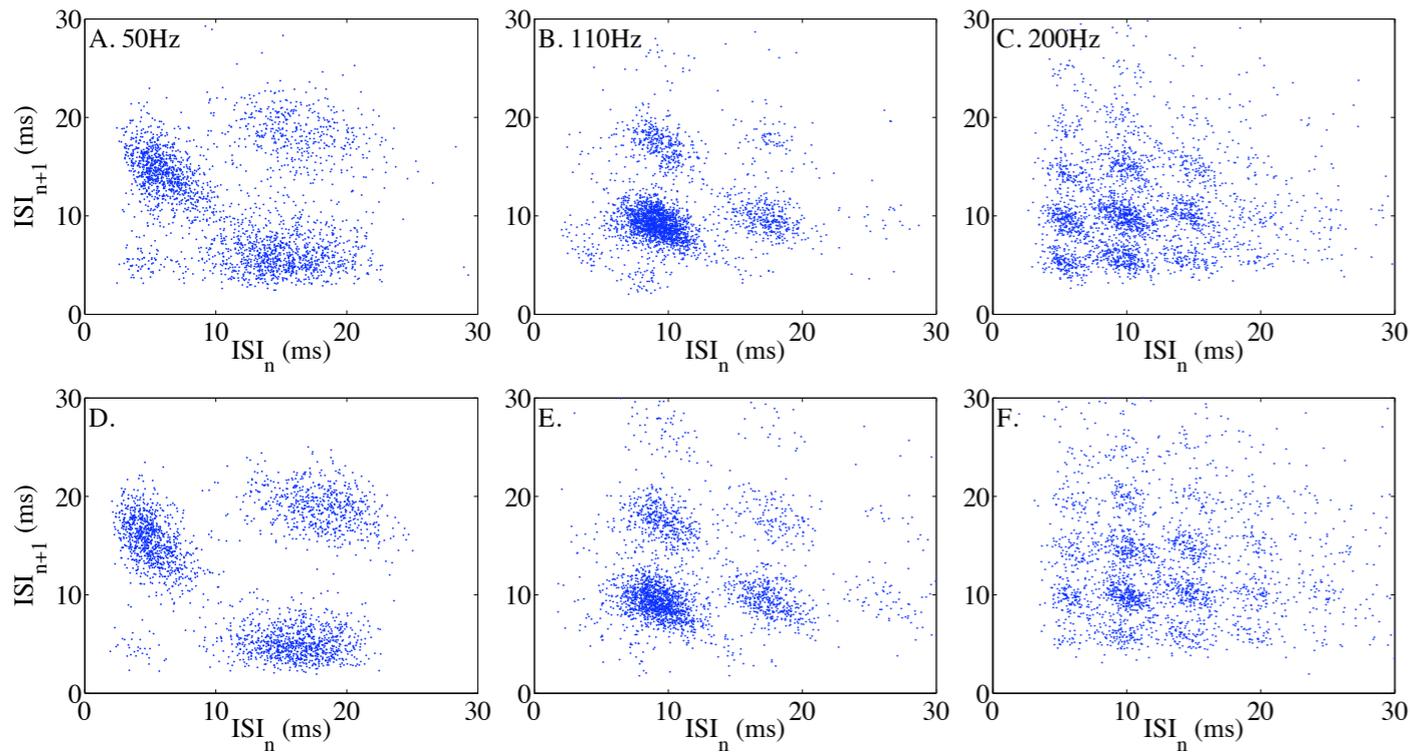


Global heteroclinic bifurcation ($N=3$) $ISI_n = T^{n+1} - T^n$

$$V_x(t) + iV_y(t) = \sum_{m=1}^3 e^{2\pi i m/3} U_m(t)$$



Fits to data

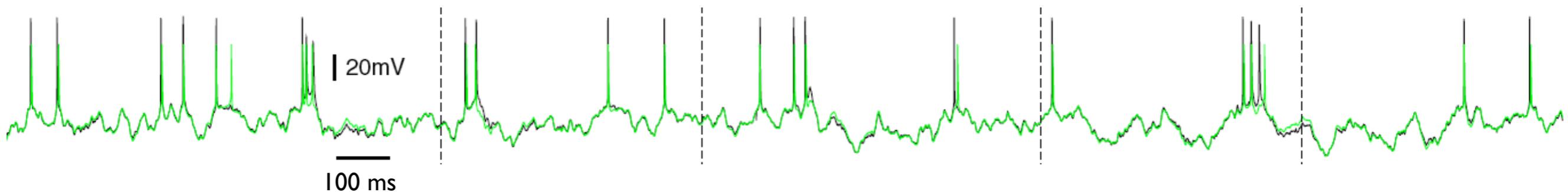


VCN stellate cell

Linear IF and threshold noise

$$ISI_n = T^{n+1} - T^n$$

J Laudanski *et al.*, Journal of Neurophysiology, 103, 2010



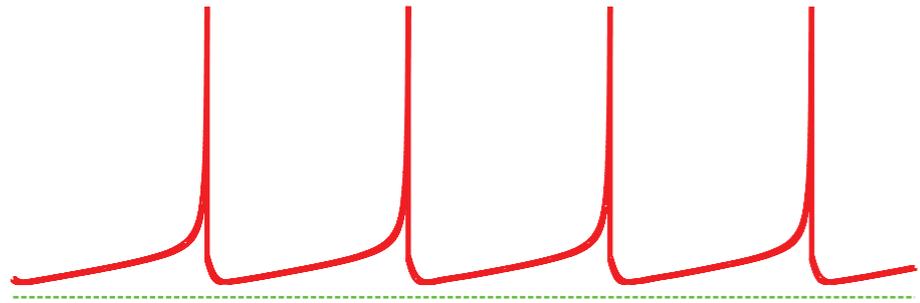
Nonlinear IF

$$-\frac{1}{\tau}(v - v_L) + \frac{\kappa}{\tau} e^{(v - v_\kappa)/\kappa}$$

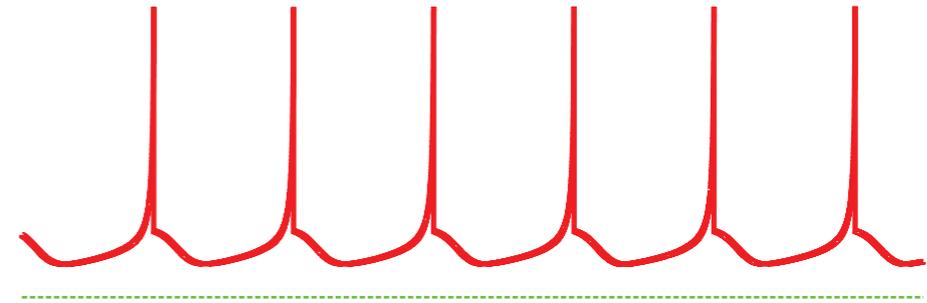
Layer V cortical pyramidal cell

Badel *et al.*, Journal of Neurophysiology, 99, 2010

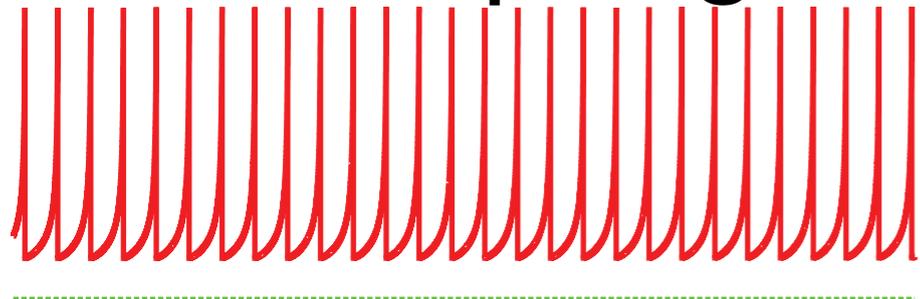
Regular spiking



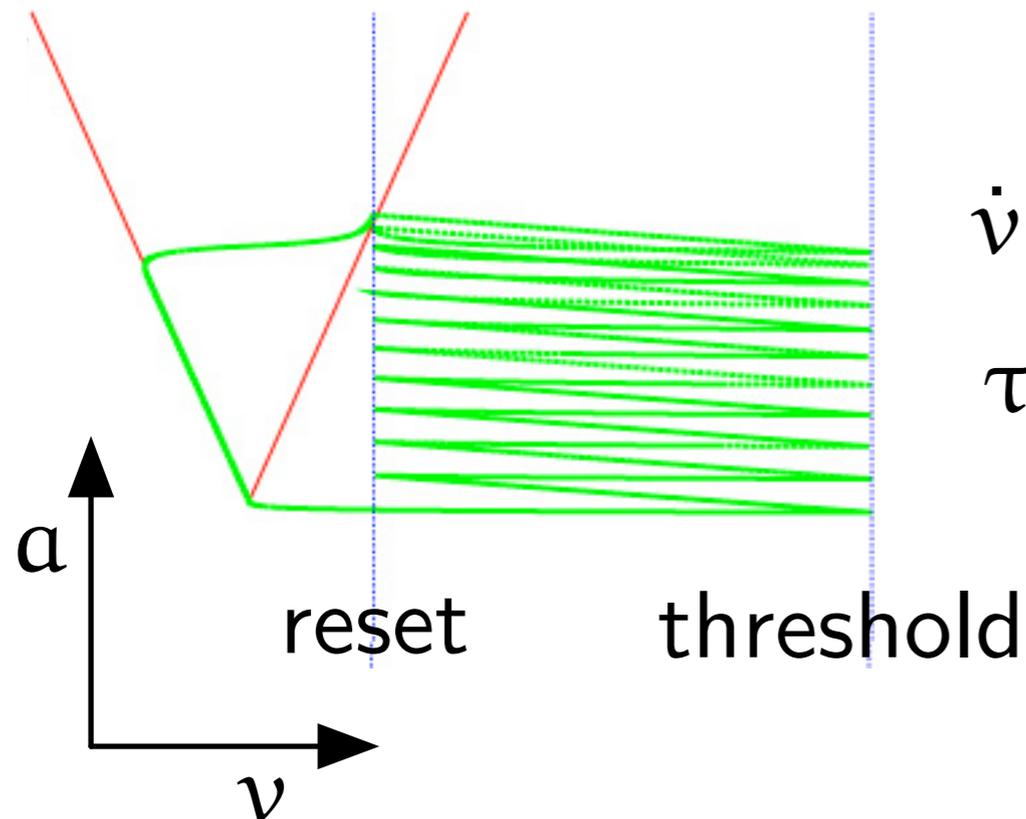
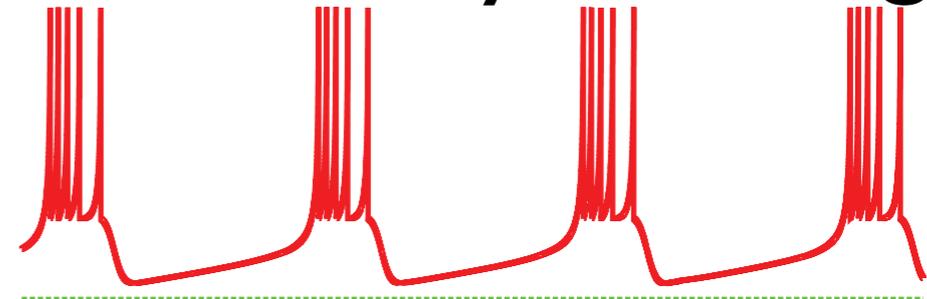
Chattering



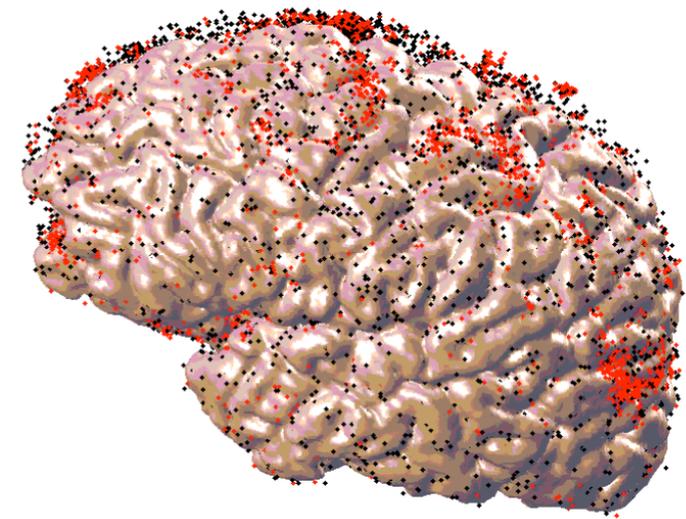
Fast spiking



Intrinsically bursting

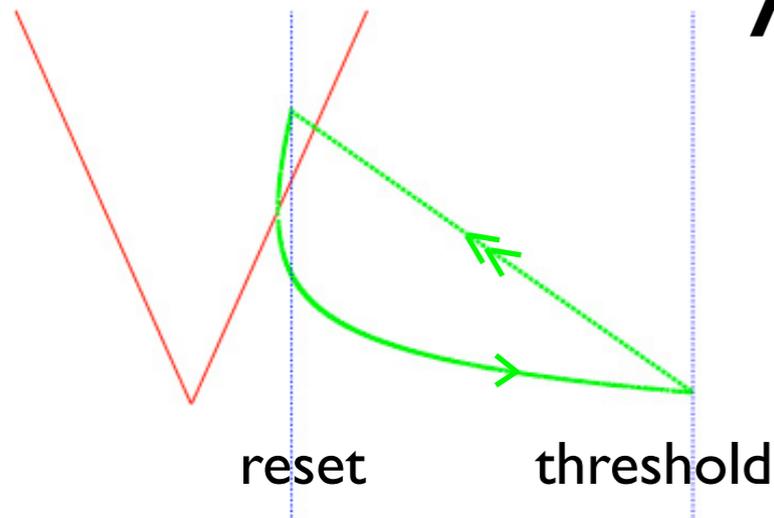


$$\dot{v} = |v| +$$
$$\tau \dot{a} = -a$$



S Coombes and M Zachariou 2009, in
Coherent Behavior in Neuronal Networks
(Ed. Rubin, Josic, Matias, Romo), Springer.

Absolute IF networks



$$a(T^m) \rightarrow a(T^m) + g_a/\tau_a$$

Orbit and PRC in closed form (pwl system)

Gap j n network: asynchronous state

network averages

time averages

$$\lim_{N \rightarrow \infty} \frac{1}{N} \sum_{j=1}^N v(t + jT/N) = \frac{1}{T} \int_0^T v(t) dt \equiv v_0$$

$$\dot{v} = |v| - \epsilon v + I - a + \epsilon v_0, \quad \dot{a} = -a/\tau_a$$

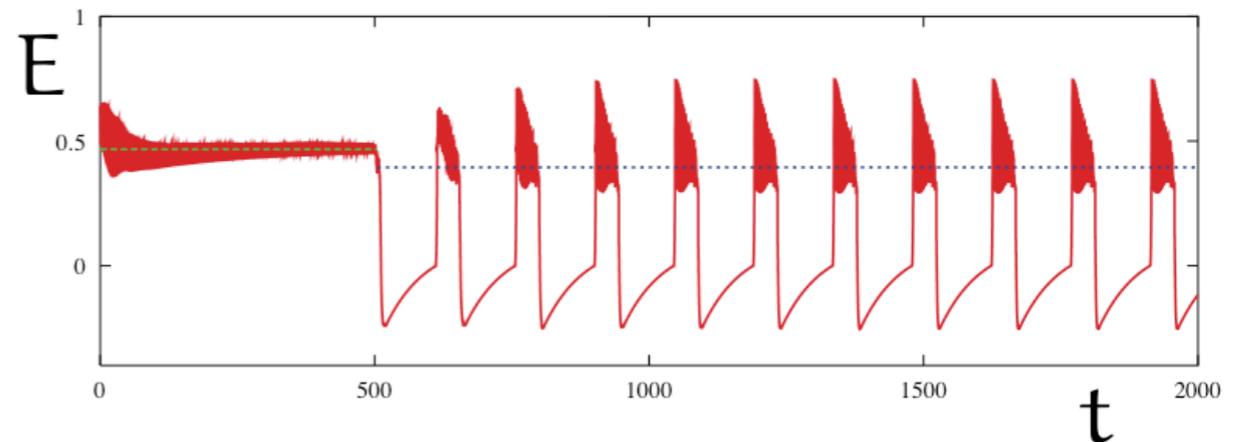
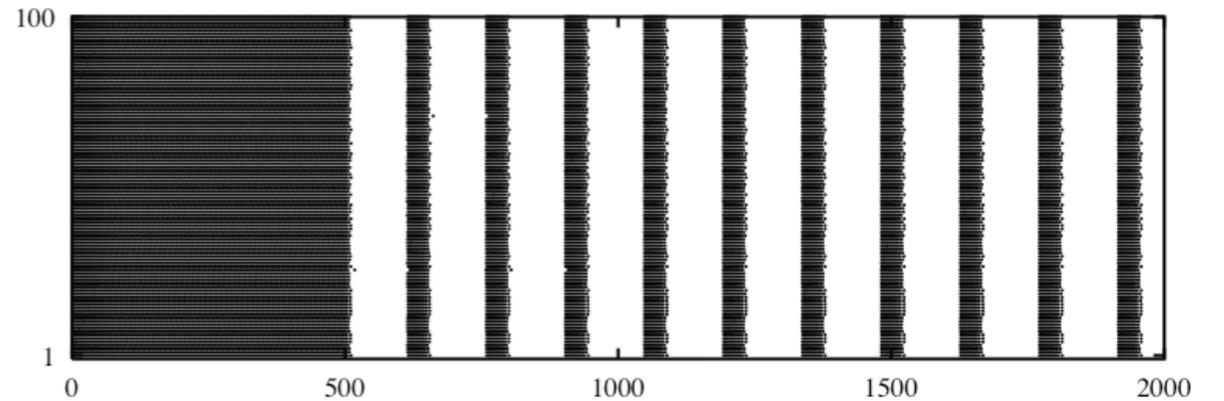
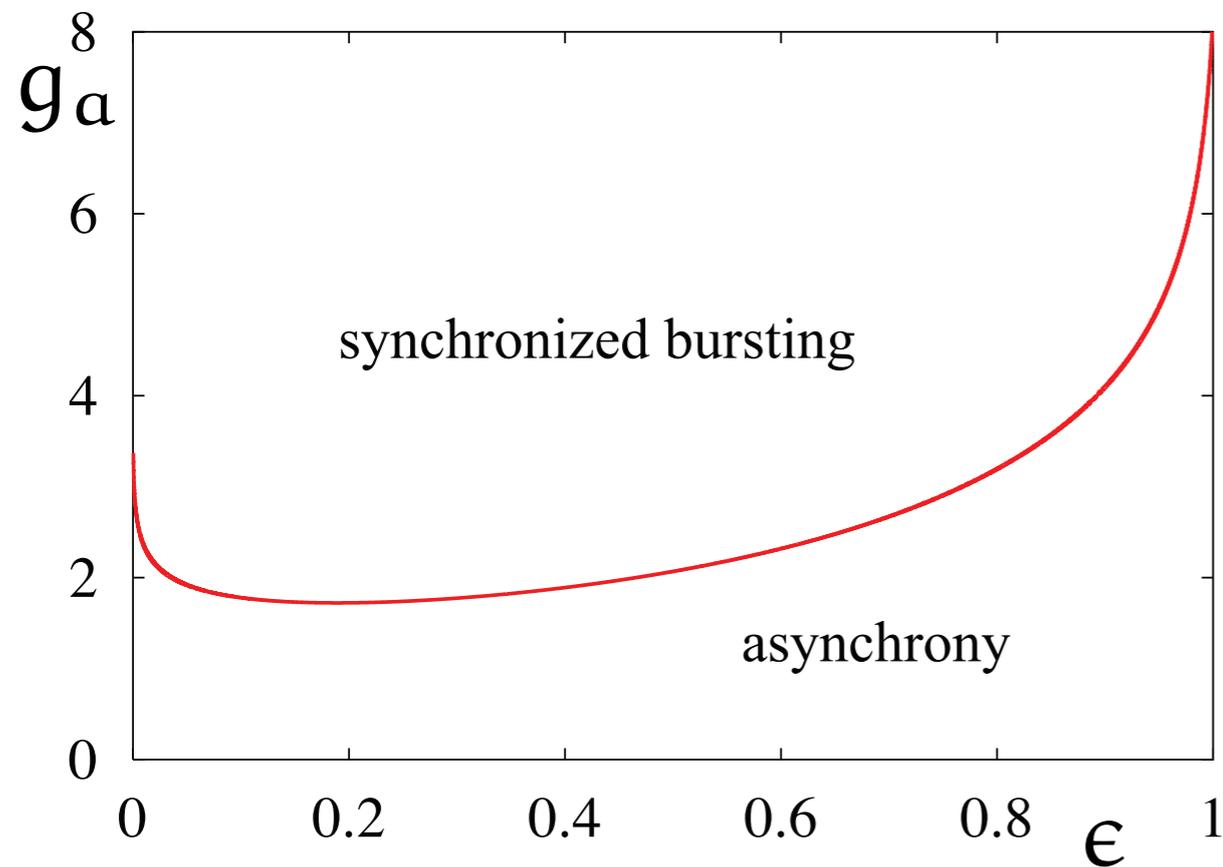
advanced-retarded ode - self-consistent periodic solution

Stability and bifurcations

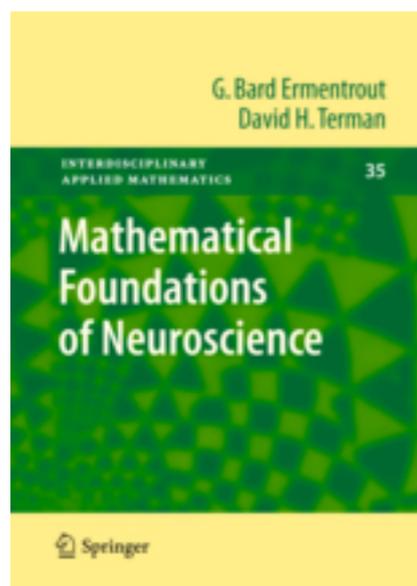
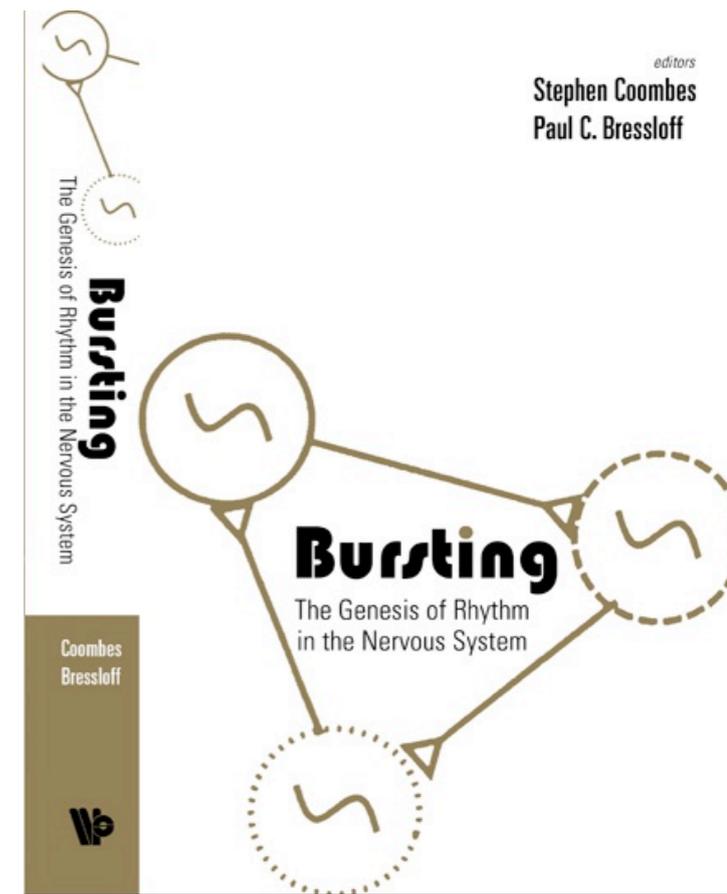
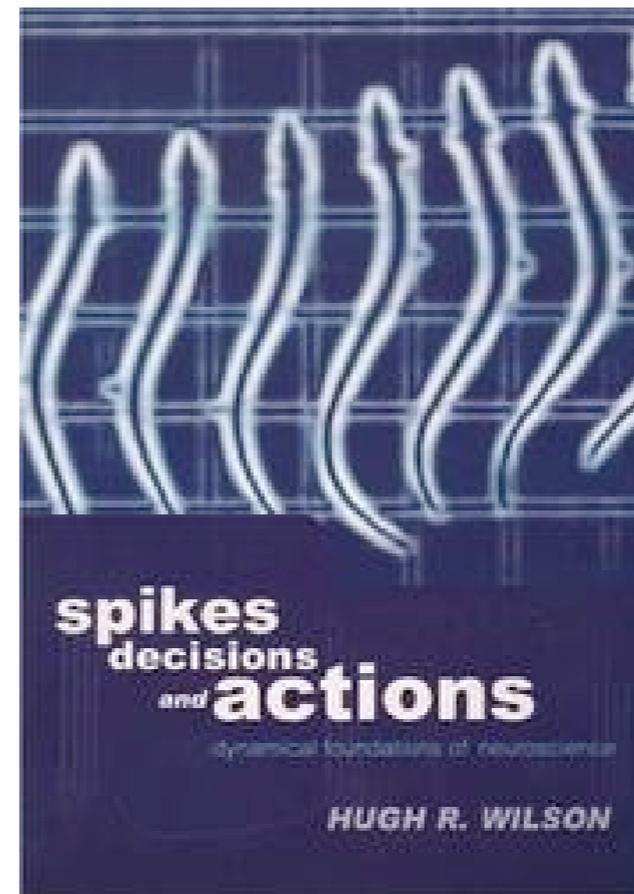
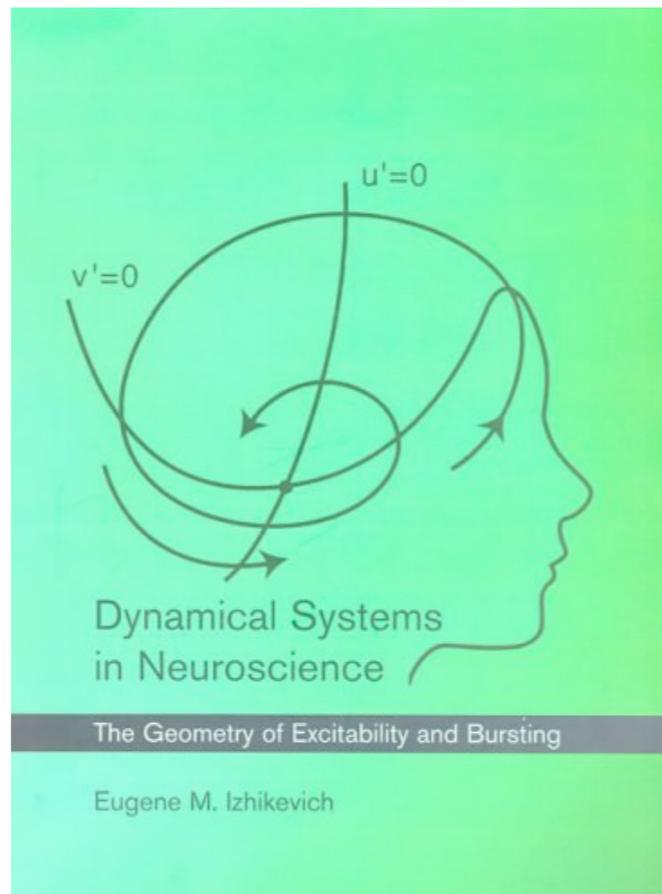
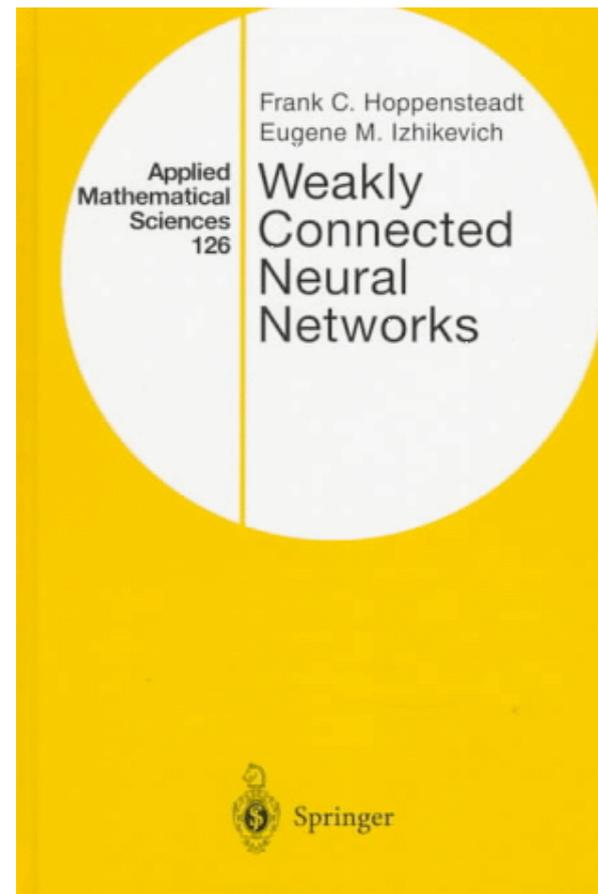
e-values as zeros of (using phase density formalism):

$$\text{LT of orbit} \longrightarrow \mathcal{E}(\lambda) = \frac{e^{\lambda T}}{\tilde{\nu}(\lambda)} + \epsilon \lambda T \int_0^1 R(\theta) e^{\lambda \theta T} d\theta$$

PRC of splay



Books



Physica D: Nonlinear Phenomena
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 pp. 727-904 (1 June 2010)
 Emergent Phenomena in Spatially Distributed Systems - In Honor of Stefan C. Müller, Emergent Phenomena in Spatially Distributed Systems

Volume 239, Issue 10

Volume 239, Issue 9, Pages 475-578 (1 May 2010)

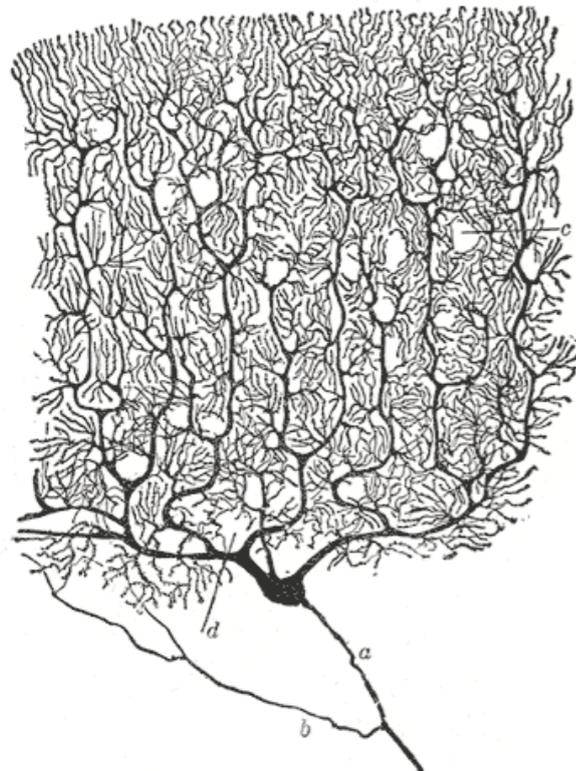
Mathematical Neuroscience
 Edited by Stephen Coombes and Yulia Timofeeva

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1. **Special issue: Mathematical neuroscience**
 Pages 475-476
 S. Coombes, Y. Timofeeva
[Preview](#) | [PDF \(410 K\)](#) | [Related Articles](#)

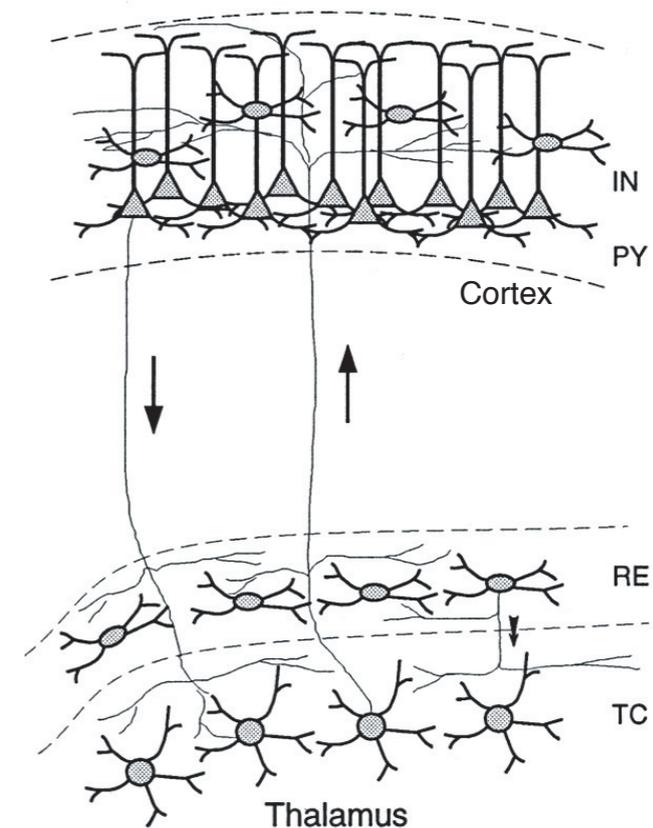
Physica D Special Issue
 Mathematical Neuroscience
 Vol 239, May 2010

Mathematical Neuroscience: from neurons to networks



Part II

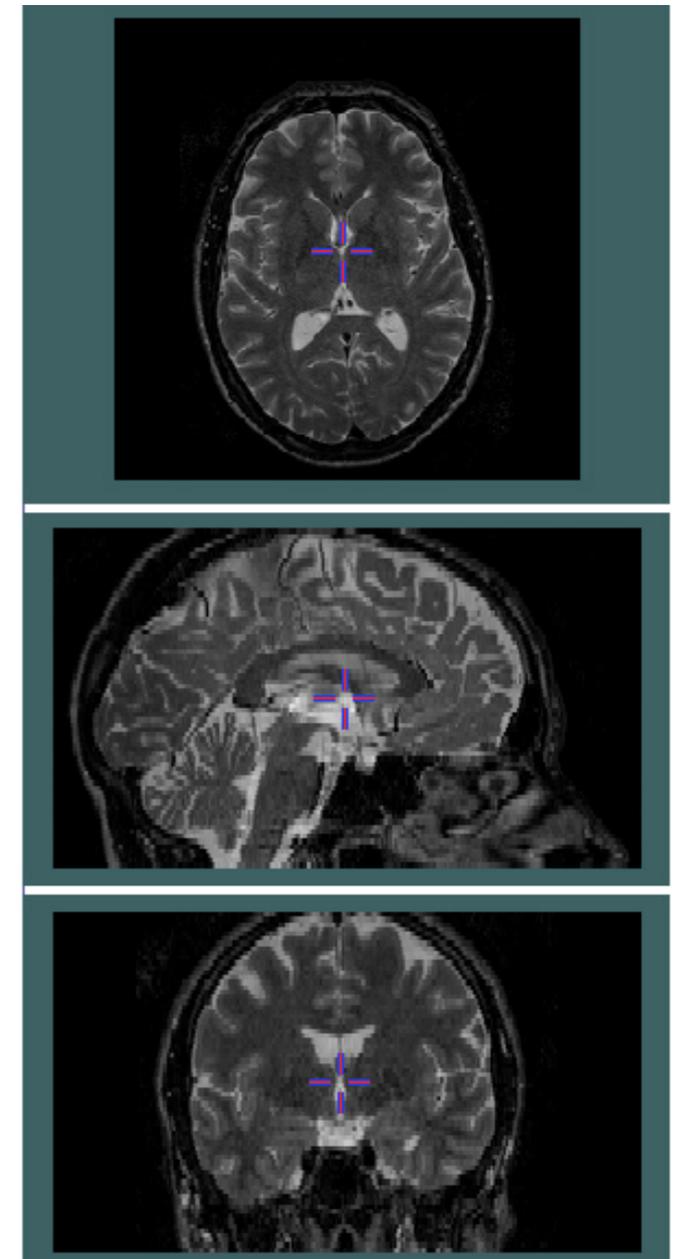
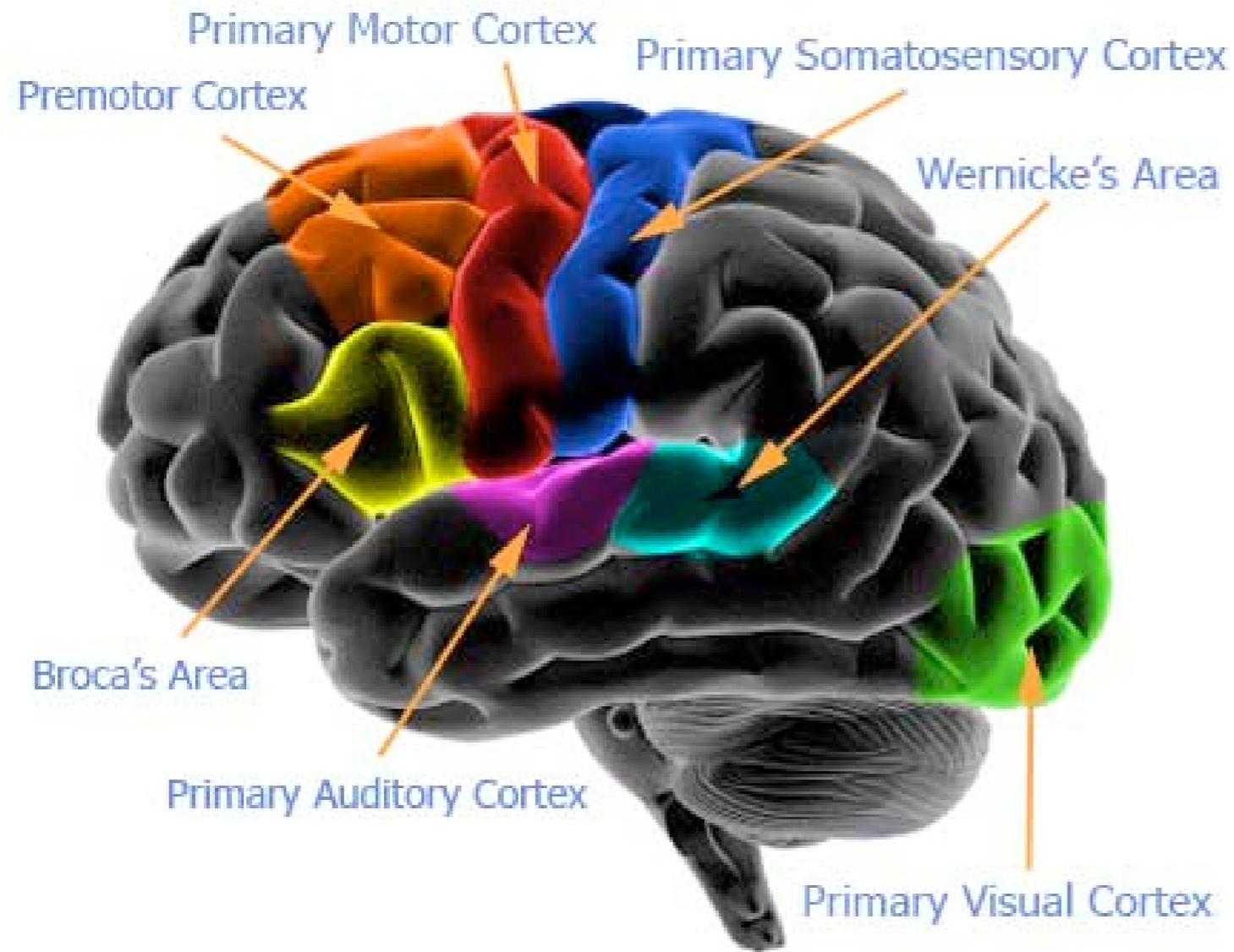
Steve
Coombes



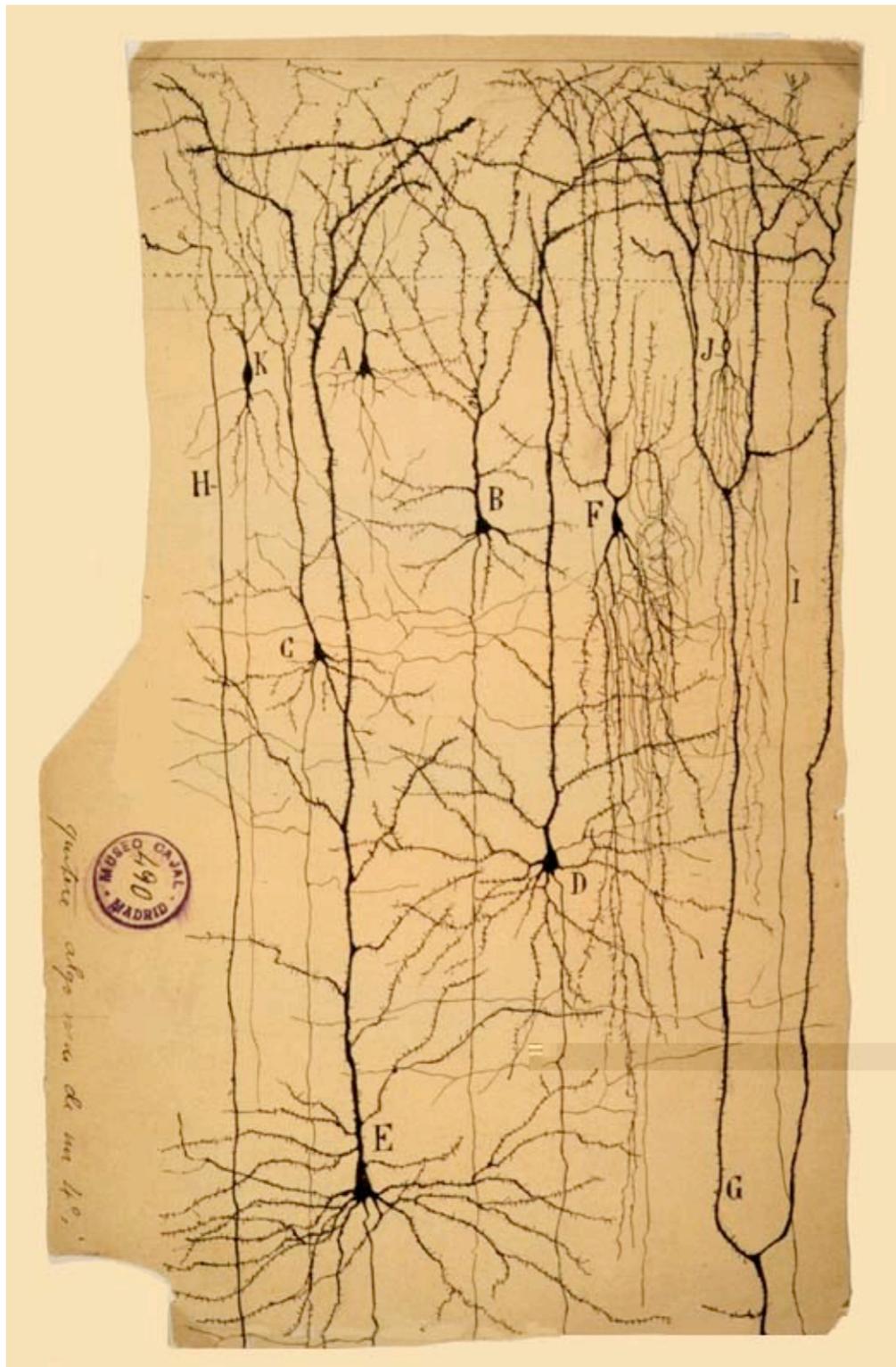
The University of
Nottingham

School of Mathematical
Sciences

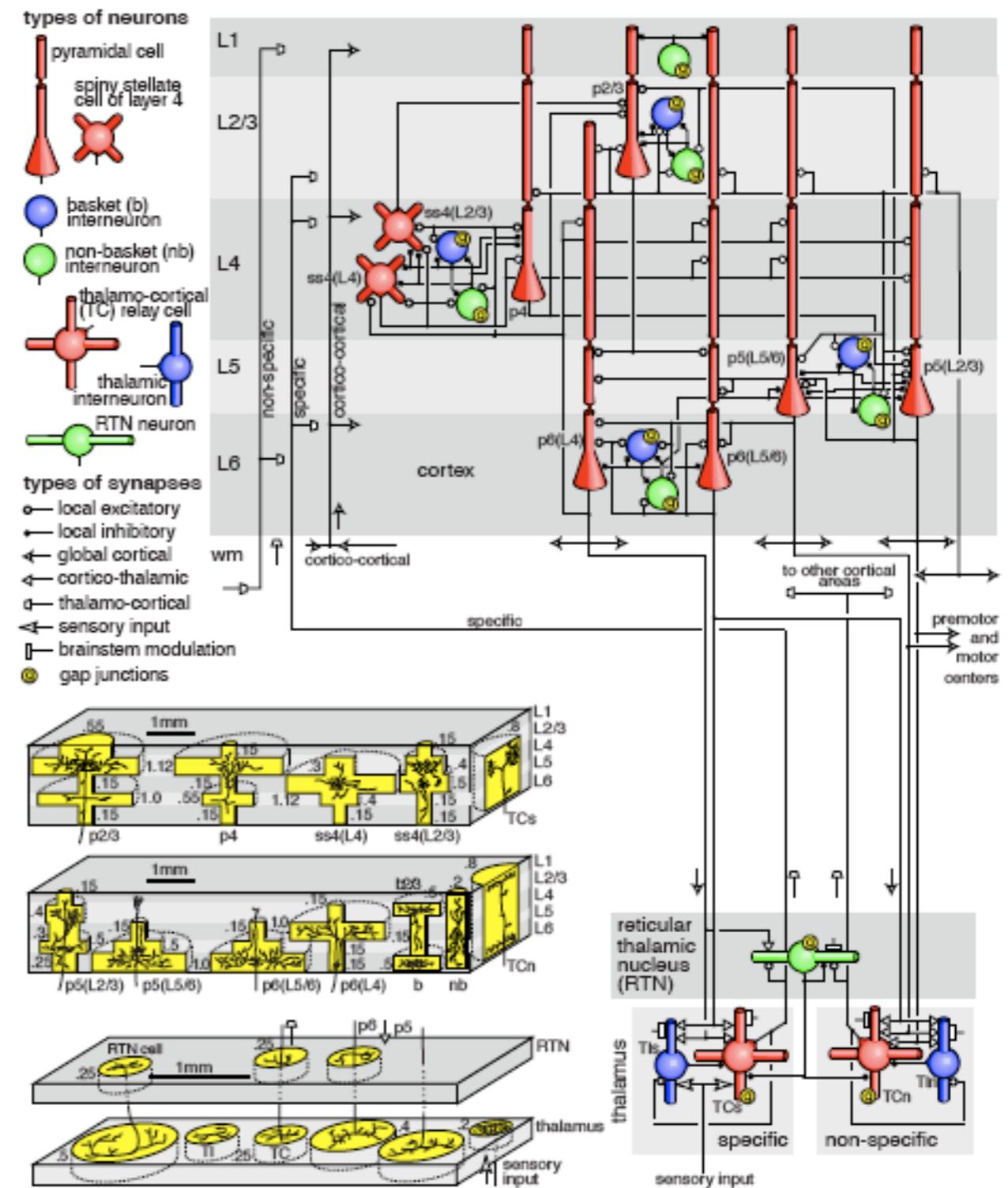
Brain and Cortex



Principal cells and interneurons

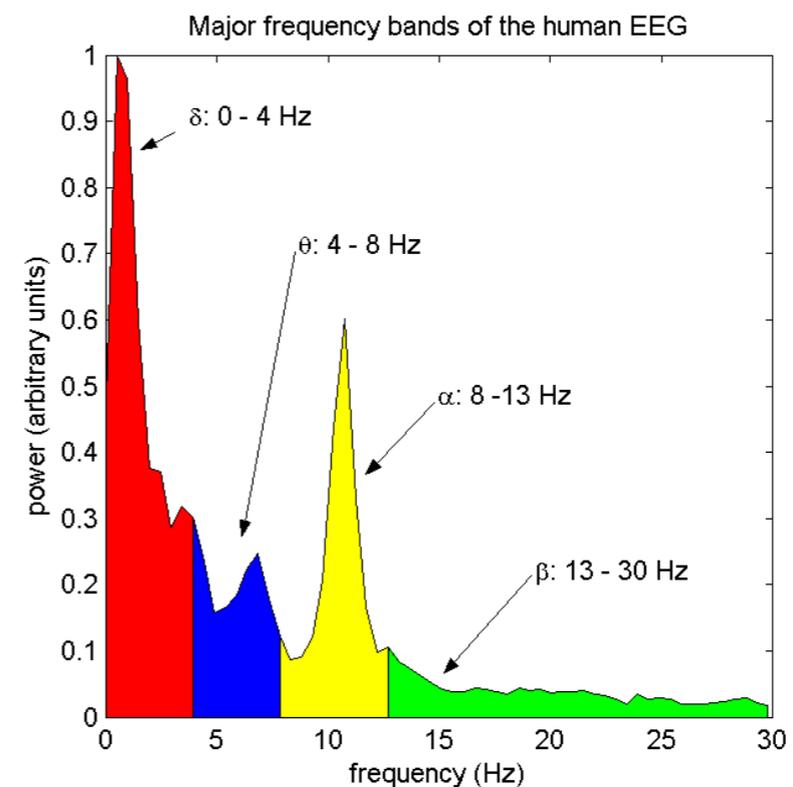
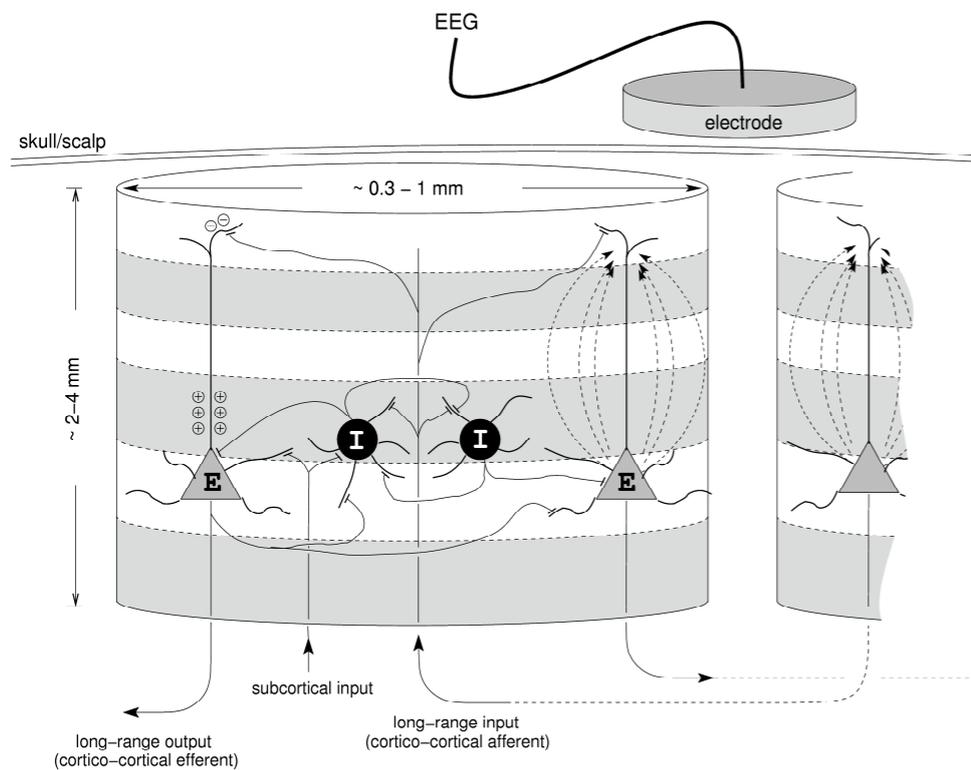
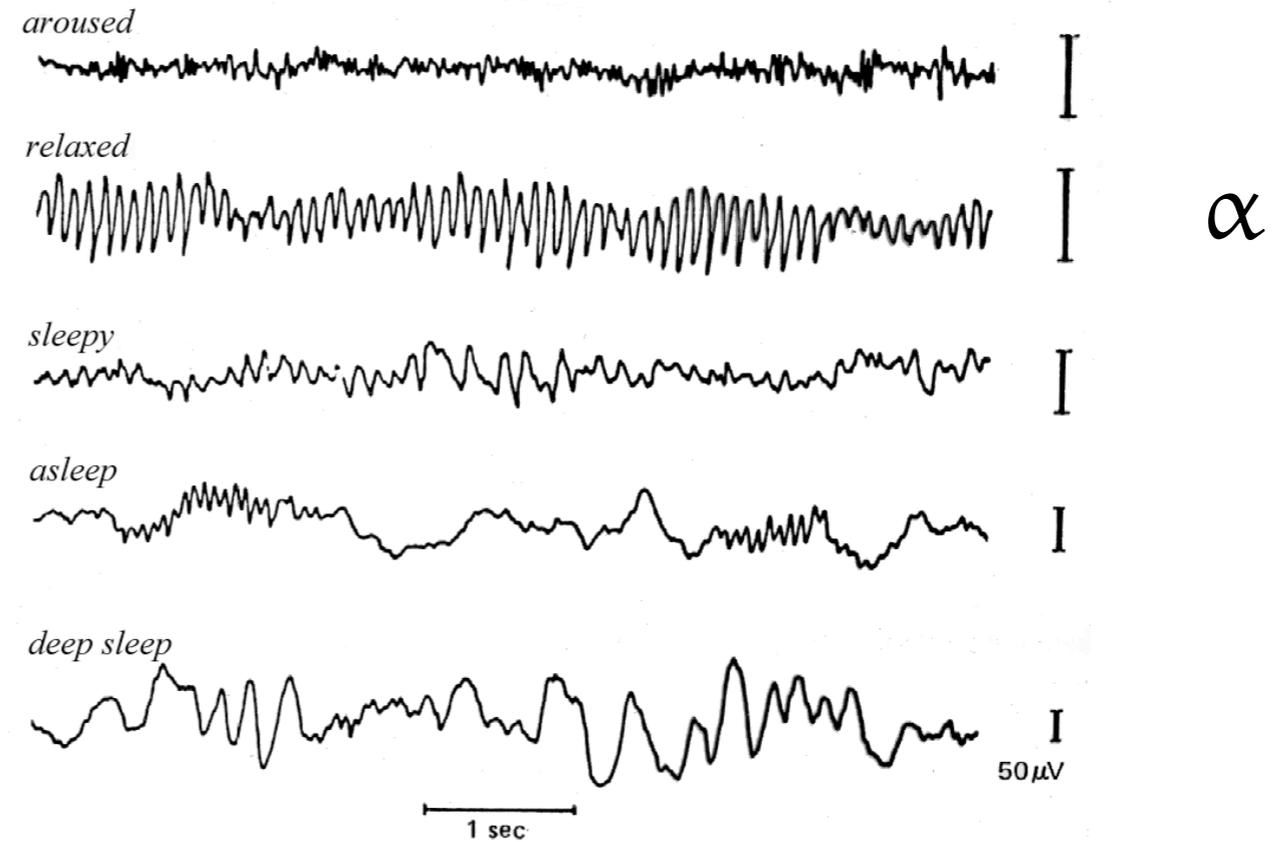


Santiago Ramón y Cajal
1900



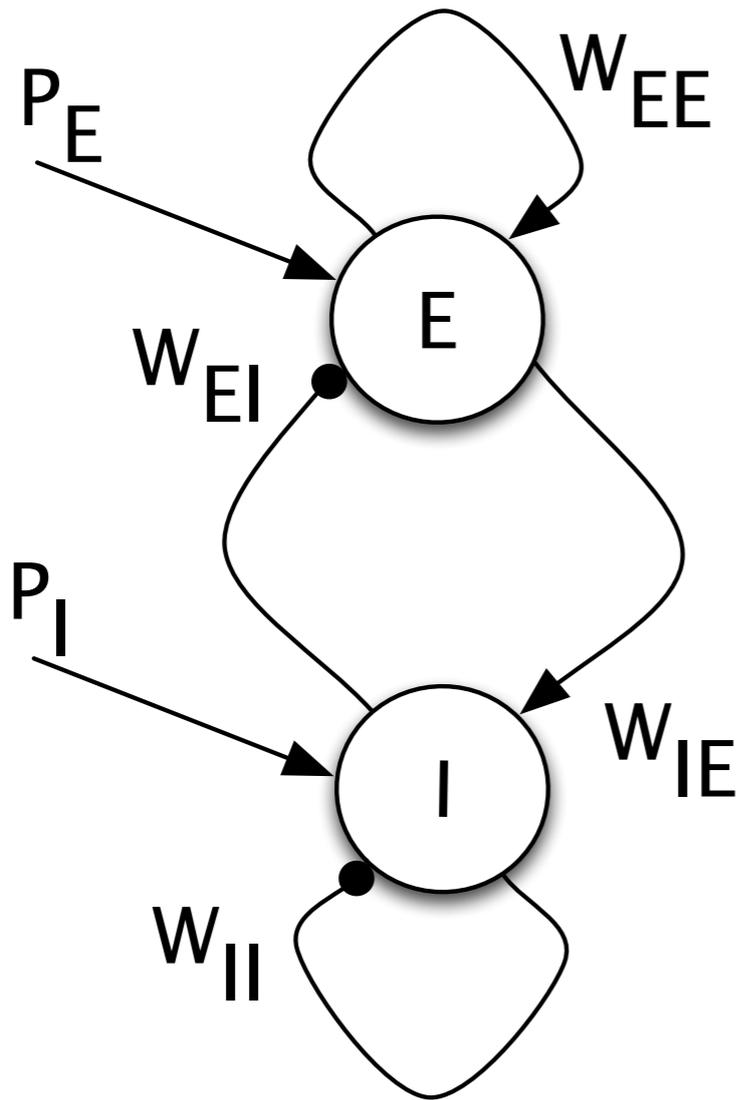
Eugene Izhikevich
2008

Electroencephalogram (EEG) power spectrum

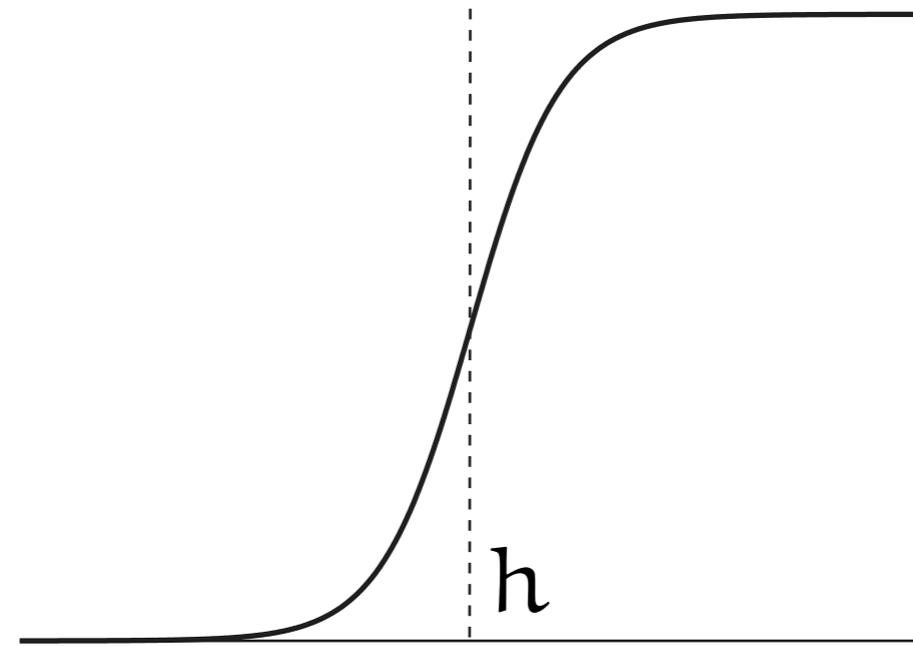


EEG records the activity of $\sim 10^6$ pyramidal neurons.

Population model



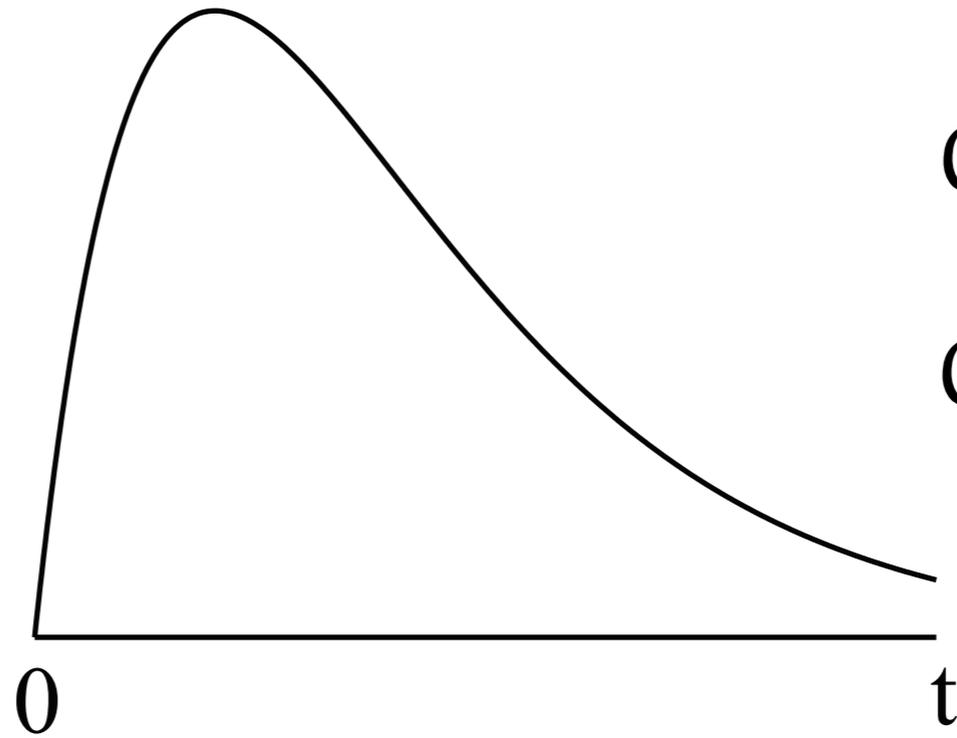
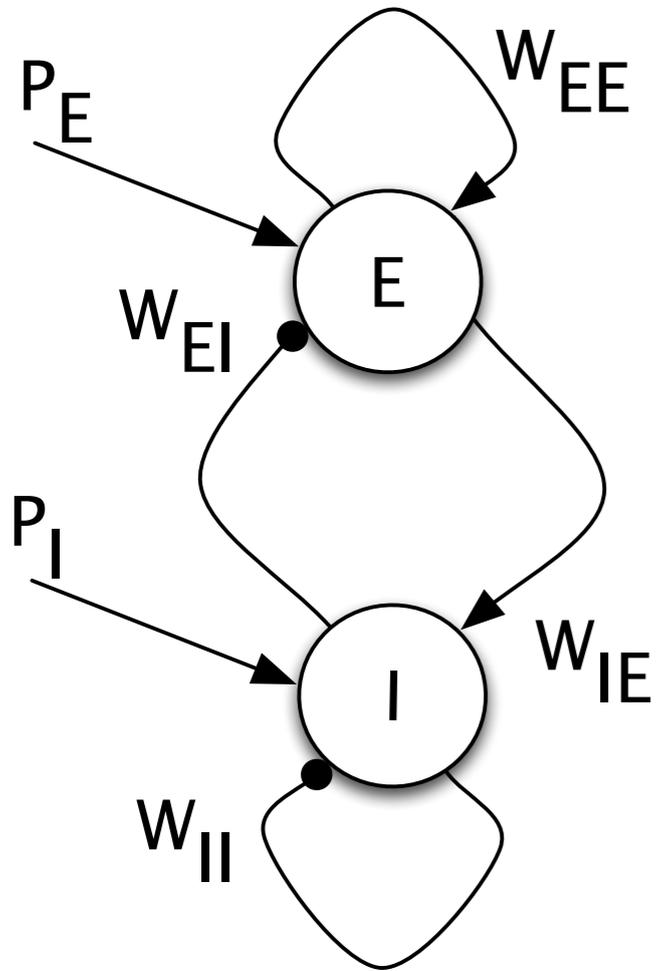
Firing rate activity $f(E)$



Firing rate activity $f(I)$

$$\dot{E} = -\frac{E}{\tau_E} + W_{EE}g_{EE}(A^+ - E) + W_{EI}g_{EI}(A^- - E) + P_E$$

$$\dot{I} = -\frac{I}{\tau_I} + W_{II}g_{II}(A^- - I) + W_{IE}g_{IE}(A^+ - I) + P_I$$



$$\eta(t) = \alpha^2 t e^{-\alpha t}$$

$$Q\eta = \delta$$

$$Q = \left(1 + \frac{1}{\alpha} \frac{d}{dt} \right)^2$$

$$Qg_{jE} = f(E)$$

$$Qg_{jI} = f(I)$$

Steady state approximation

$$E = E(g_{EE}, g_{EI})$$

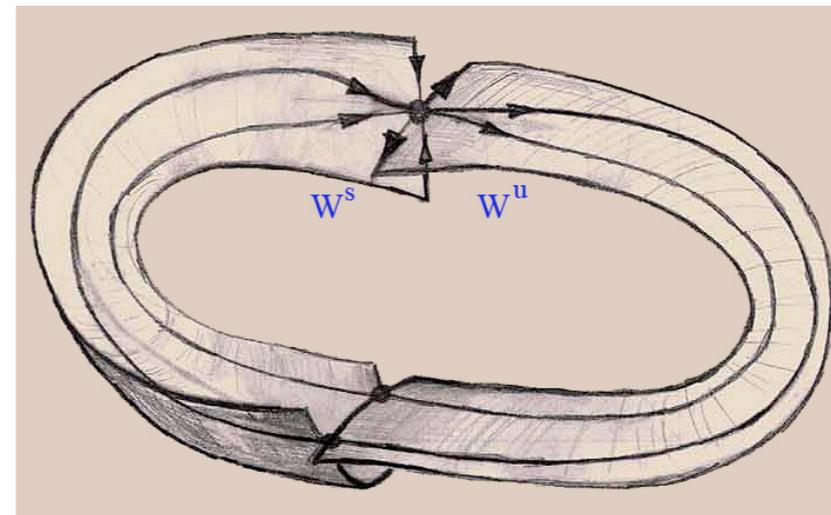
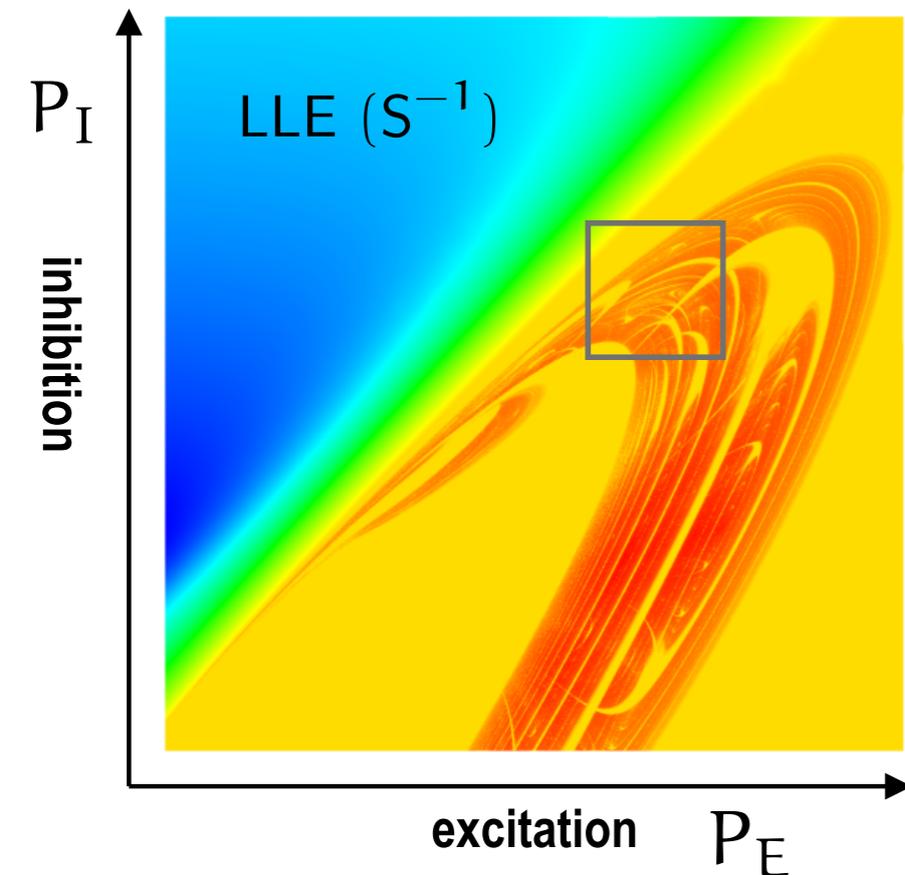
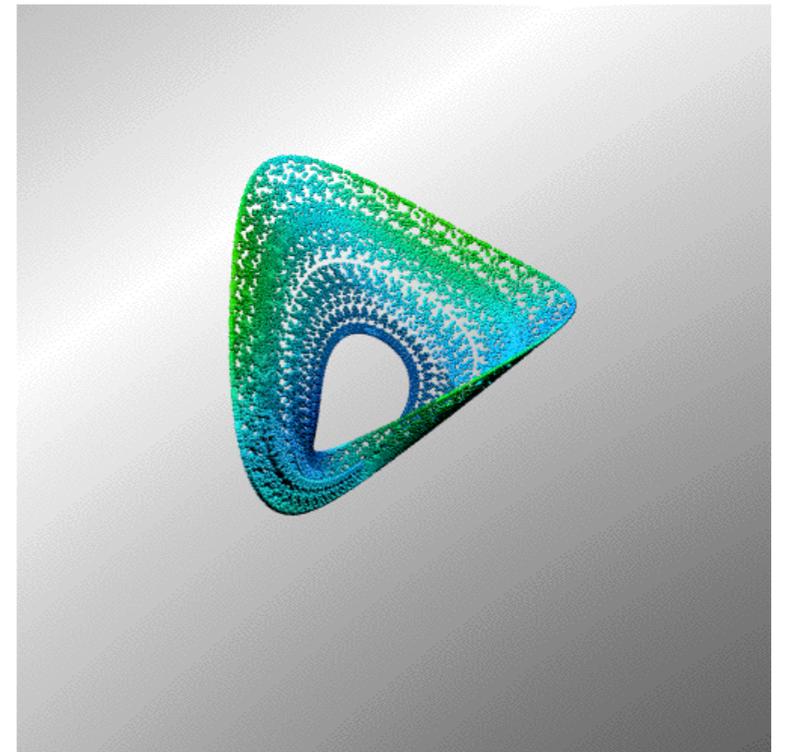
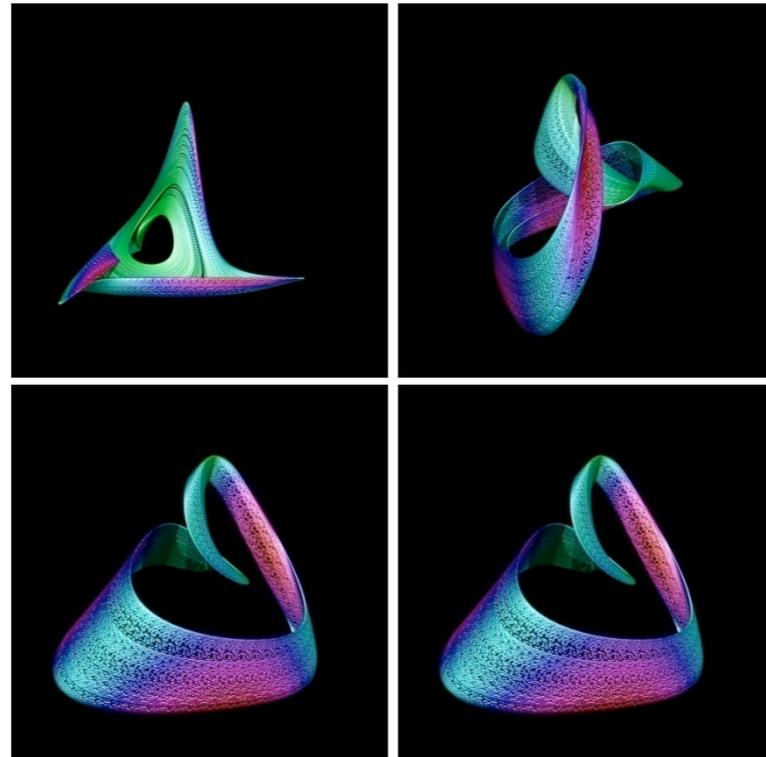
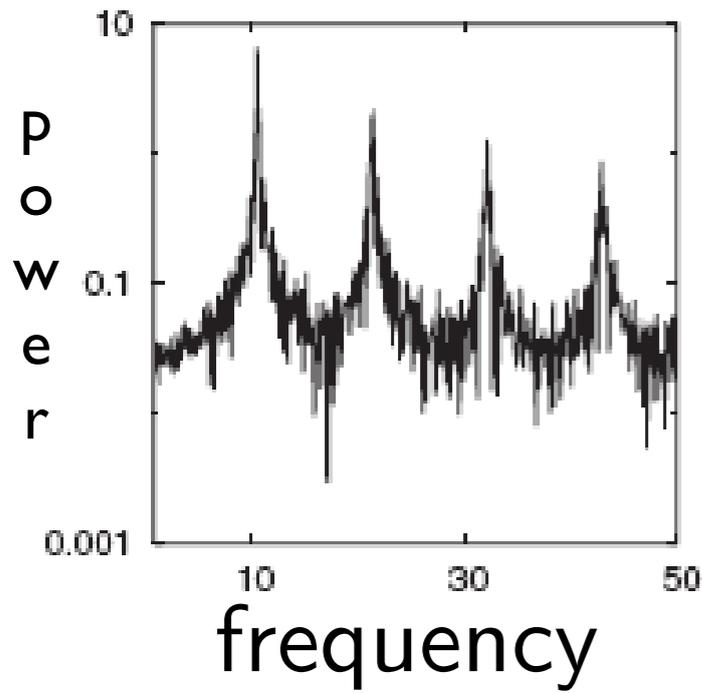
$$I = I(g_{II}, g_{IE})$$

$$Qg = f$$

$$f = f(\{g\})$$

$$g = \eta * f$$

Alphoid chaos (10 D)

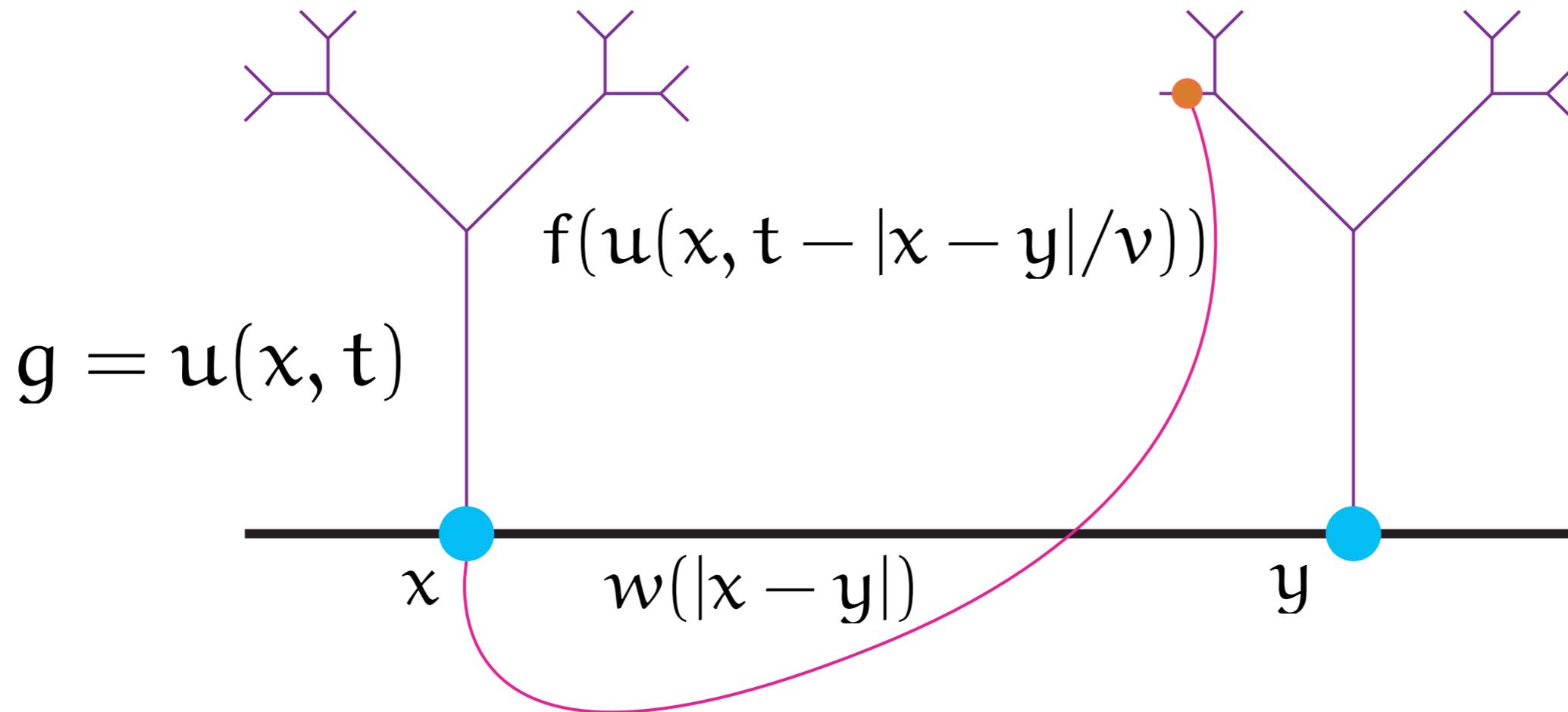


Shilnikov saddle-node route to chaos
van Veen and Liley, PRL, **97**, 208101 (2006)

Spatially extended models

$$g = w \otimes \eta * f$$

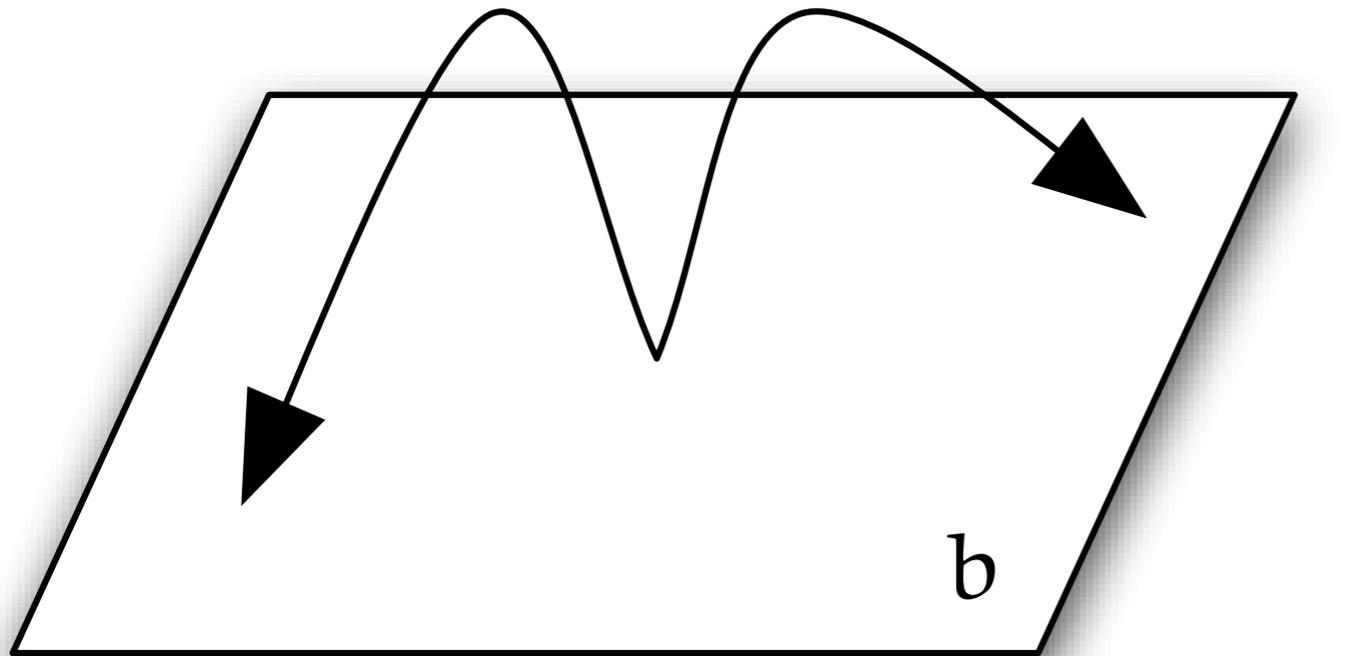
Simplest neural field model: Wilson-Cowan ('72), Amari ('77)



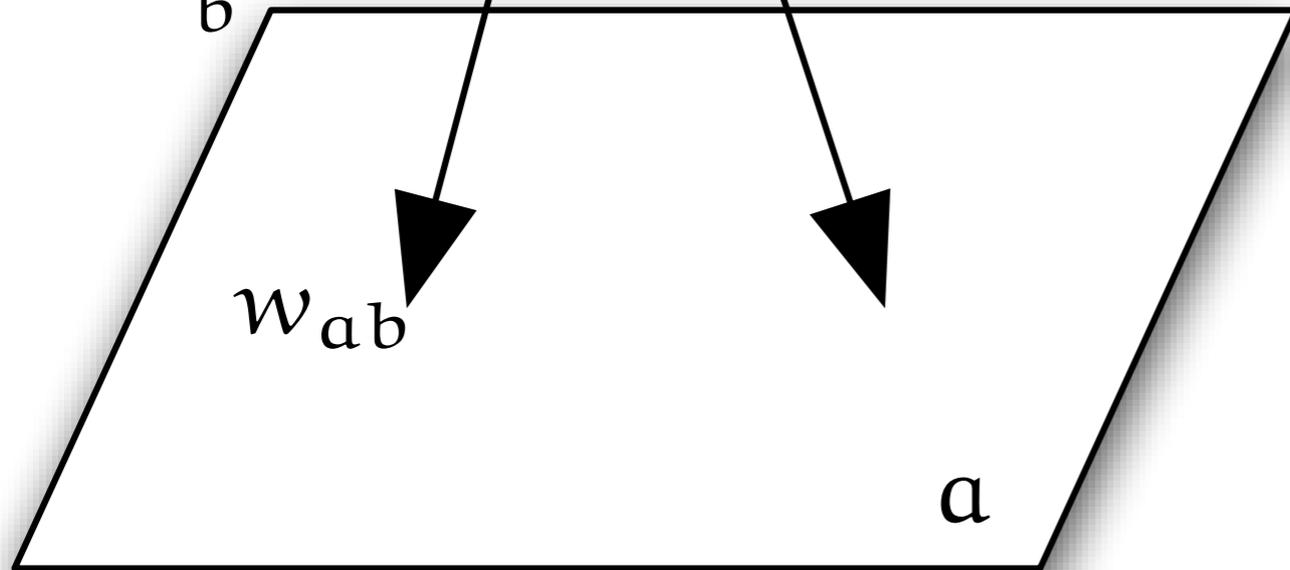
$$u(x, t) = \int_{-\infty}^{\infty} dy w(y) \int_0^{\infty} ds \eta(s) f(u(x - y, t - s - |y|/v))$$

2D layers

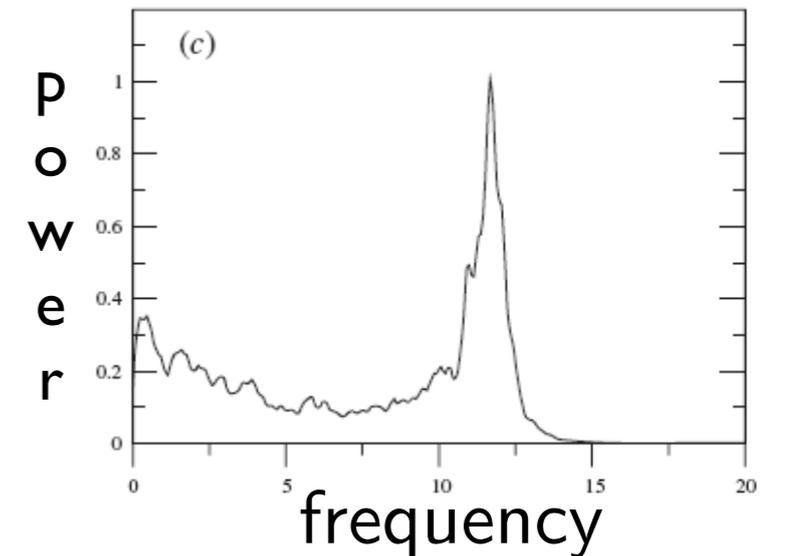
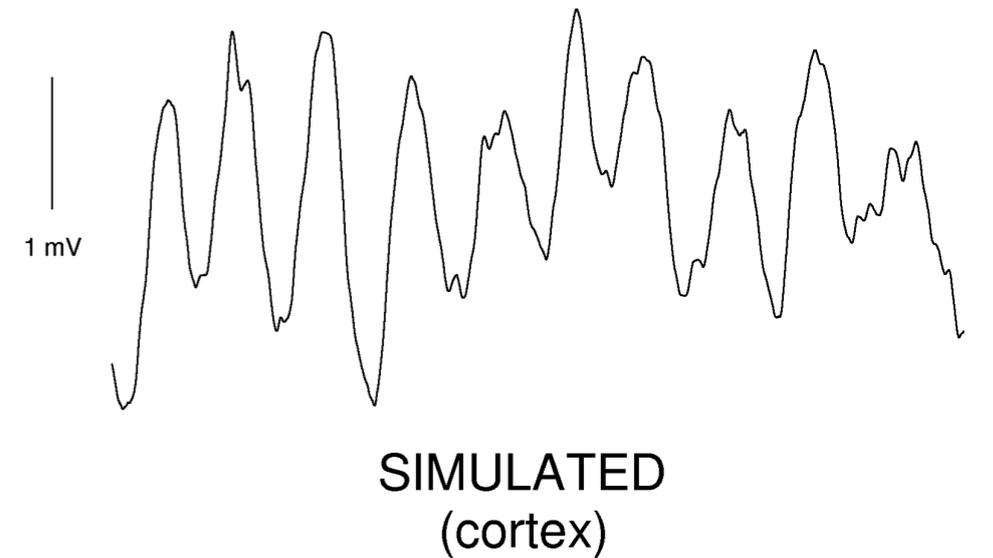
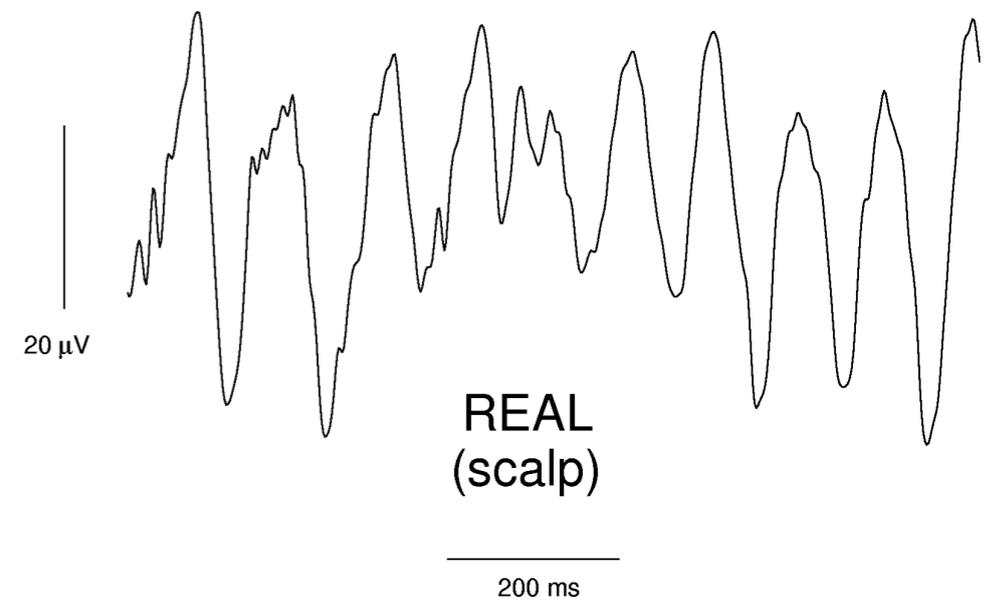
$$u_{ab} = \eta_{ab} * \psi_{ab}$$



$$h_a = \sum_b u_{ab}$$

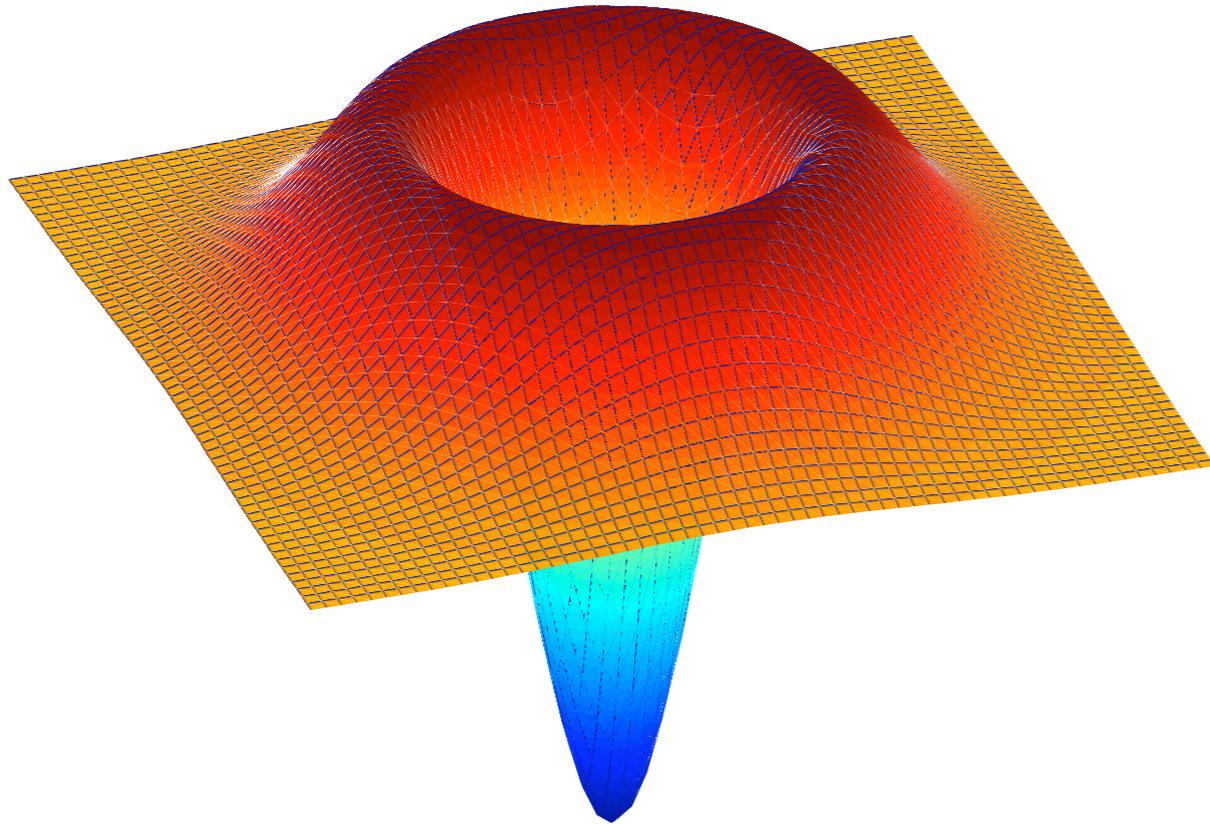


$$\psi_{ab}(\mathbf{r}, t) = \int_{\mathbb{R}^2} d\mathbf{r}' w_{ab}(\mathbf{r}, \mathbf{r}') f_b \circ h_b(\mathbf{r}', t - |\mathbf{r} - \mathbf{r}'|/v_{ab})$$



Turing instability analysis

E layer and I layer



$$e^{i\mathbf{k}\cdot\mathbf{r}} e^{\lambda t}$$

Continuous spectrum

$$\det(\mathcal{D}(\mathbf{k}, \lambda) - I) = 0$$

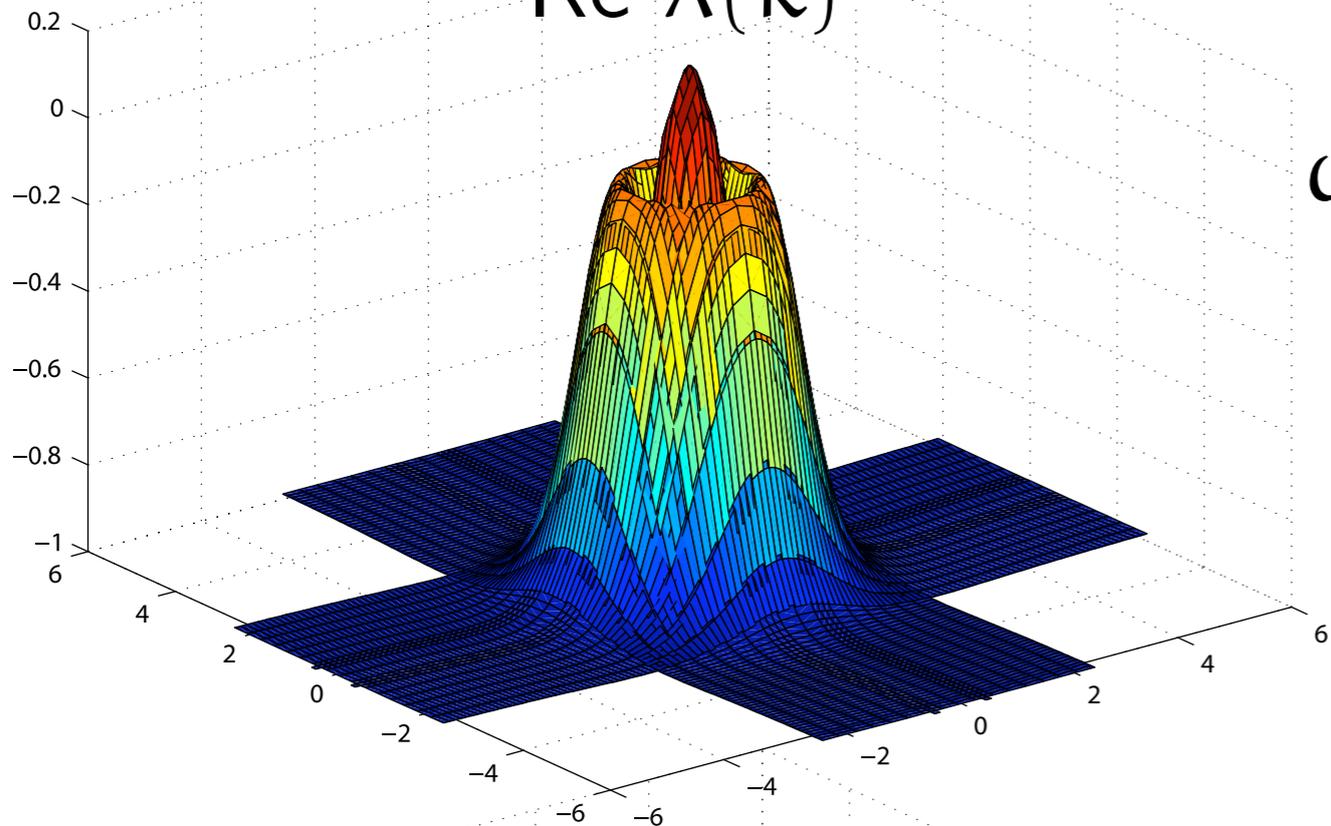
$$[\mathcal{D}(\mathbf{k}, \lambda)]_{ab} = \tilde{\eta}_{ab}(\lambda) G_{ab}(\mathbf{k}, -i\lambda) \gamma_b$$

$$\tilde{\eta} = \text{LT } \eta$$

$$G = \text{FLT } w(r) \delta(t - r/v)$$

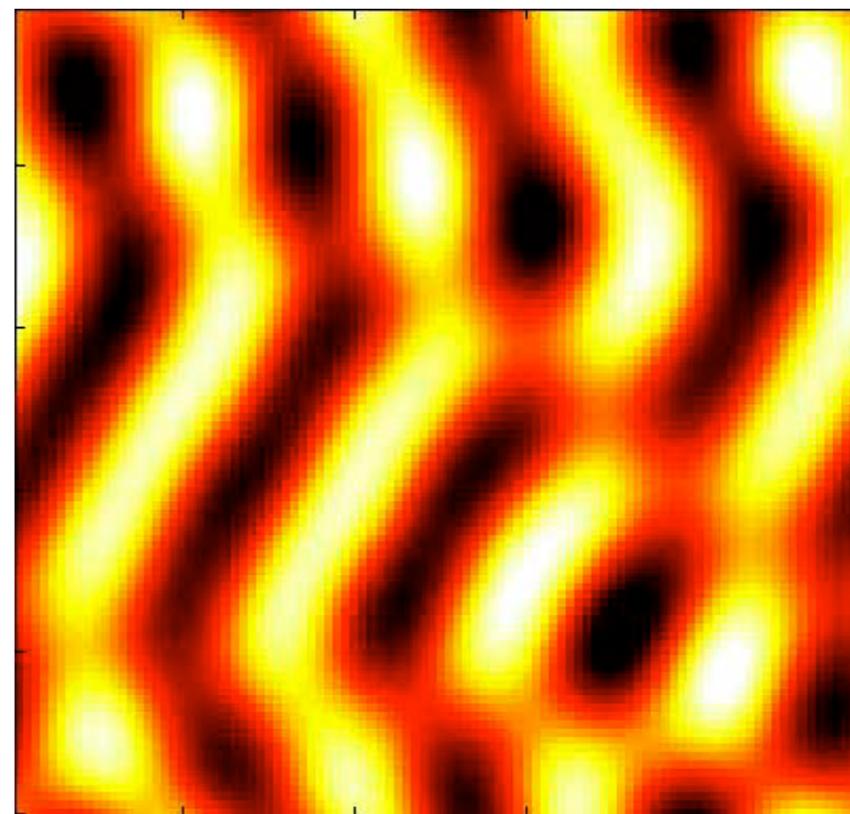
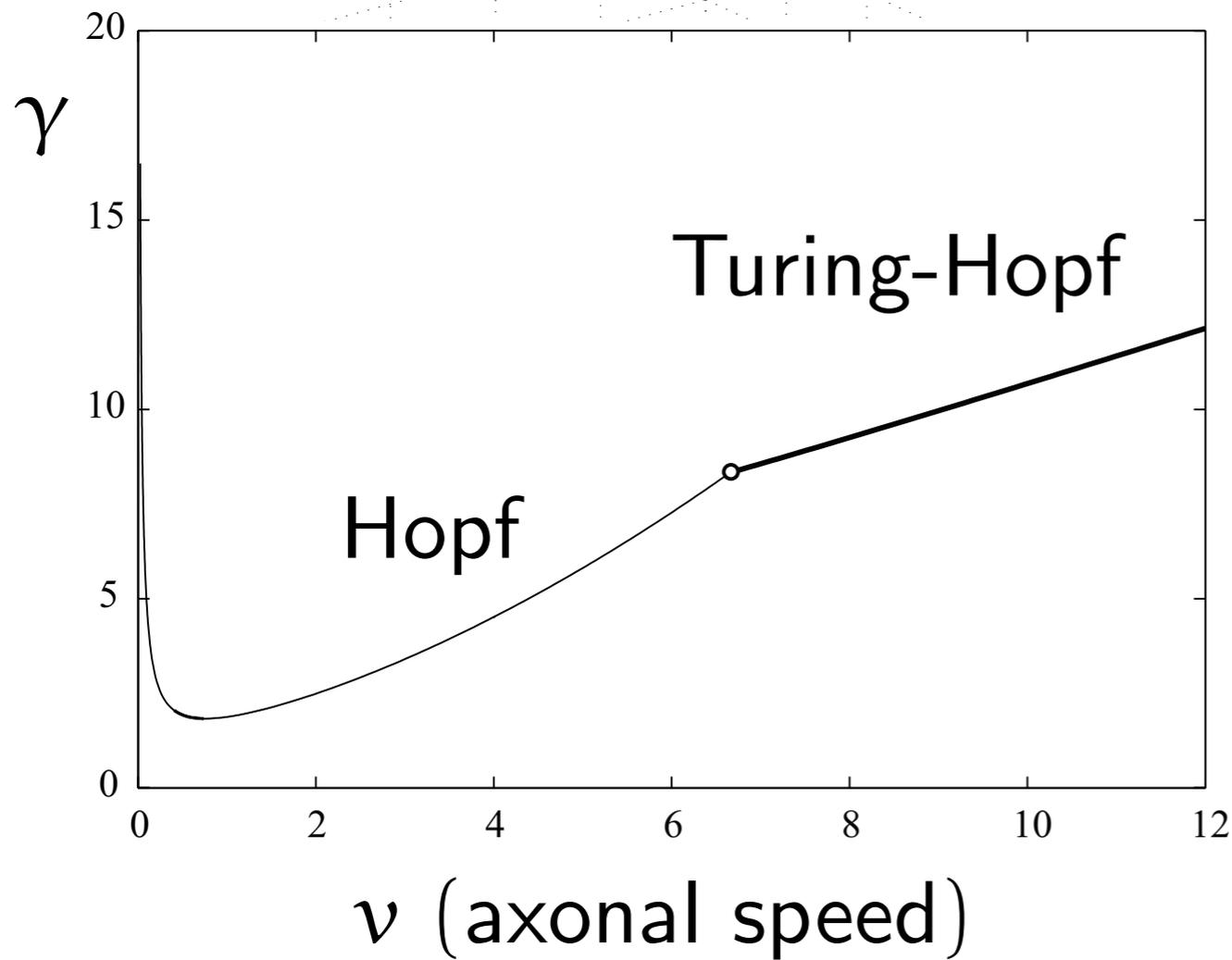
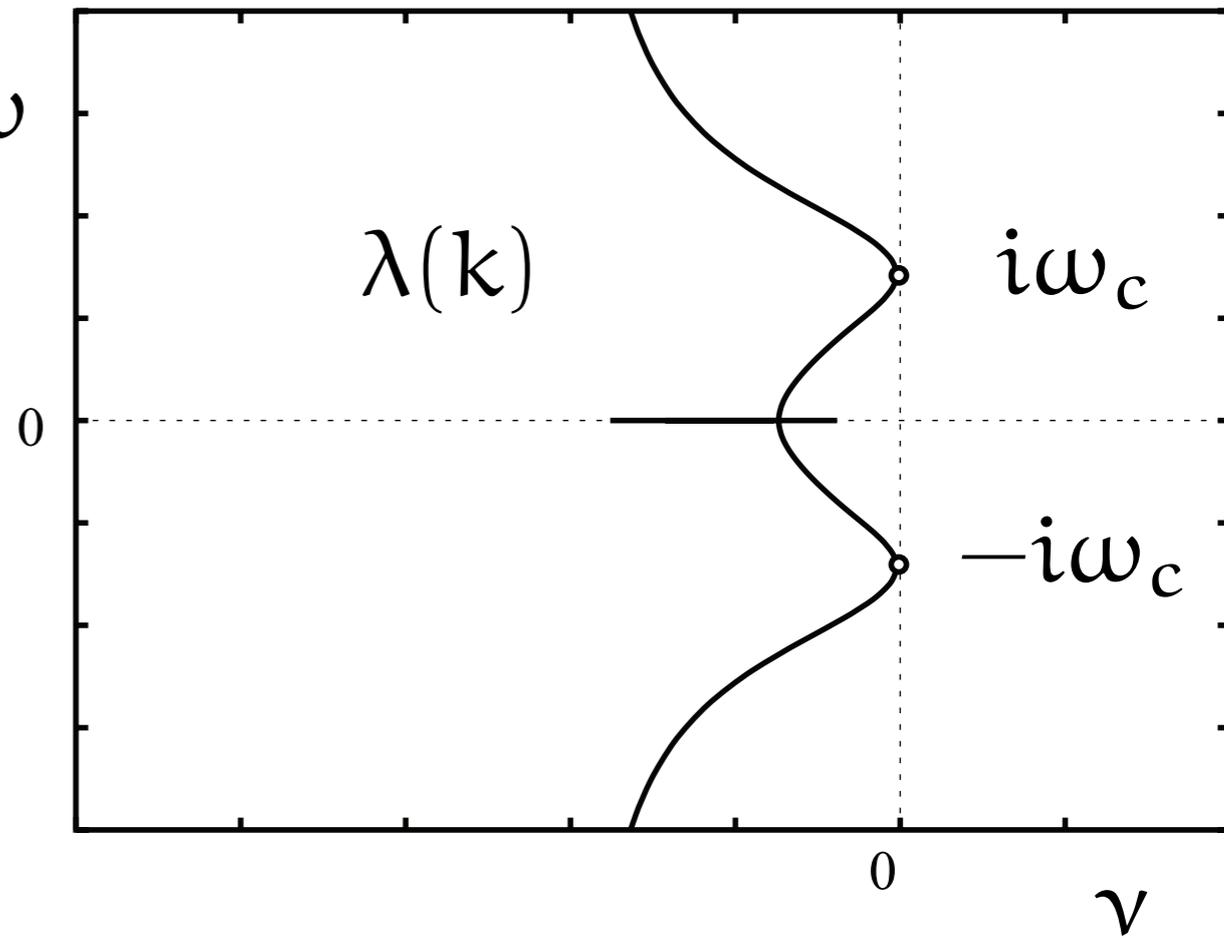
$$\gamma = f'(ss)$$

$\text{Re } \lambda(\mathbf{k})$



$$\lambda = \nu + i\omega$$

ω



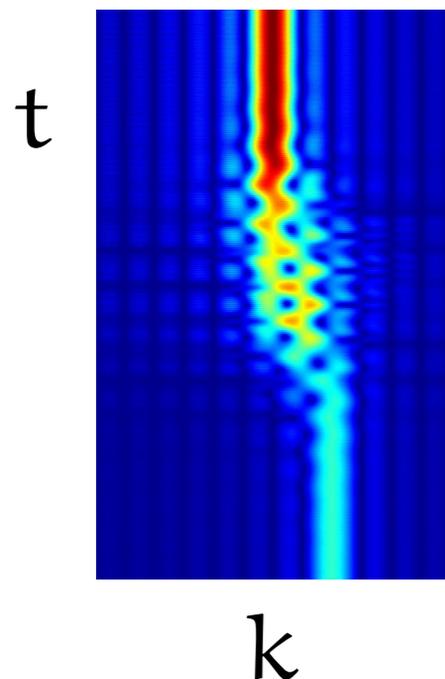
Amplitude Equations (one D)

Coupled mean-field Ginzburg–Landau equations describing a Turing–Hopf bifurcation with modulation group velocity of $O(1)$.

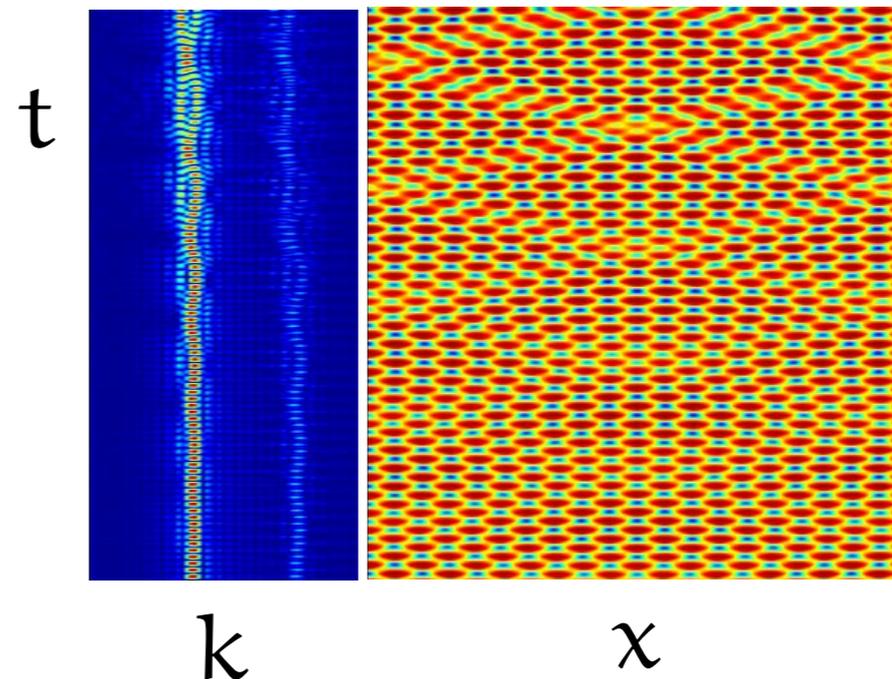
$$\frac{\partial A_1}{\partial \tau} = A_1 (a + b|A_1|^2 + c\langle |A_2|^2 \rangle) + d \frac{\partial^2 A_1}{\partial \xi_+^2}$$

$$\frac{\partial A_2}{\partial \tau} = A_2 (a + b|A_2|^2 + c\langle |A_1|^2 \rangle) + d \frac{\partial^2 A_2}{\partial \xi_-^2}$$

Benjamin–Feir (BF)

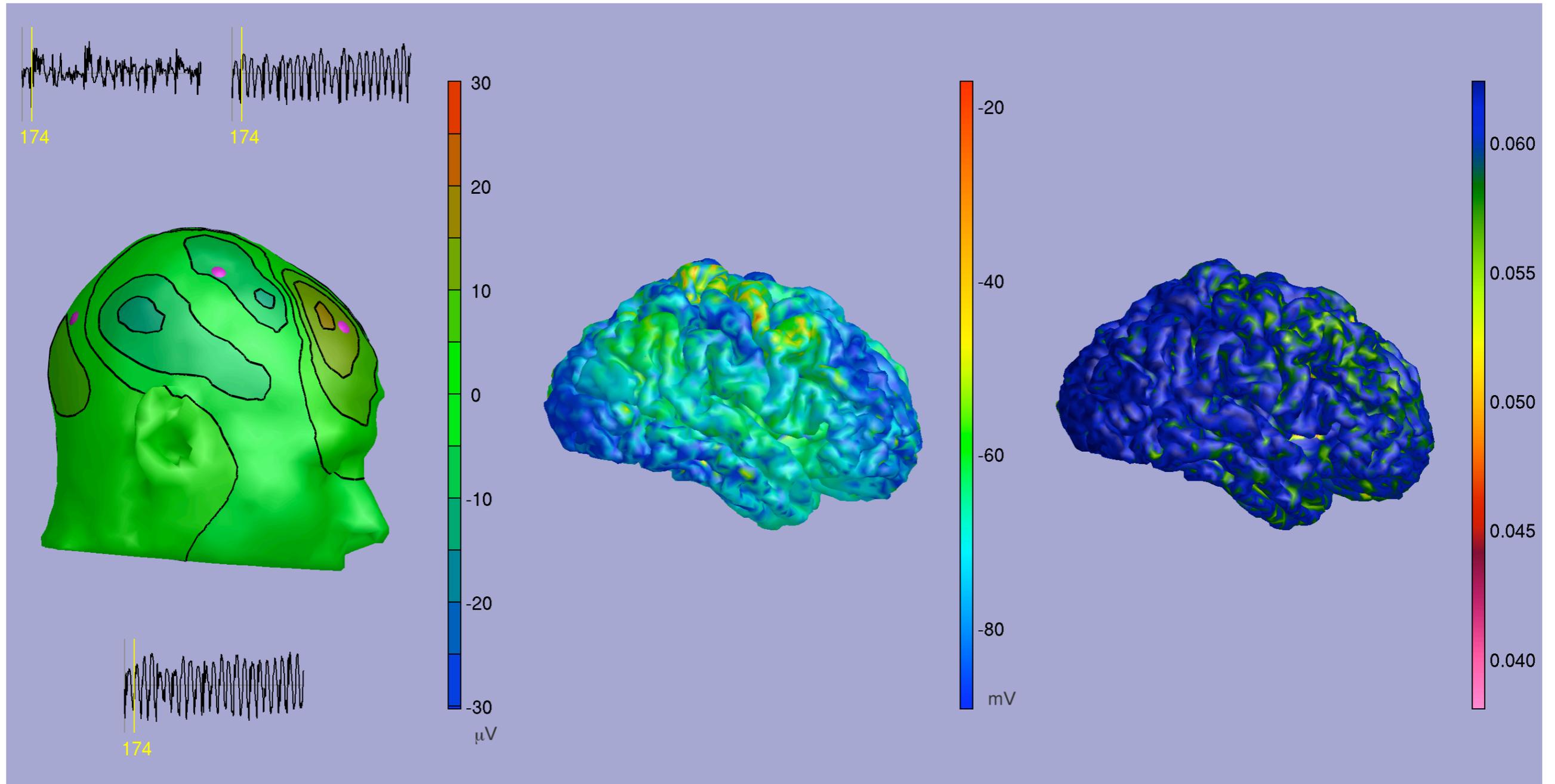


BF-Eckhaus instability



Coefficients in terms of integral transforms of w and η .

Applications to co-registered EEG/fMRI

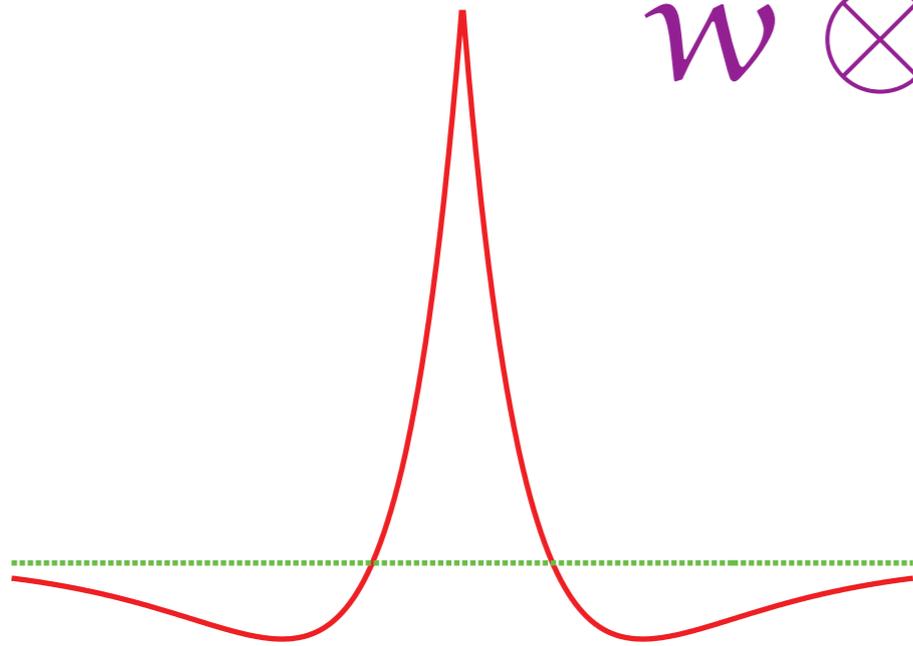


Bojak, I., Oostendorp, T. F., Reid, A. T., Kotter, R., 2009. Realistic mean field forward predictions for the integration of co-registered EEG/fMRI. BMC Neuroscience 10, L2.

Time independent localised solutions

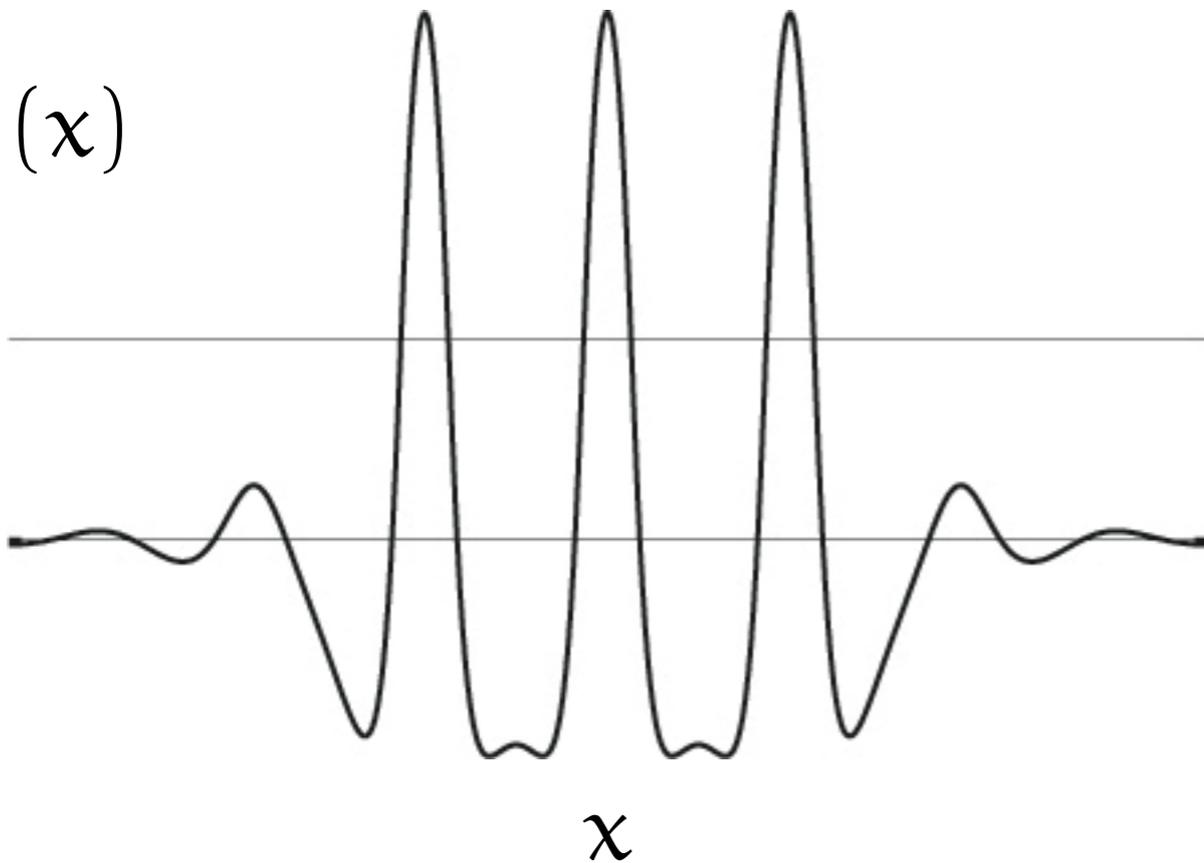
$$w \otimes \eta * f \rightarrow w \otimes f$$

$$q(x) = \int_{\mathbb{R}} dy w(x-y) f \circ q(y)$$

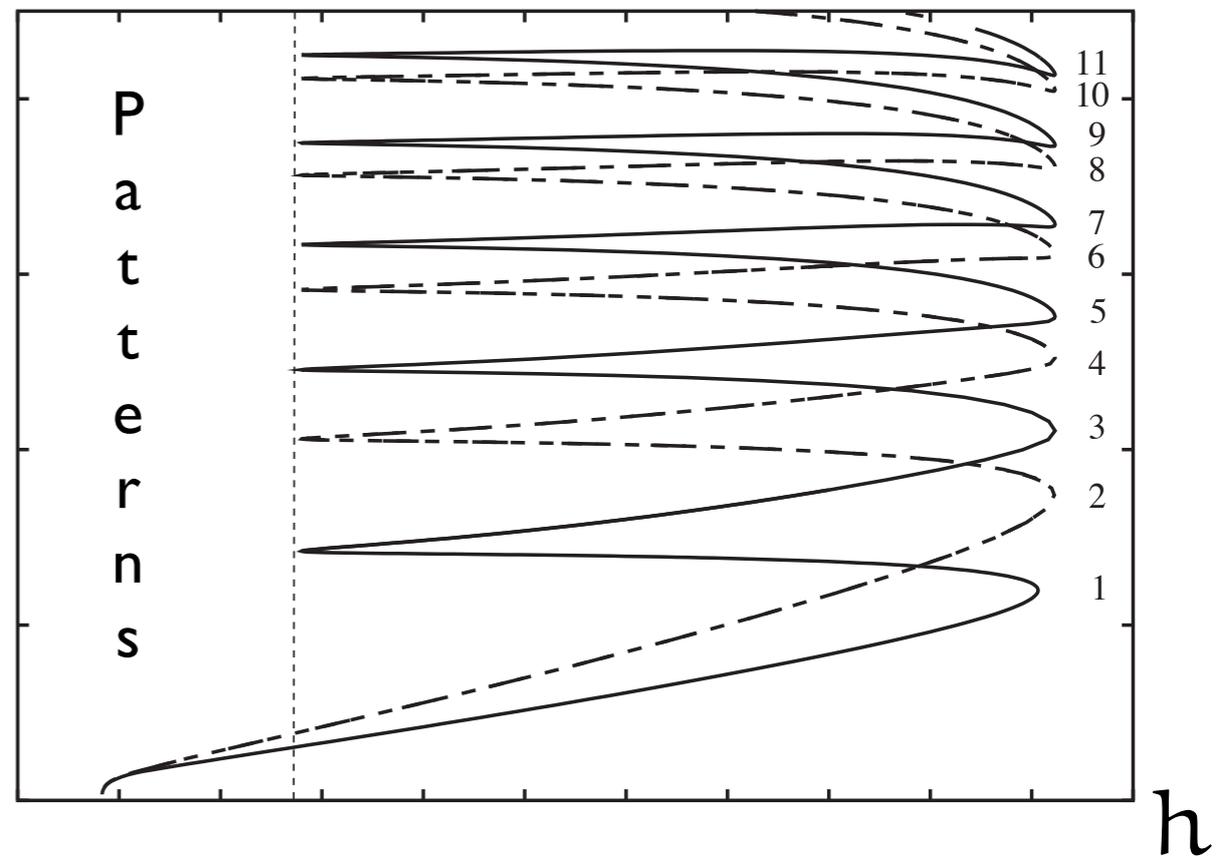


$$w(x) = (1 - |x|)e^{-|x|}$$

$q(x)$

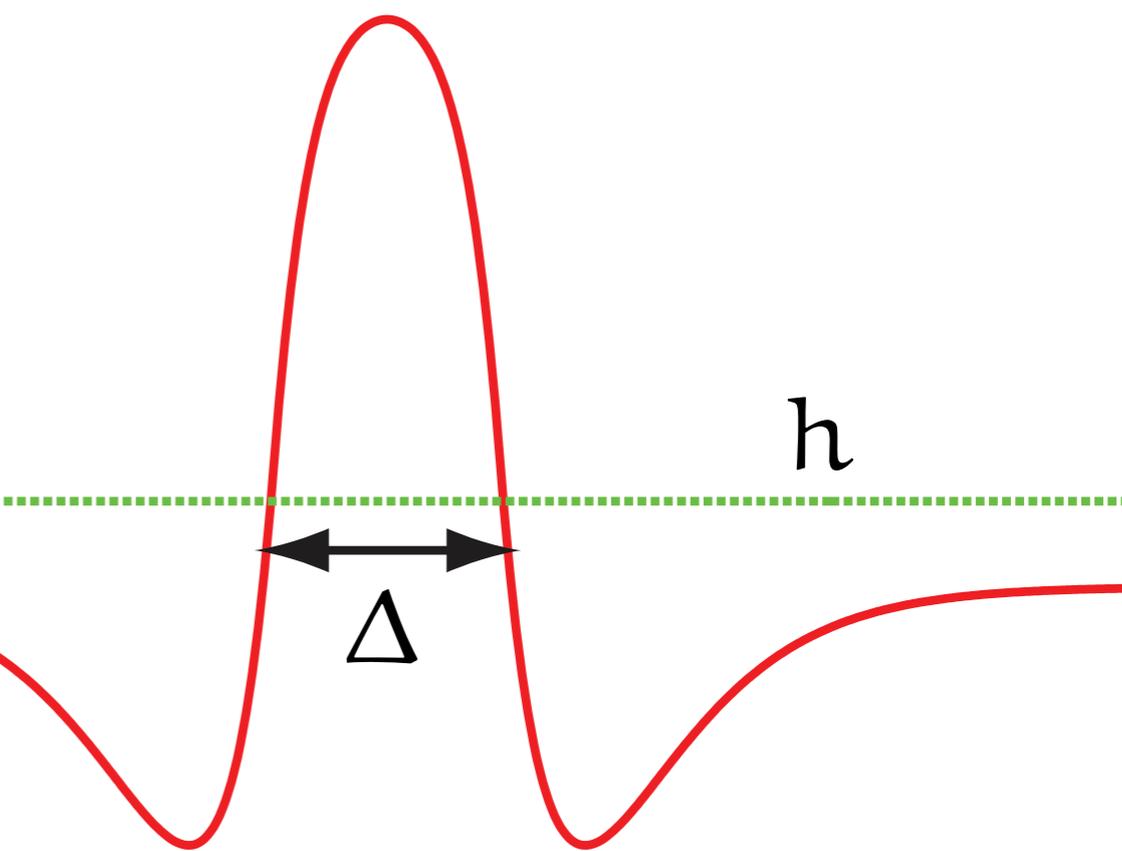
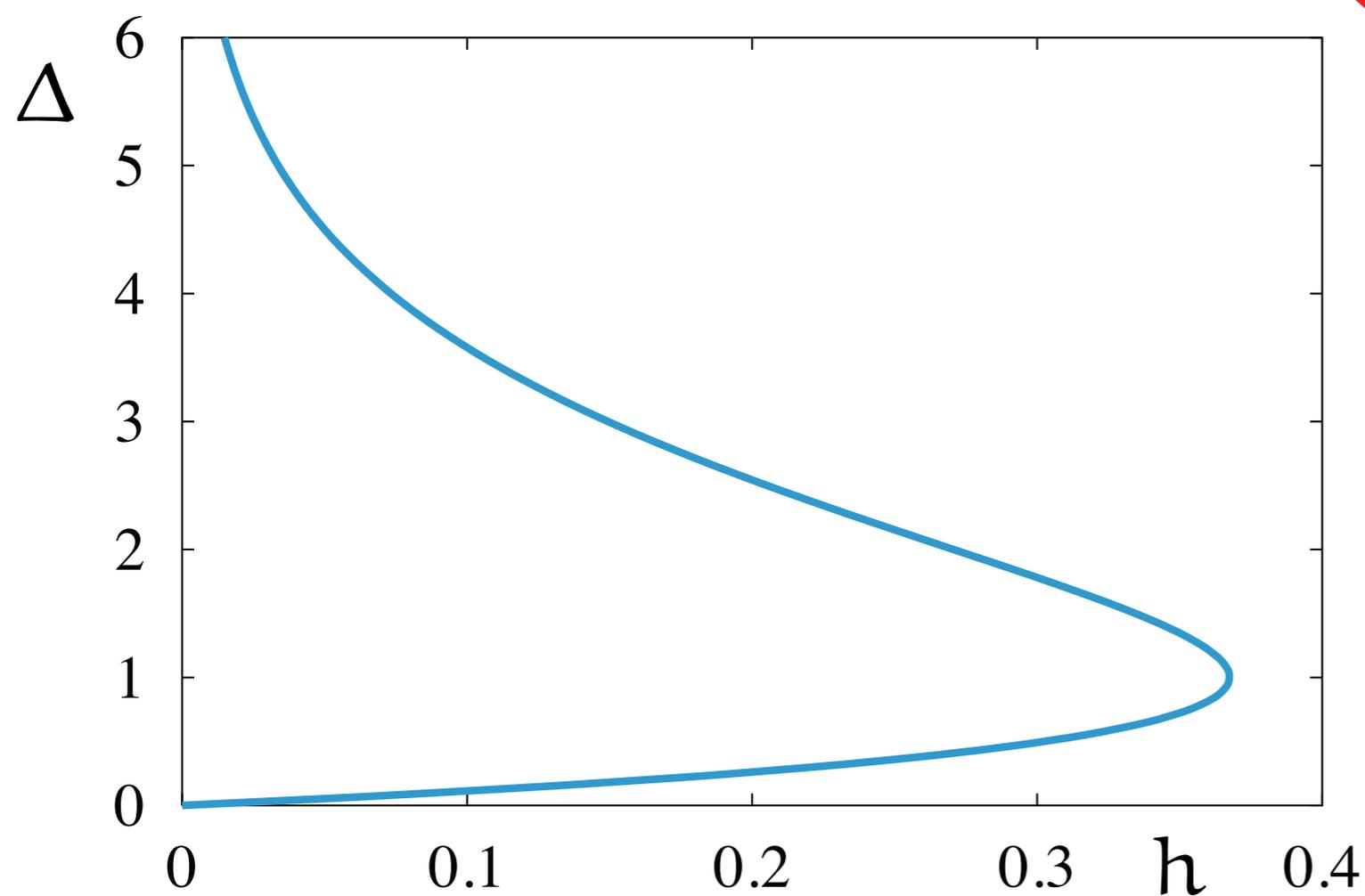


L_2



Exact result for I-bump: $f(u) = H(u - h)$

$$q(x) = \int_0^\Delta dy w(x - y)$$



$$q(0) = h = q(\Delta)$$

$$\Delta e^{-\Delta} = h$$

working memory

Stability

Examine eigenspectrum of the linearization about a solu

Solutions of form $u(x)e^{\lambda t}$ satisfy $\mathcal{L}u(x) = u(x)$

$$\mathcal{L}u(x) = \tilde{\eta}(\lambda) \int_{-\infty}^{\infty} dy w(x-y) f'(q(y) - h) u(y)$$

For Heaviside firing rate

$$f'(q(x)) = \frac{\delta(x)}{|q'(0)|} + \frac{\delta(x - \Delta)}{|q'(\Delta)|}$$

so

$$u(x) = \frac{\tilde{\eta}(\lambda)}{|w(0) - w(\Delta)|} [w(x)u(0) + w(x - \Delta)u(\Delta)]$$

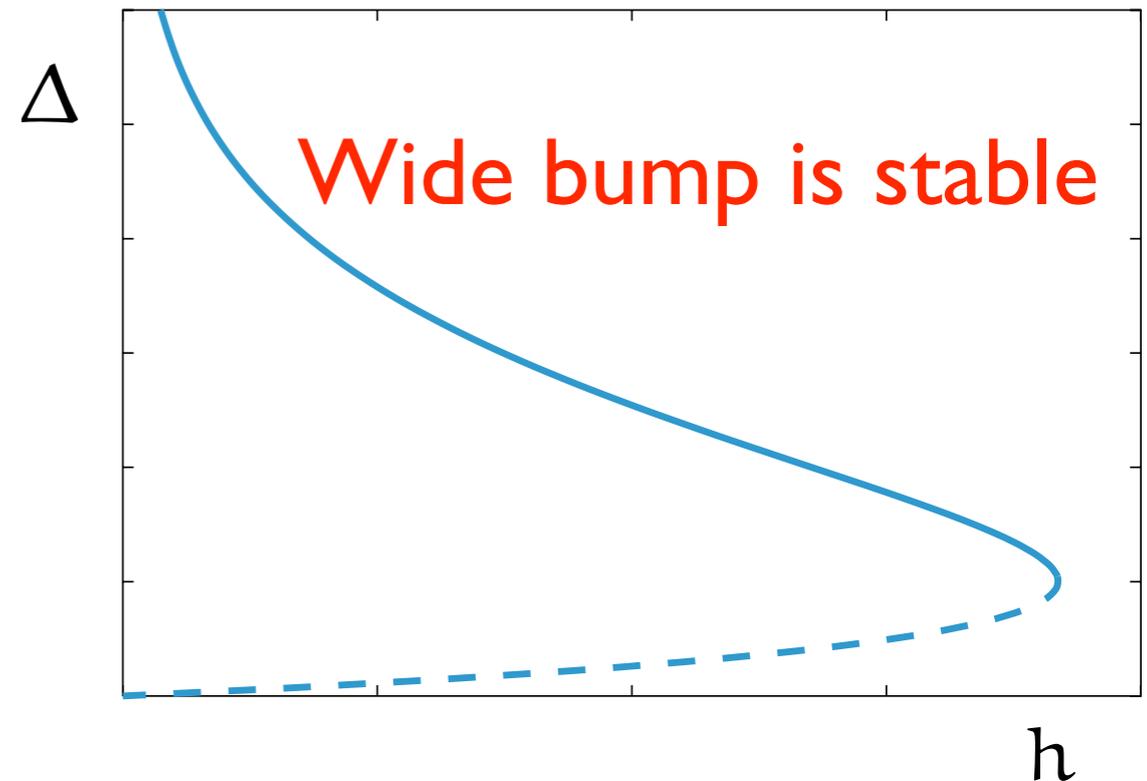
System of linear equations for perturbations at threshold

$$\begin{bmatrix} u(0) \\ u(\Delta) \end{bmatrix} = \mathcal{A}(\lambda) \begin{bmatrix} u(0) \\ u(\Delta) \end{bmatrix}, \quad \mathcal{A}(\lambda) = \frac{\tilde{\eta}(\lambda)}{|w(0) - w(\Delta)|} \begin{bmatrix} w(0) & w(\Delta) \\ w(\Delta) & w(0) \end{bmatrix}$$

Non trivial solution if

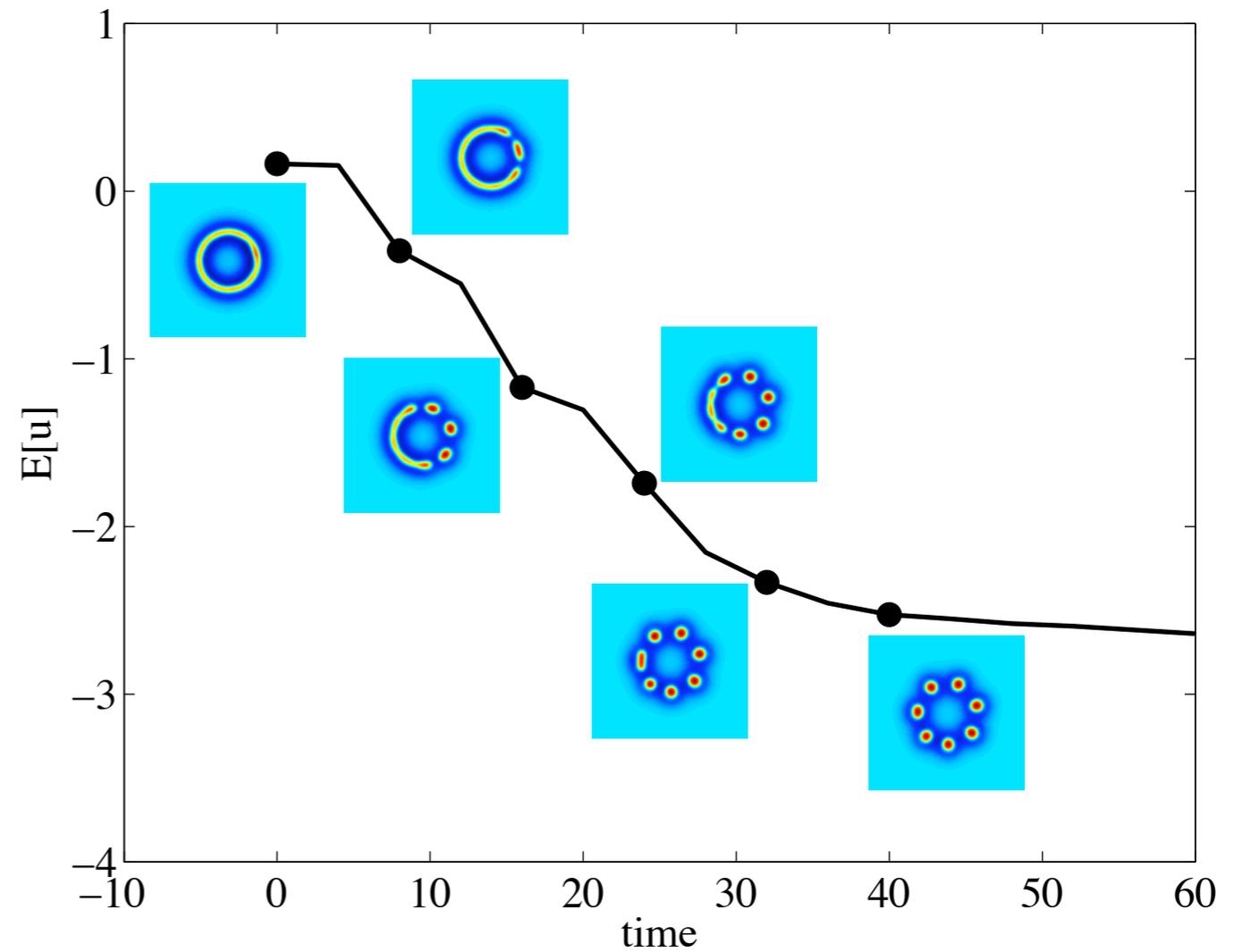
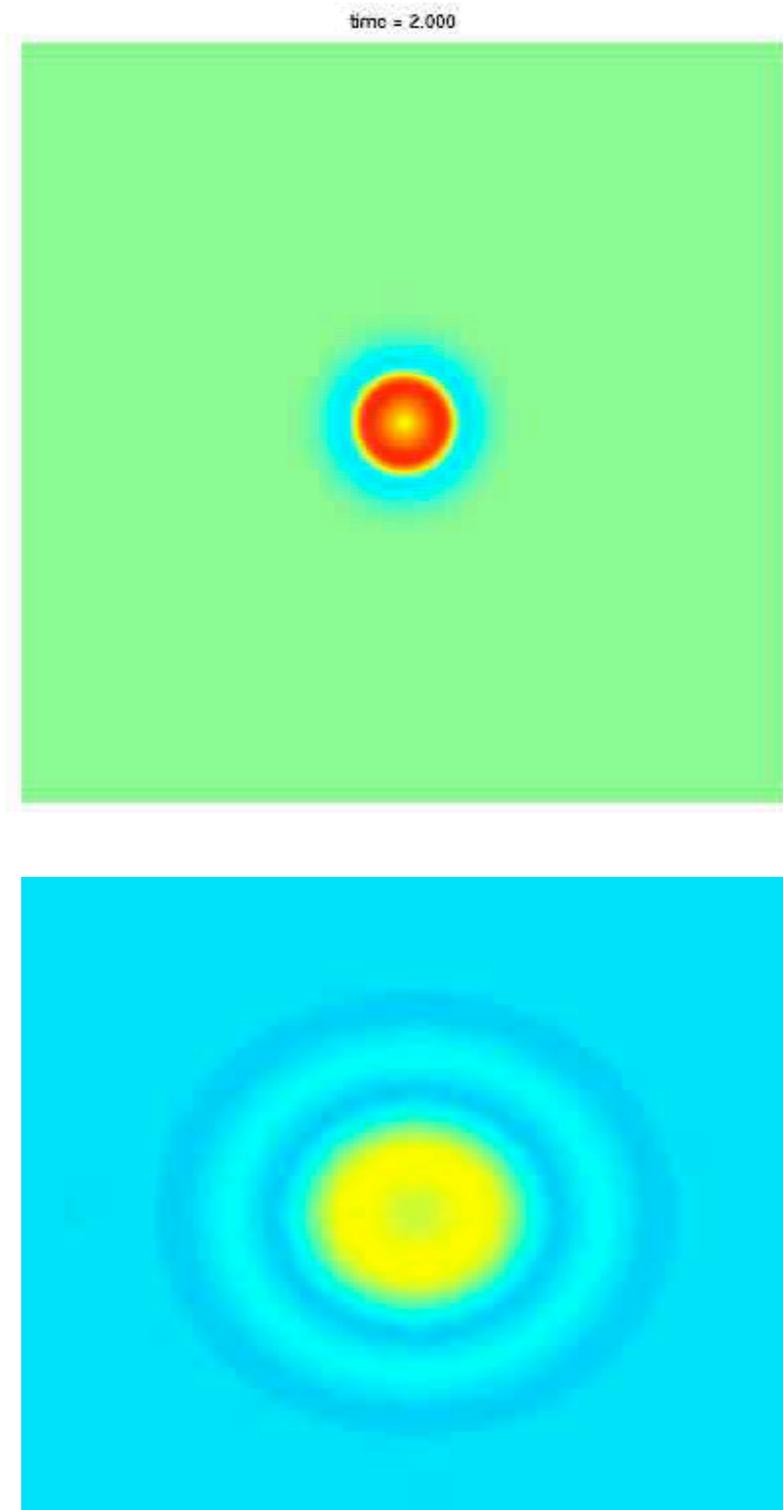
$$\mathcal{E}(\lambda) = \det(\mathcal{A}(\lambda) - I) = 0$$

Solutions stable if $\text{Re } \lambda < 0$



Evans function for integral neural field equation

Predictions of Evans function



Threshold accommodation

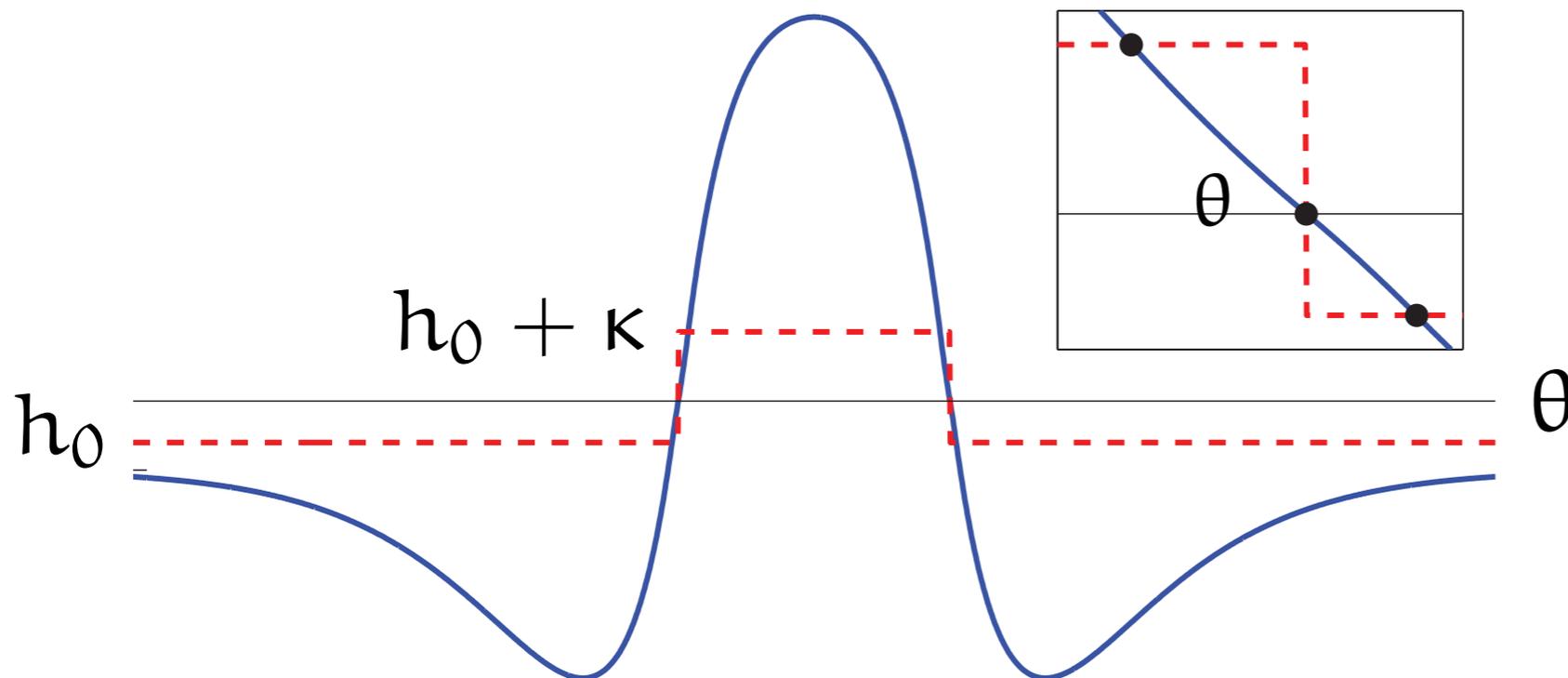
Hill (1936), “... the threshold rises when the *local potential* is maintained ... and reverts gradually to its original value when the nerve is allowed to rest.”

$$\frac{\partial h}{\partial t} = -(h - h_0) + \kappa H(u - \theta)$$

One bump $(u, h) = (q(x), p(x))$

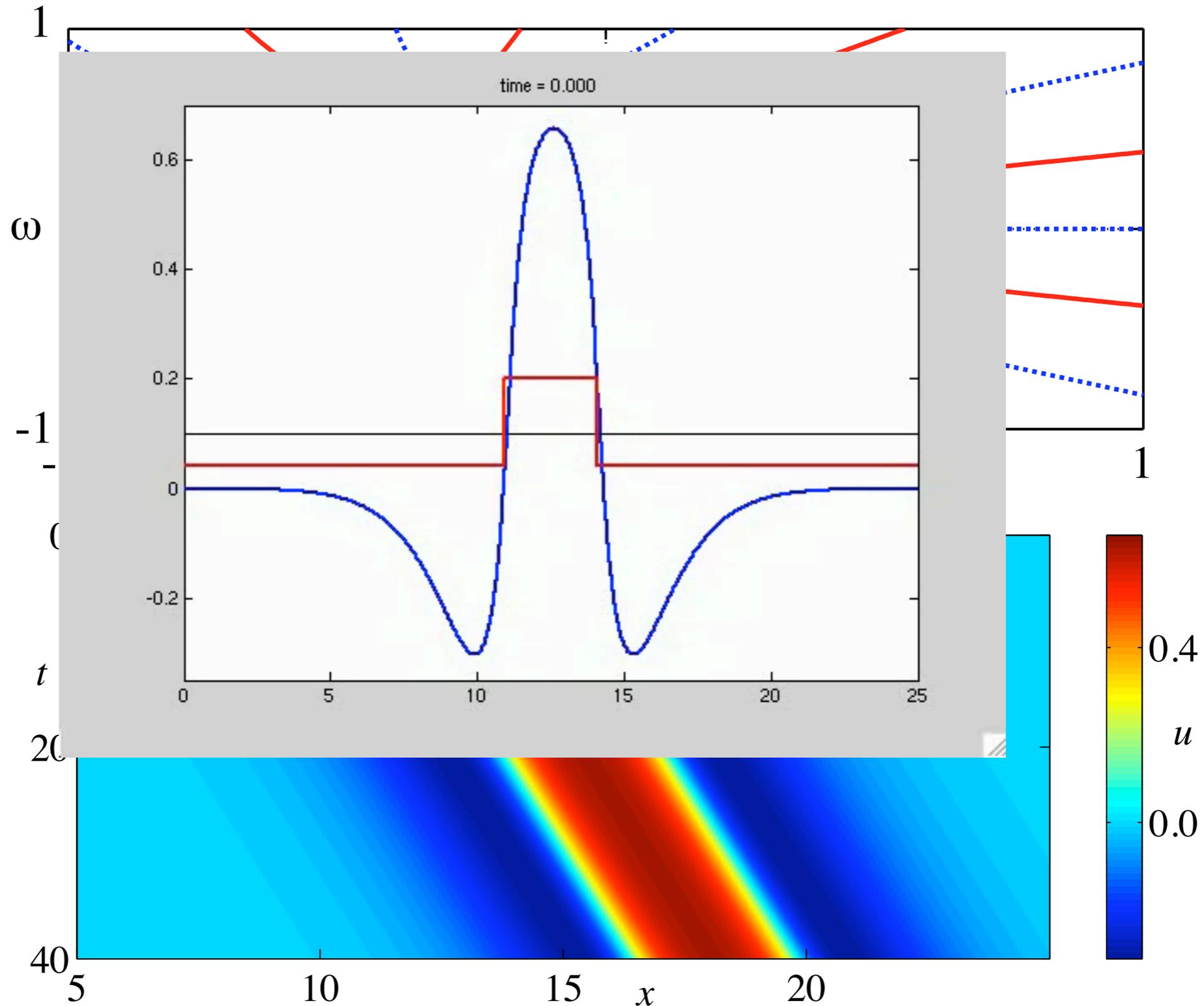
$$q = w \otimes H(q - p)$$

$$p = \begin{cases} h_0 & q < \theta \\ h_0 + \kappa & q \geq \theta \end{cases}$$



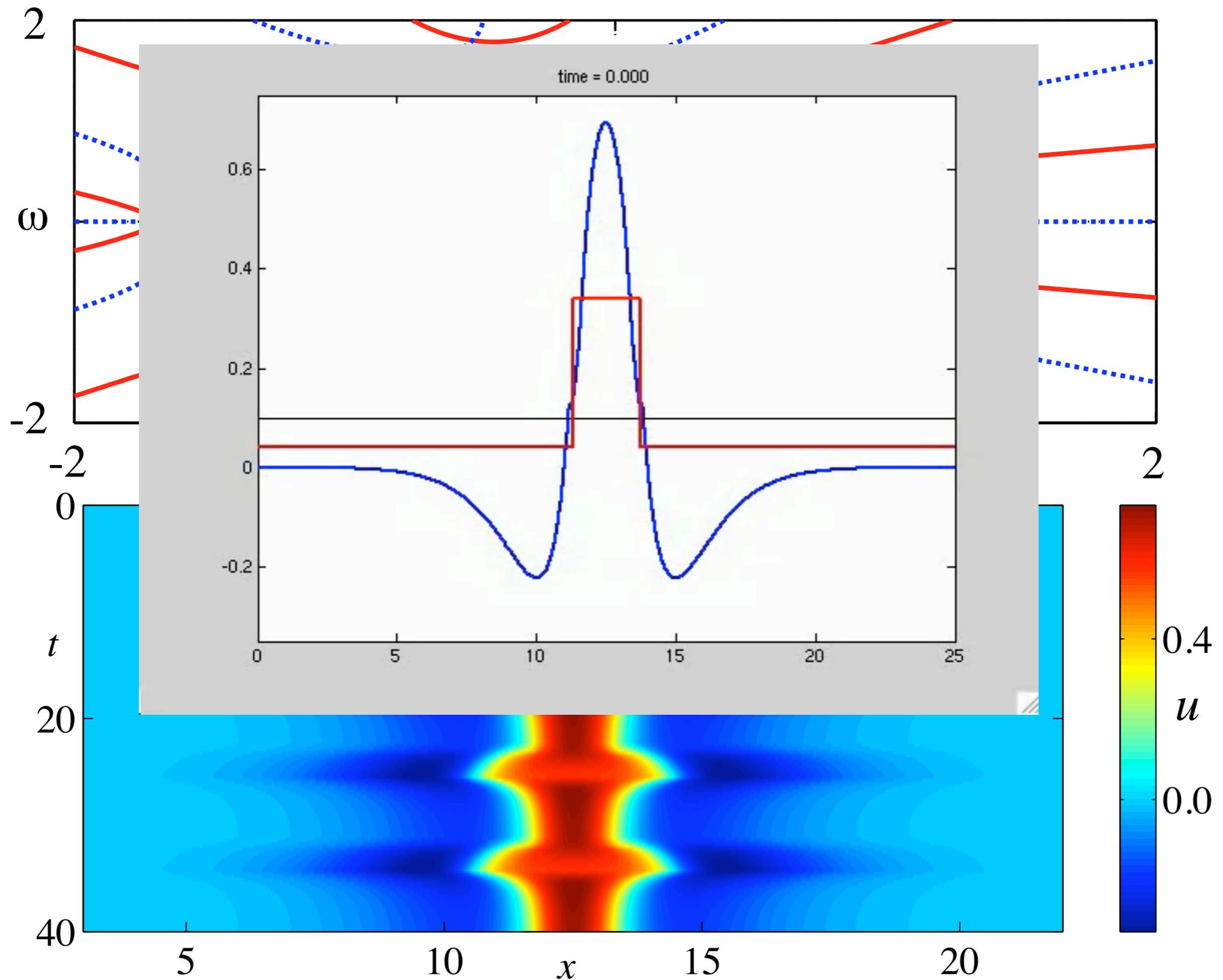
Bump Stability I: $\eta(t) = \alpha^2 t e^{-\alpha t}$

Low κ instability on Re axis (increasing α)

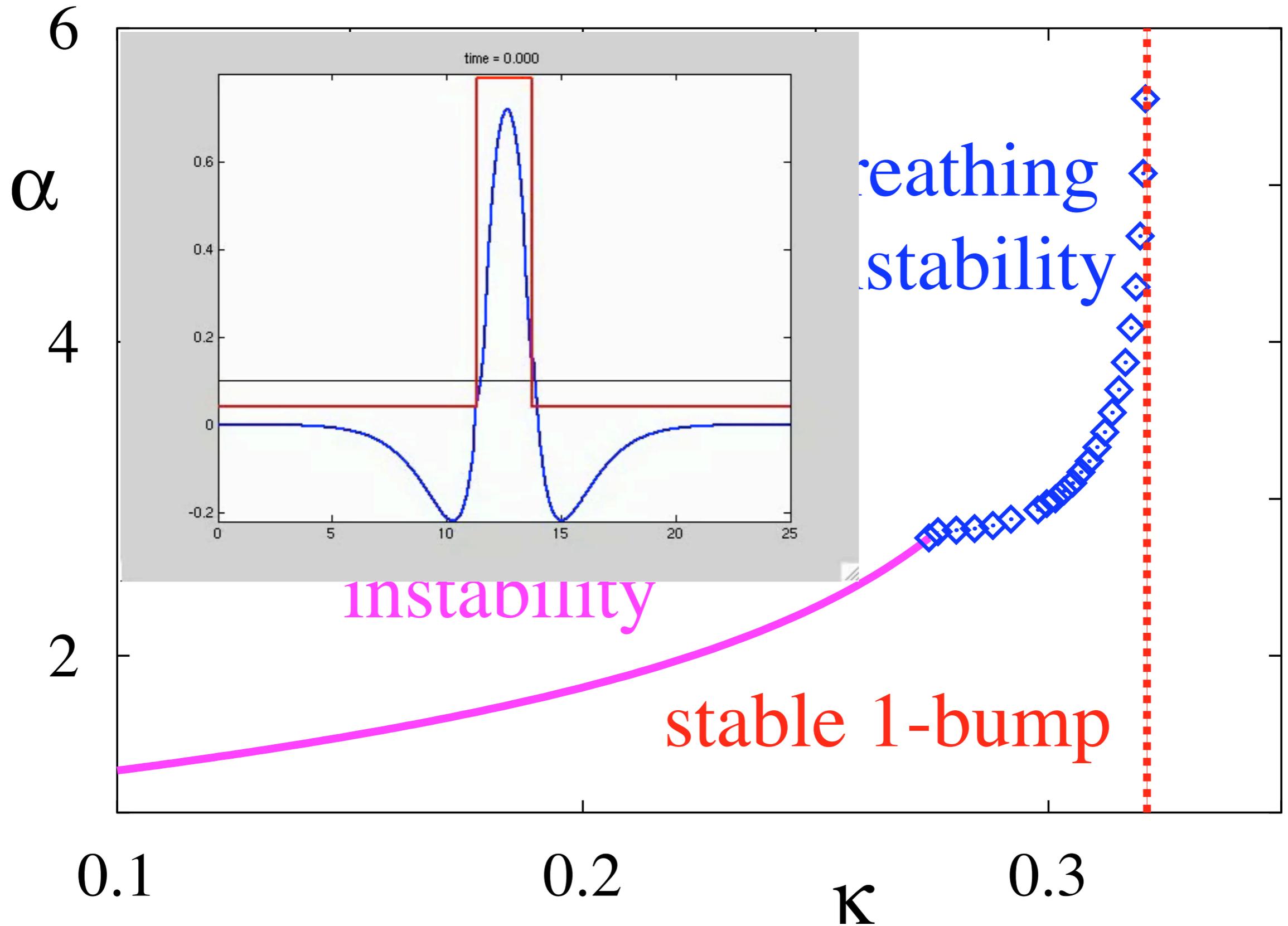


Bump Stability II

High κ instability on Im axis (increasing α) gives a breather

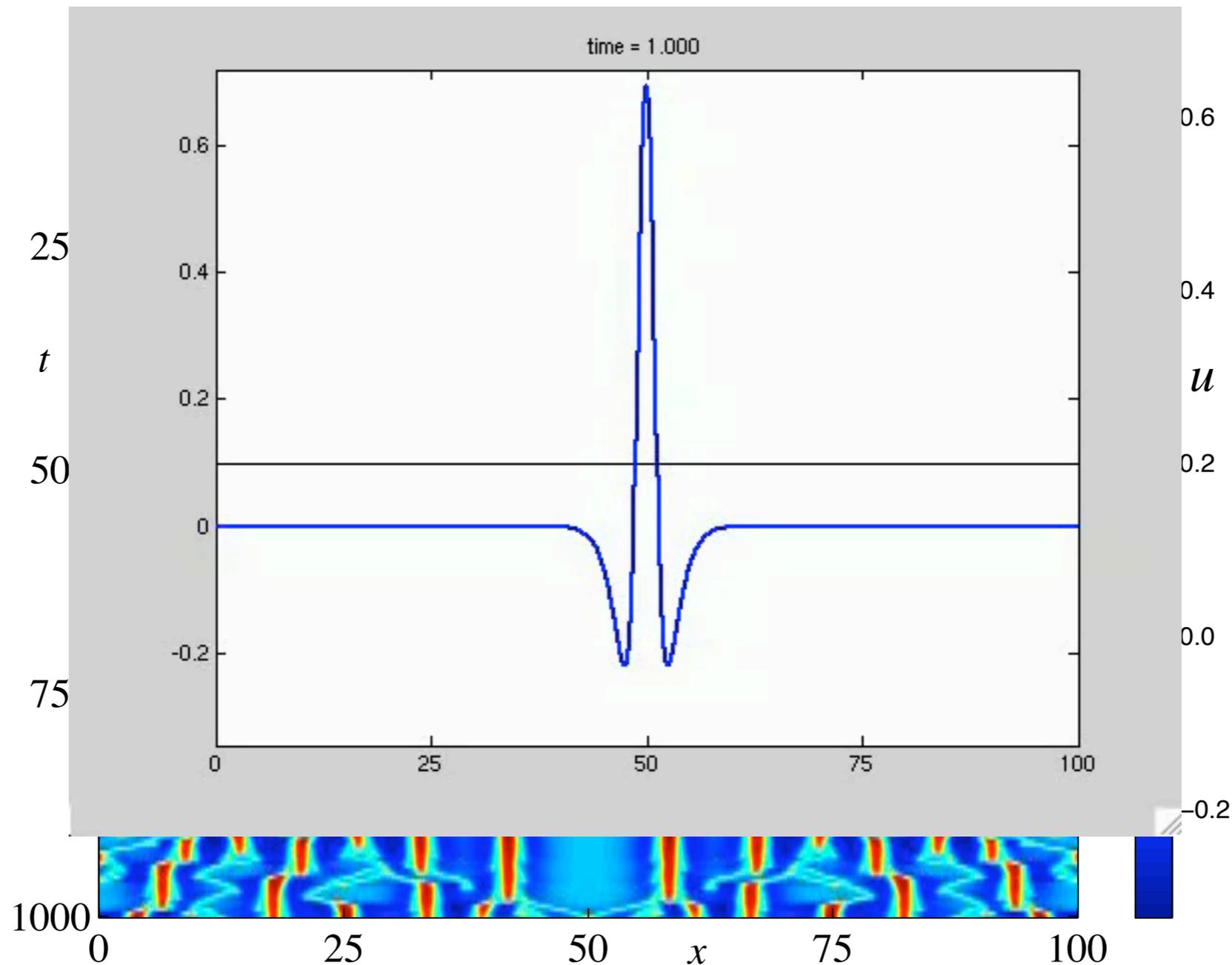


Summary of Bump instabilities



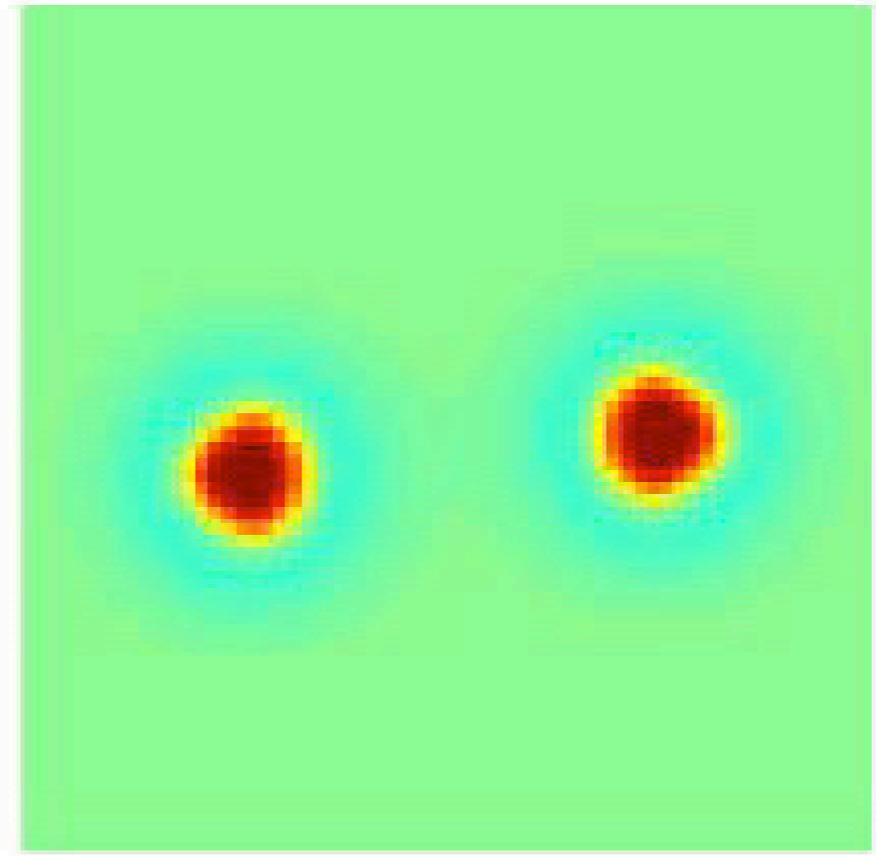
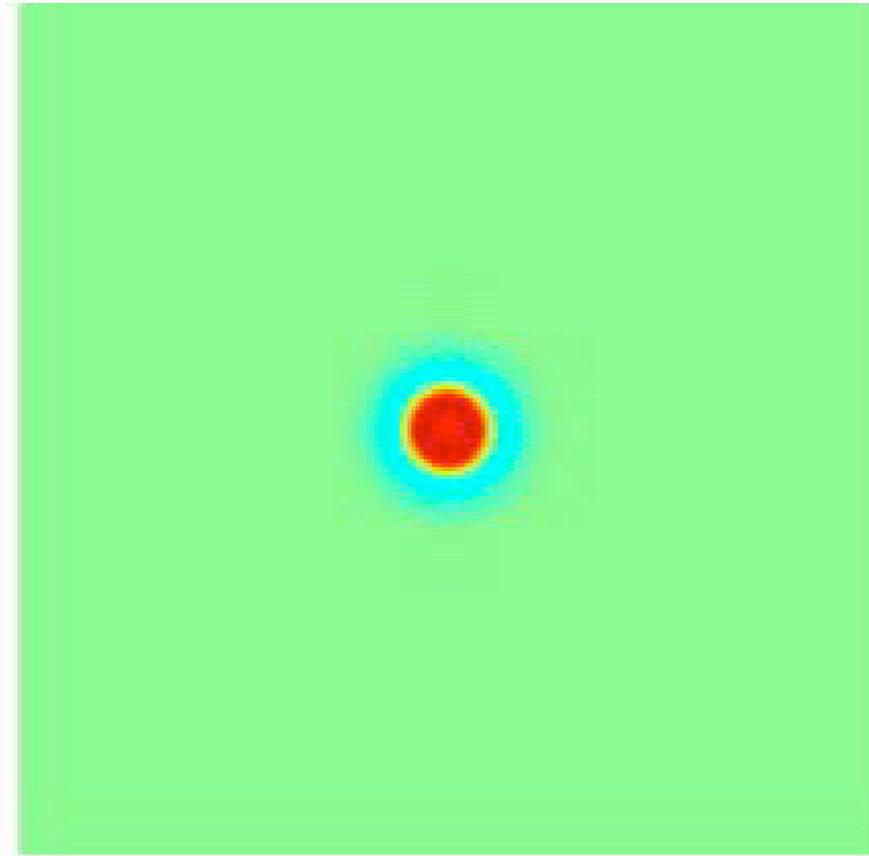
Exotic Dynamics

... including asymmetric breathers, multiple bumps, multiple pulses, periodic traveling waves, and bump-splitting instabilities that appear to lead to spatio-temporal chaos.



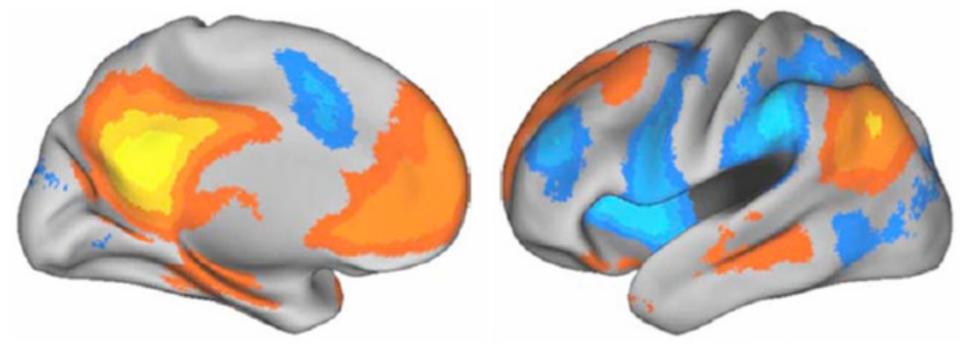
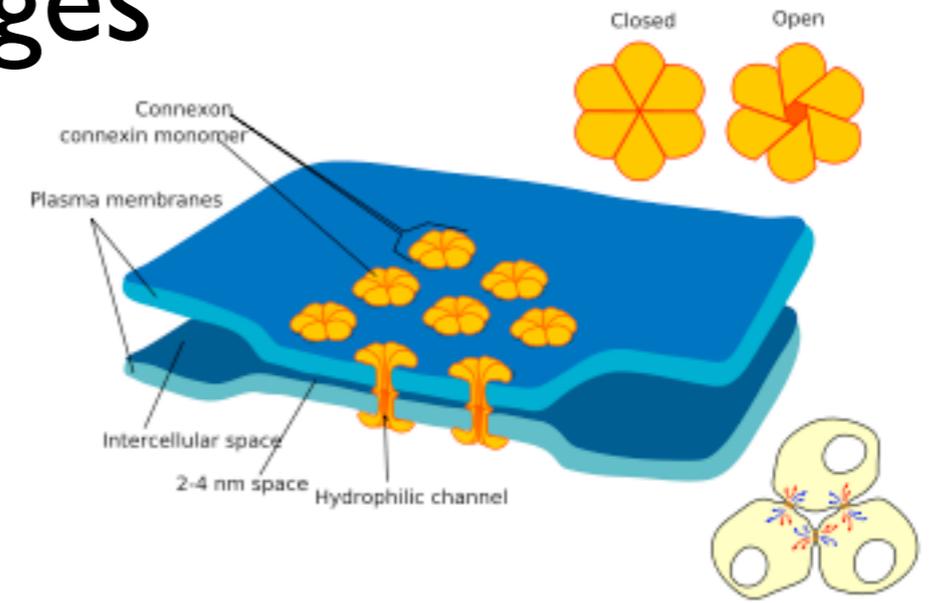
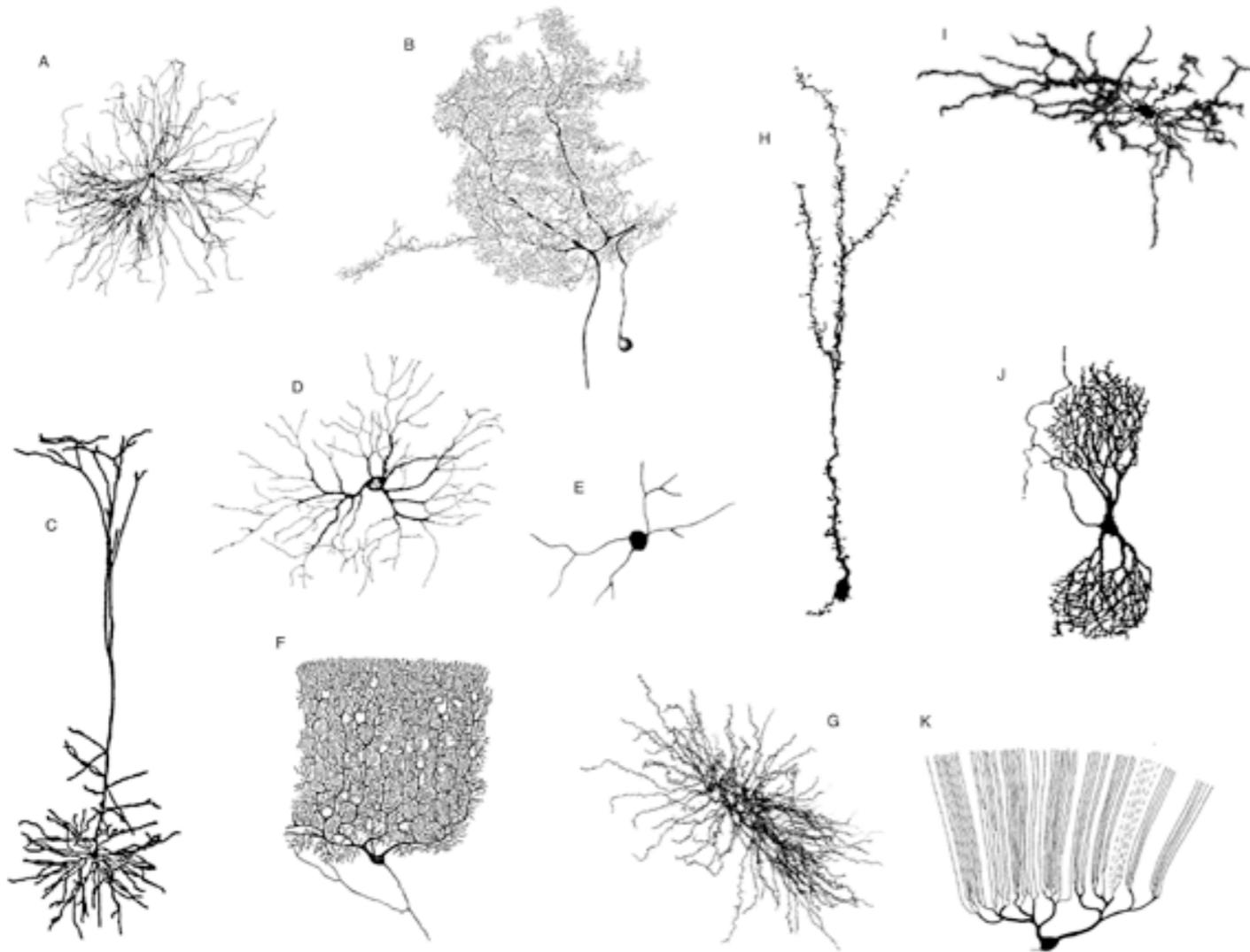
S Coombes and M R Owen: Bumps, breathers and waves in a neural network with spike frequency adaptation. PRL, 94, 148102, (2005).

Splitting and scattering

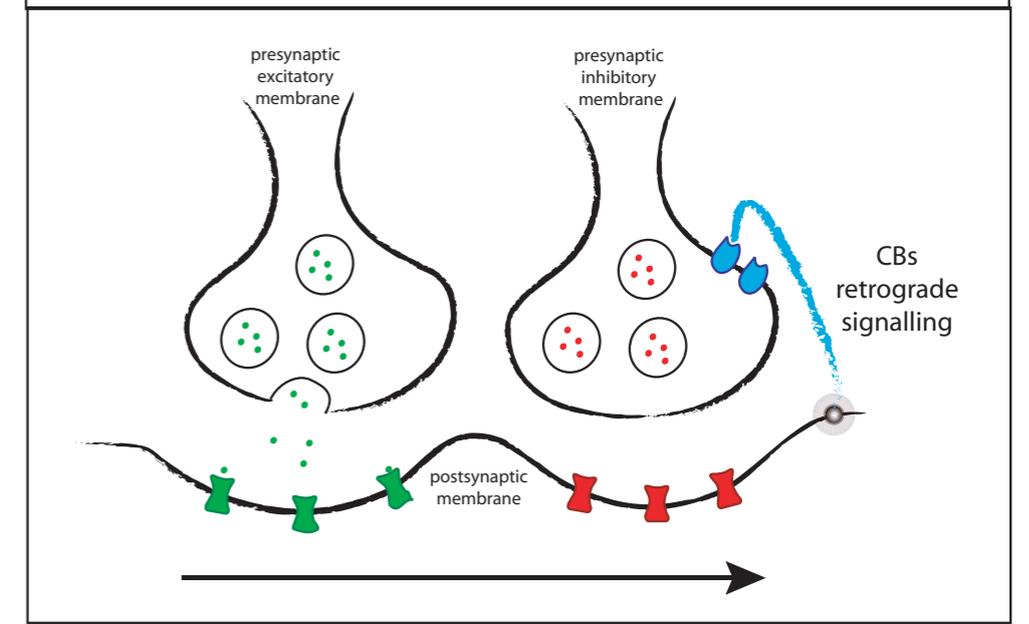
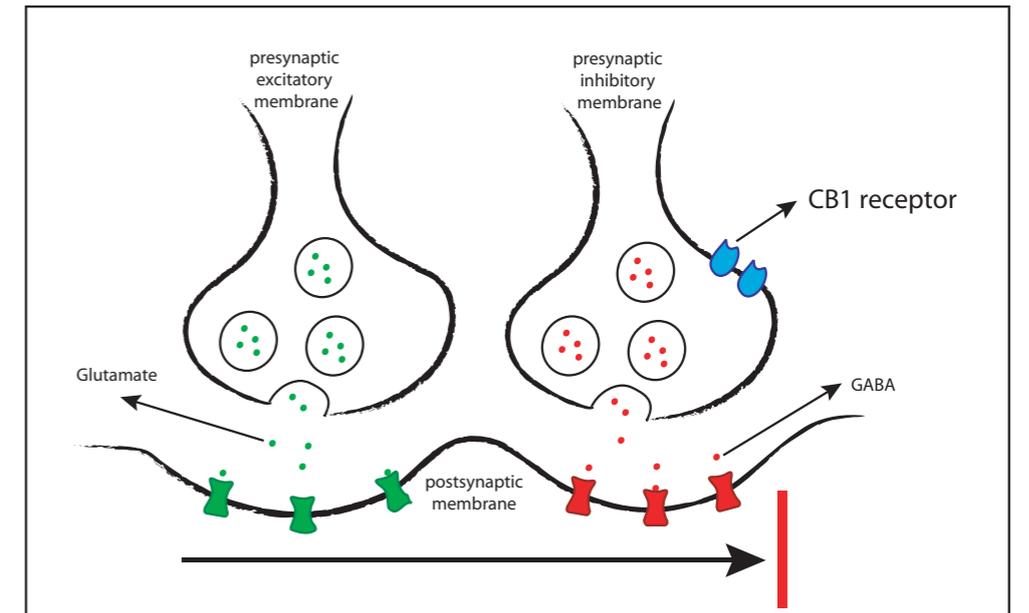


Auto/dispersive solitons as seen in coupled cubic complex Ginzburg-Landau systems and three component reaction-diffusion systems.

Further Challenges



Default mode network and ultra slow coherent oscillations



In collaboration with

Nikola Venkov
(Notts)



Gabriel Lord
(Heriot-Watt)



Yulia Timofeeva
(Warwick)



David Liley
(Melbourne)

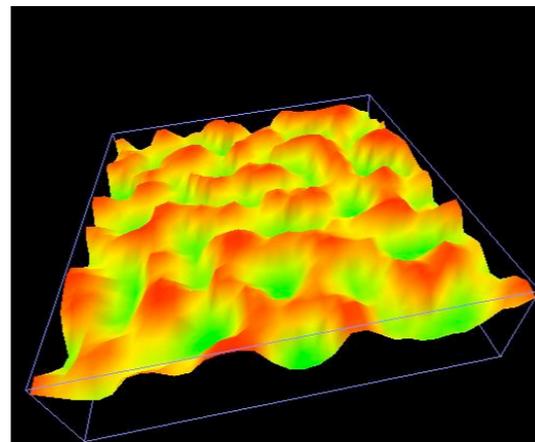


Engineering and Physical Sciences
Research Council

Markus Owen (Notts)



Ingo Bojak (Nijmegen)



Carlo Laing (Massey, NZ)

